

Data modeling and interpretation of interferometric data

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Tone of the talk

- Oriented toward instrumental aspects with important effects on the science data
- Not exhaustive
- Through different examples, gives general methods





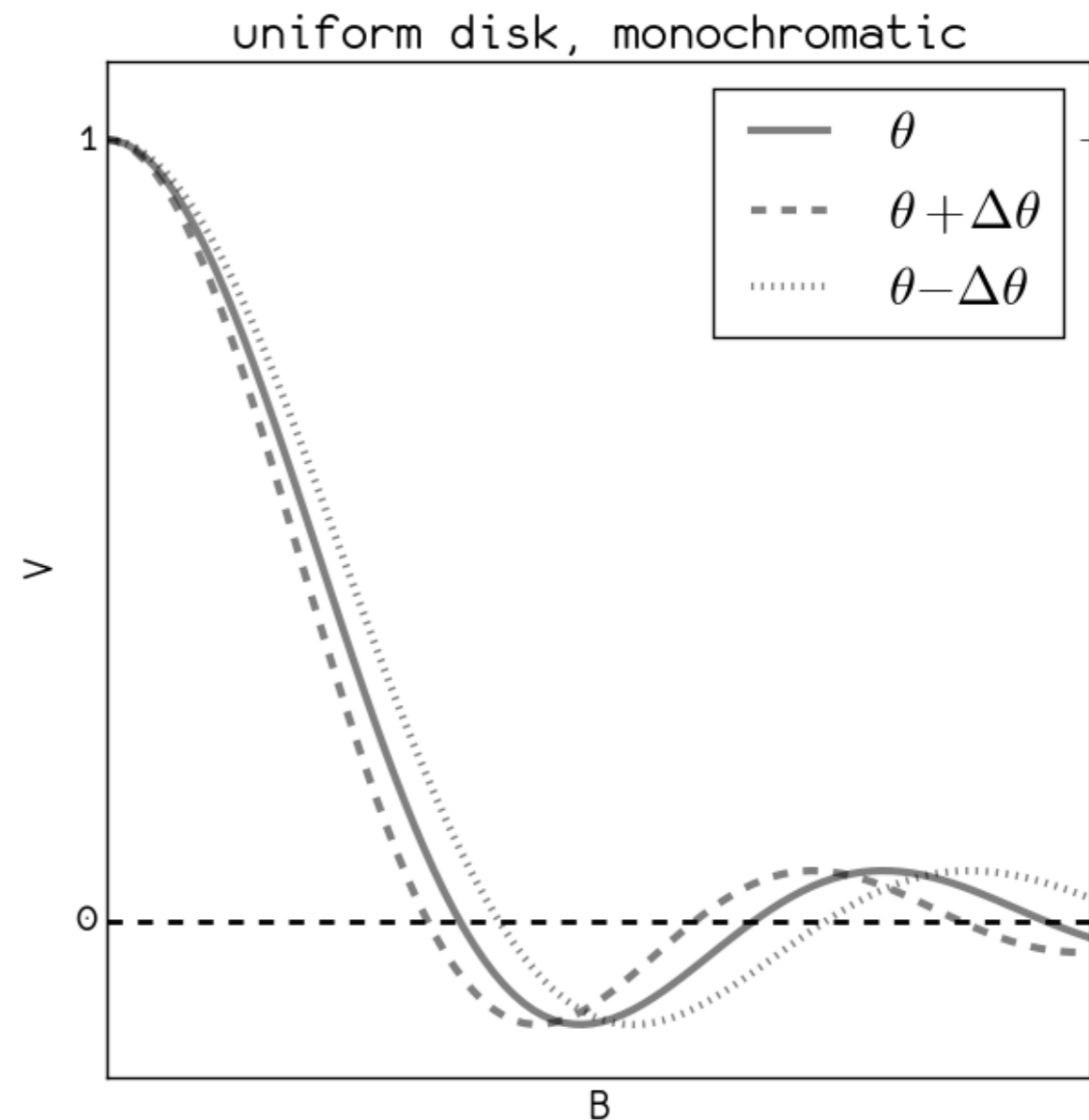
Bandwidth smearing

- The effects of the wavelength bandpass on the interferometric observables
- Most beam combiners have crude spectral resolution to favor sensitivity
- e.g. whole K band, $R = \lambda/\Delta\lambda = 2.2/(2.4-2.0) = 4.4$



our setting

- single baseline interferometer
- we measure only the visibility (amplitude of the fringes)
- we want to measure stellar angular diameters using the first null of the visibility curve



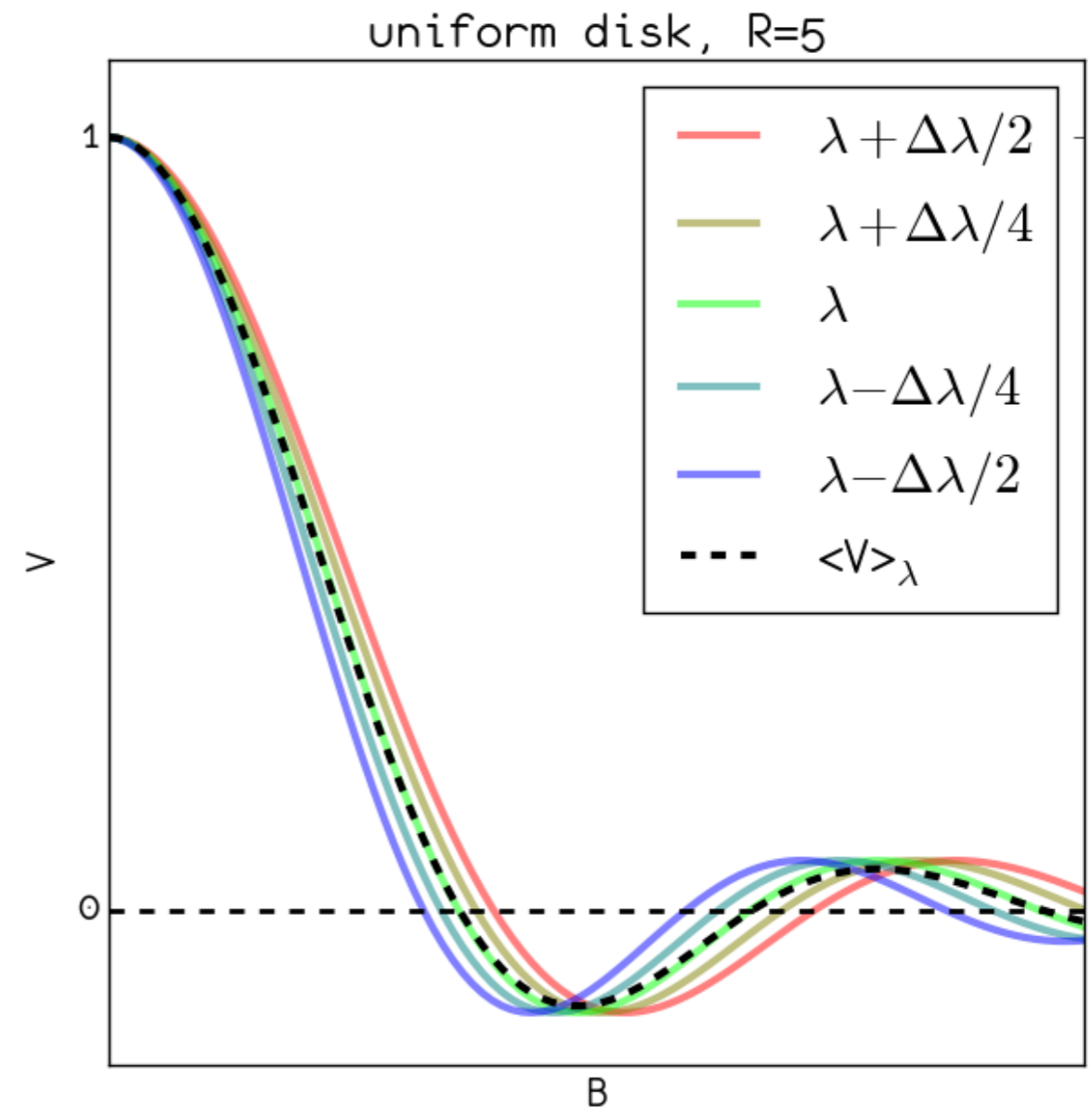
$$V = 2 \frac{J_1(\pi B \theta / \lambda)}{\pi B \theta / \lambda}$$





in broadband

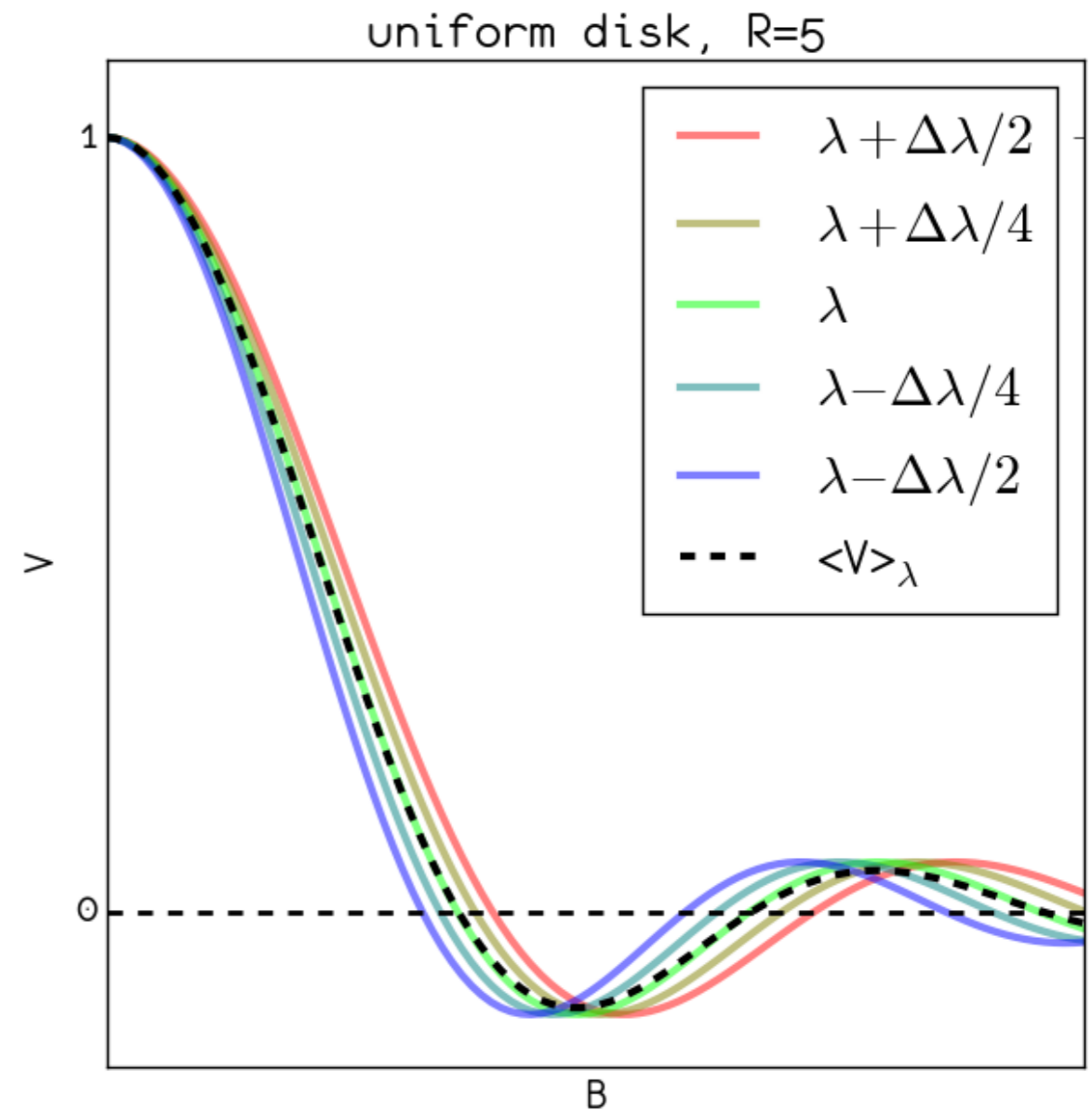
- our beam combiner has $R = \lambda / \Delta\lambda = 5$
- What is the observed visibility in broadband?

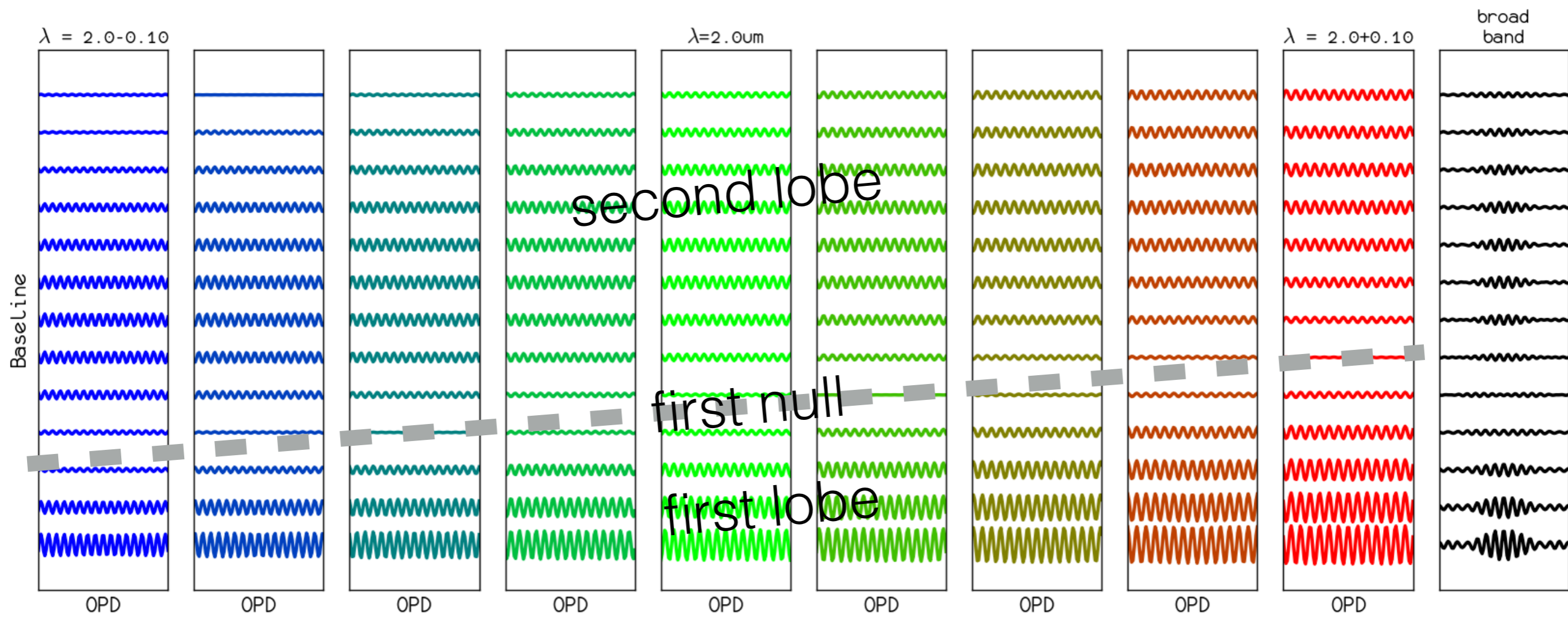


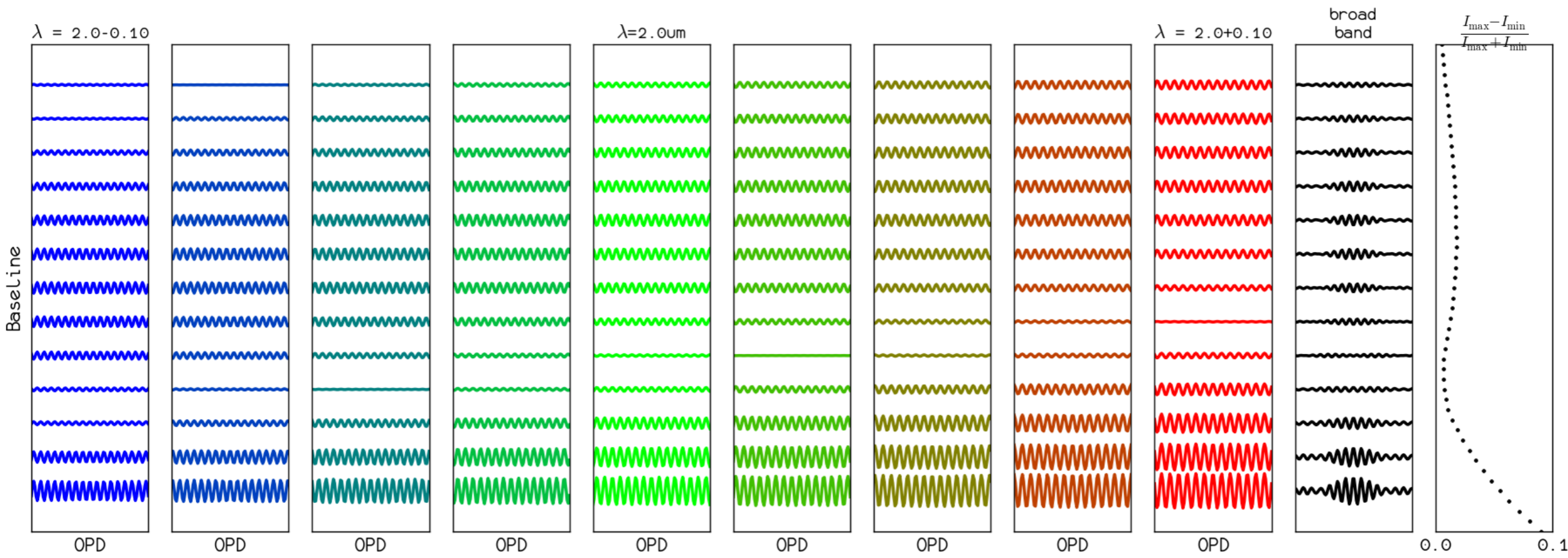


Reasoning

- Beam combiner sees the **sum** of fringes for each wavelength inside the band
- The observed visibility must be the **average** of the visibility in the band

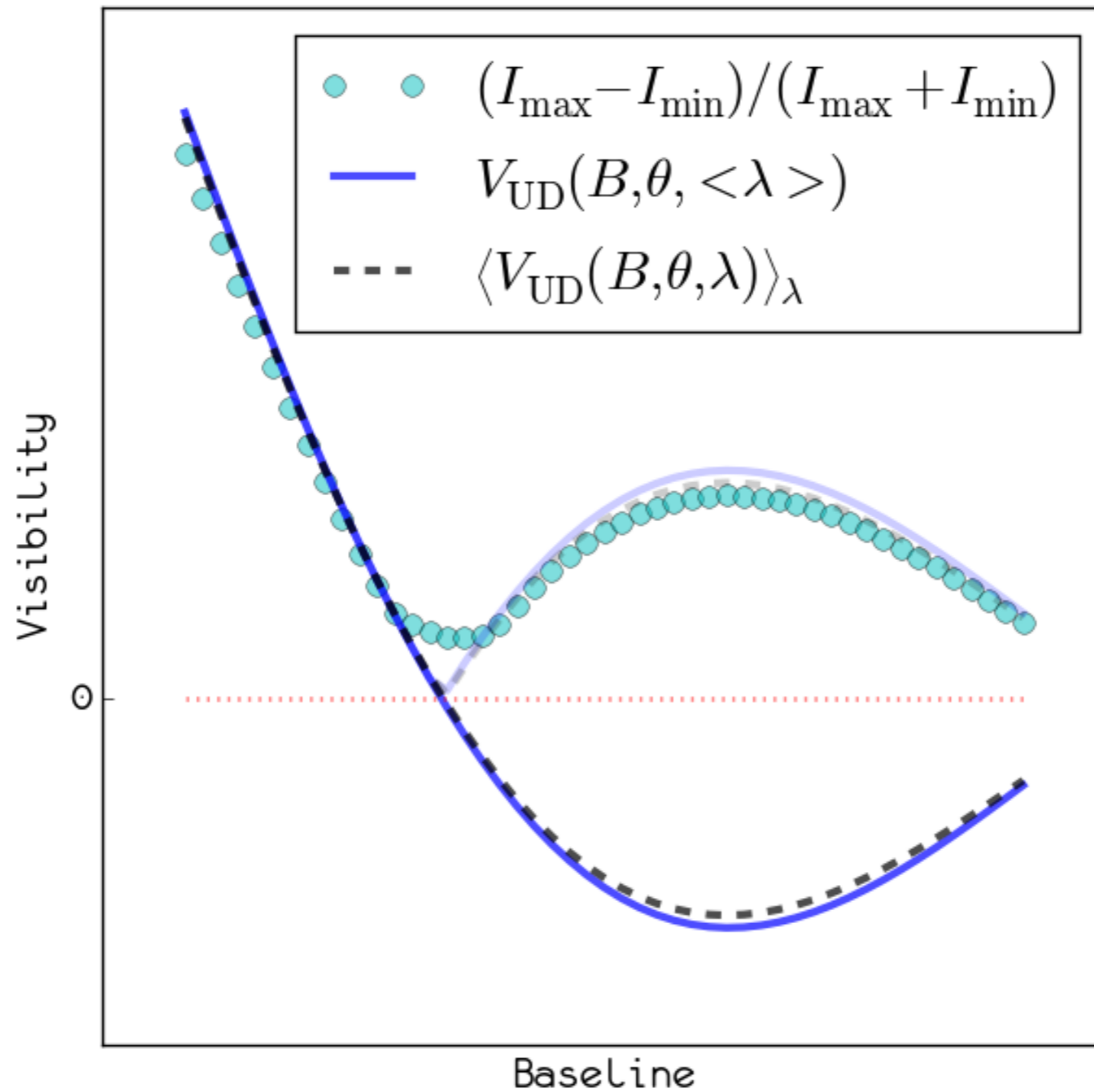






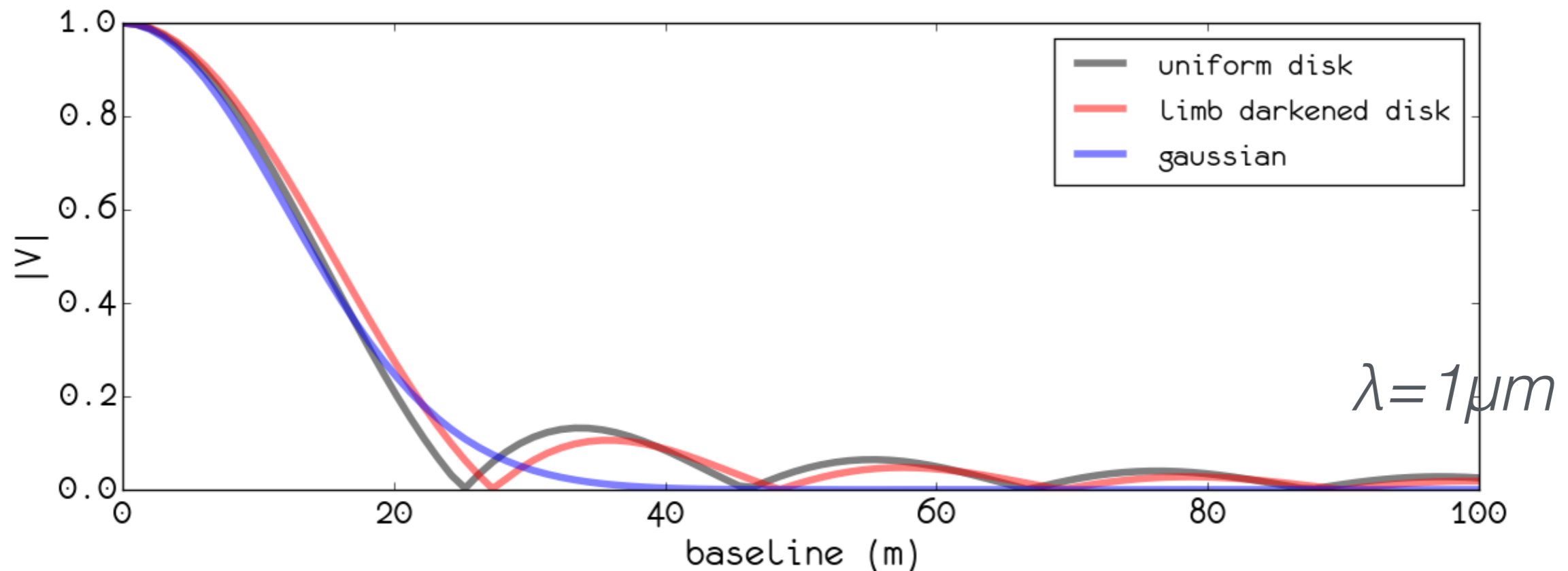
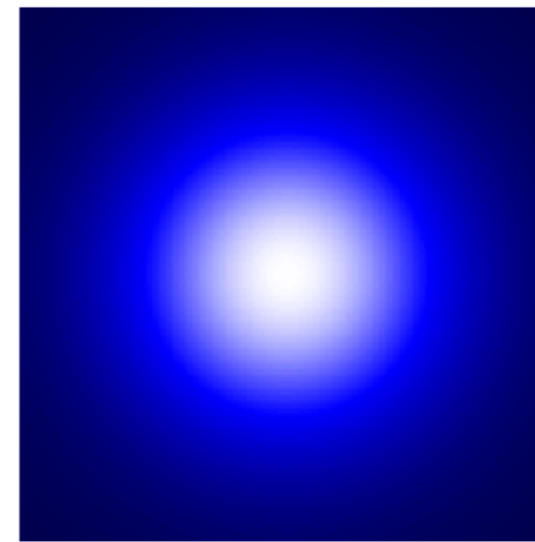
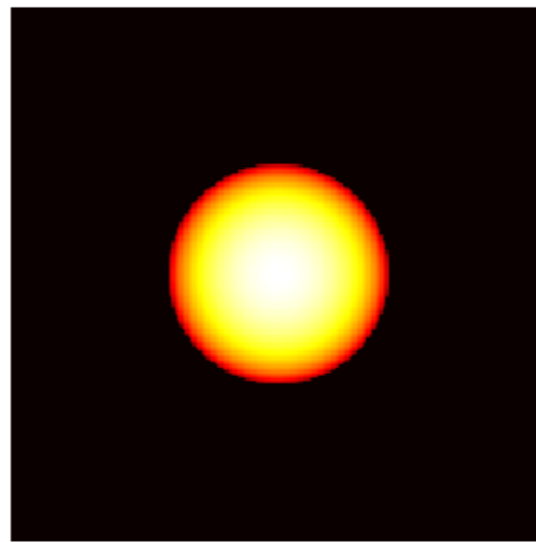
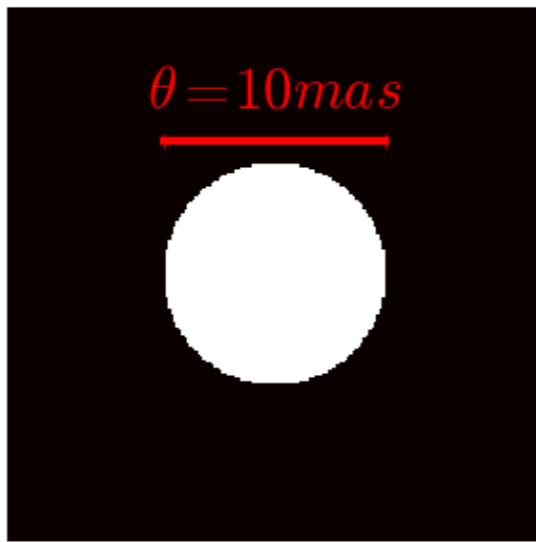


What is going on ?!?

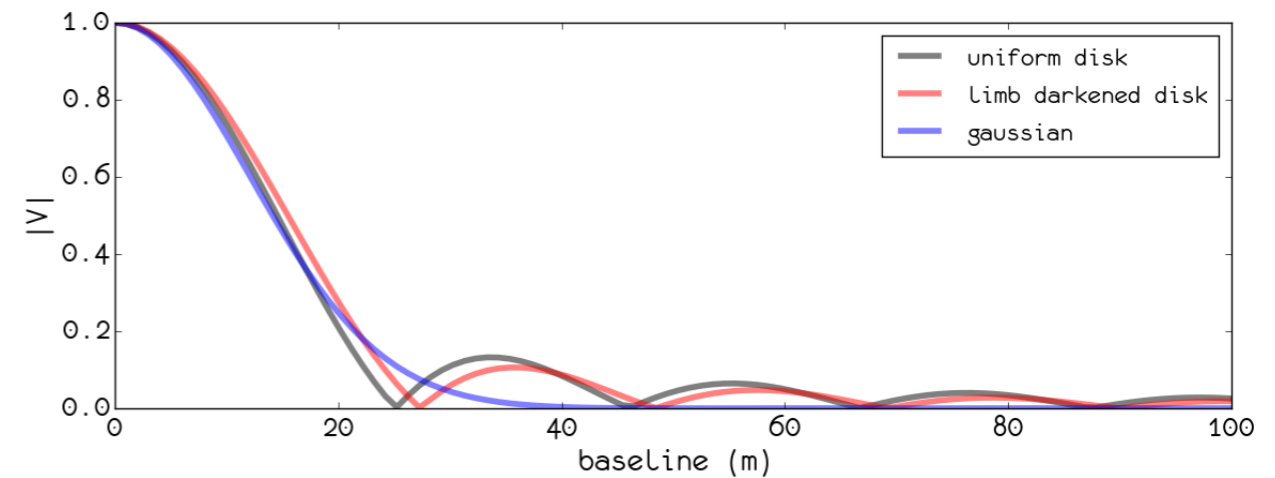
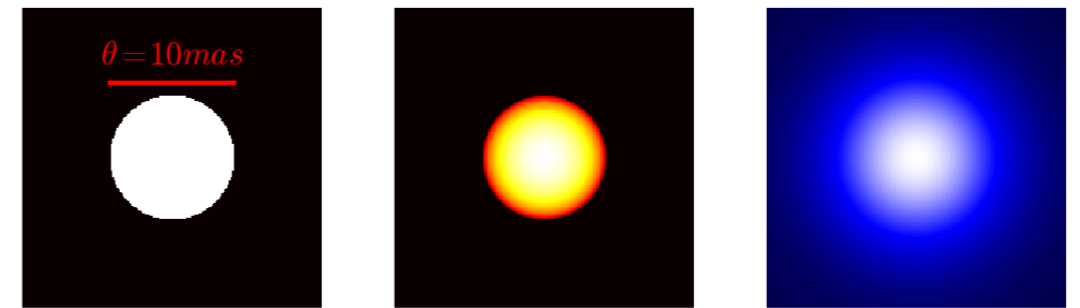


There is information in the $V(B)$ curve

$$V_{\text{UD}}(B, \lambda) = \left| 2 \frac{J_1(x = \pi B \theta / \lambda)}{x} \right| \quad \text{centro-symmetric:} \\ \text{Hankel transform} \quad V(B, \lambda) = \frac{\int I(r, \lambda) J_0(r B / \lambda) r dr}{\int I(r, \lambda) dr}$$



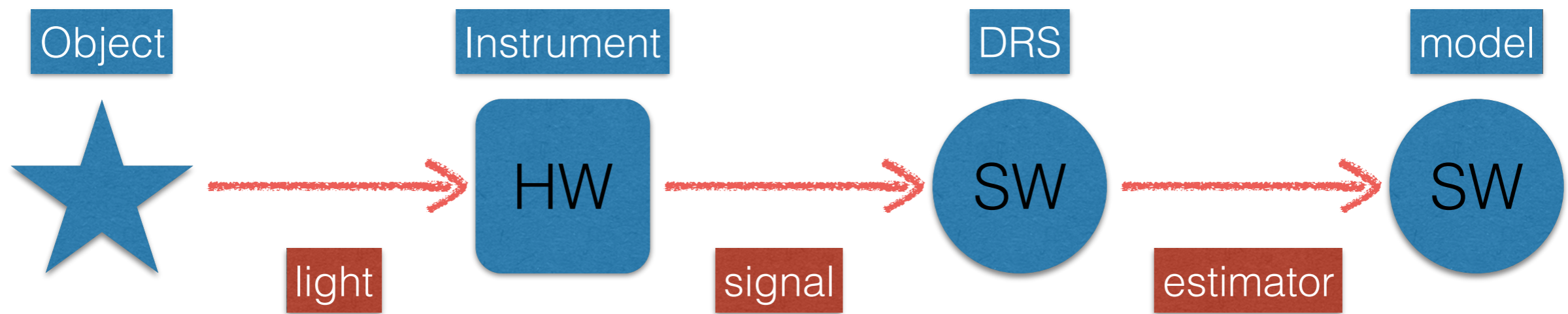
- Measuring diameters == inverting $V(B, \Theta, \lambda)$
- True stars are NOT uniform disks
- limb darkening
 - lowers the visibility lobes
 - bias the diameter measurements





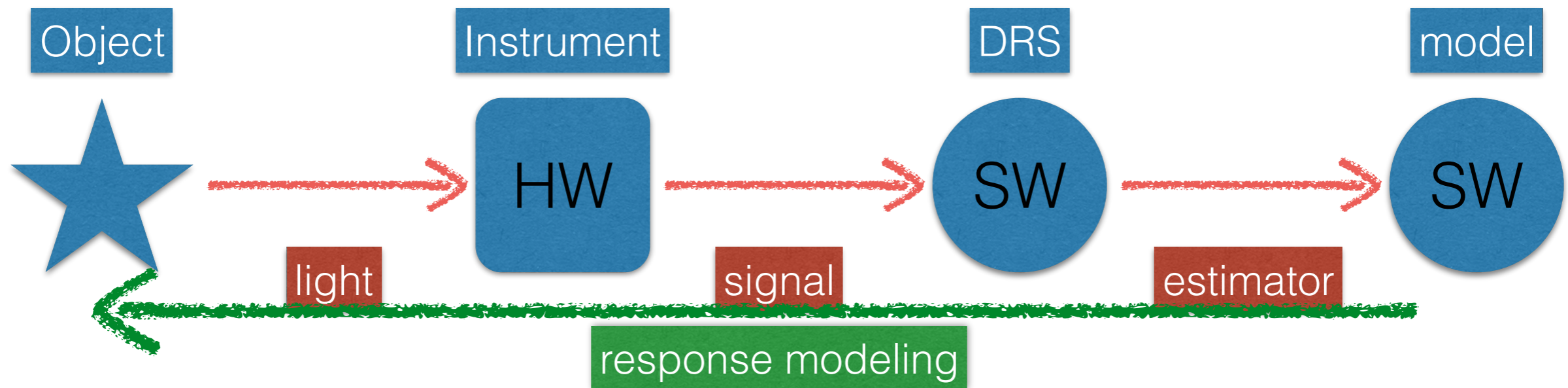
General considerations

- The instrument does not observe visibility, it observes **fringes**
- An **estimator** is used to derived the visibility from the fringes using a data reduction software (DRS)





The correct approach



The numerical model should match

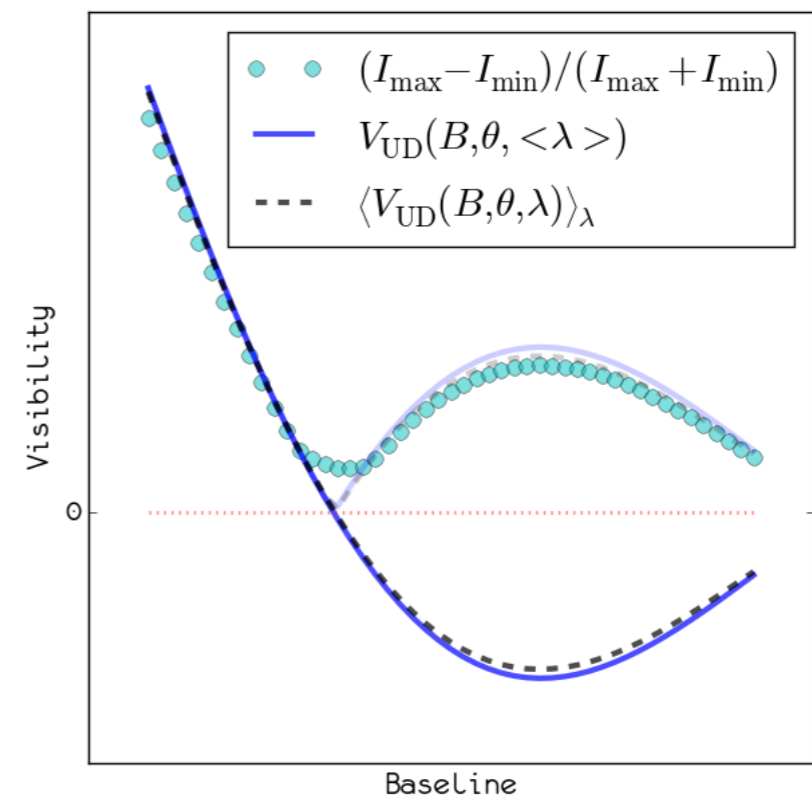
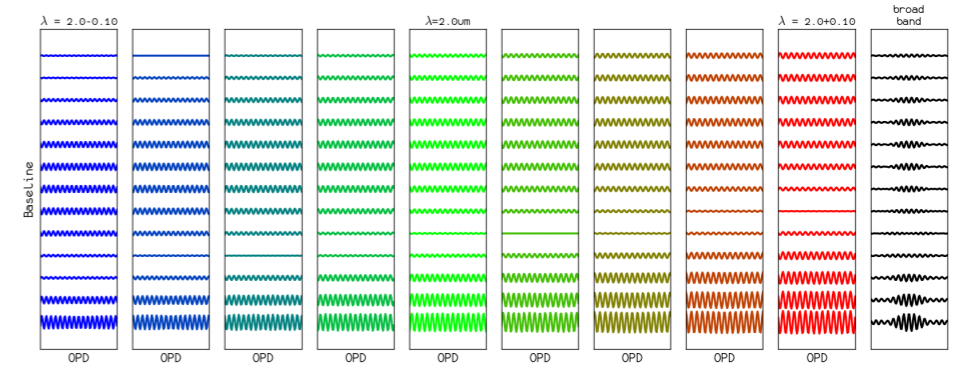
- the estimator
- the instrumental characteristics
- the object characteristics





we were on the right direction...

- we **synthesized** a signal using an instrumental characteristic: bandwidth smearing
- we used an **estimator** of the visibility:
 $\rightarrow V \sim (I_{\max} - I_{\min}) / (I_{\max} + I_{\min})$



better yet, do it analytically



Let's start all over again...

- What is the visibility estimator?
- What instrument's characteristics should I take into account?
- What object's characteristics should I take into account?





Visibility estimator

- visibility has additive noise: $V + n$

- we measure fringe's contrast $\mu = |V+n|$

- averaging:

- $\langle \mu \rangle = \langle |V+n| \rangle$

- $\langle \mu \rangle$ is biased

- what about $\langle \mu^2 \rangle$?

- $\langle \mu^2 \rangle = \langle |V+n|^2 \rangle$

- $\langle \mu^2 \rangle = \langle |V|^2 \rangle + \langle 2\text{Re}\{Vn\} \rangle + \langle |n|^2 \rangle$

- assuming V and n are uncorrelated:
 $\langle \text{Re}\{Vn\} \rangle = 0$

- $\langle \mu^2 \rangle = \langle |V|^2 \rangle + \langle |n|^2 \rangle$

- $\langle \mu^2 \rangle$ is biased but can be unbiased if $\langle |n|^2 \rangle$ is estimated





Fourier estimator

- Remember Parseval's identity?

- $$\int_{-\infty}^{\infty} |f(t)|^2 dt = \int_{-\infty}^{\infty} |TF(f)(\sigma)|^2 d\sigma$$

- The average squared amplitude of the signal ==
The average PSD



Analytical fringe signal

Fringes signal $F(\delta)$:
function of
monochromatic
fringes $f(\delta, \lambda)$,
*function of OPD (δ)
and wavelength (λ)*

$$F(\delta) = \frac{\int_0^\infty B(\lambda)T(\lambda)f(\delta, \lambda)\lambda d\lambda}{\int_0^\infty B(\lambda)T(\lambda)\lambda d\lambda}$$

$B(\lambda) \rightarrow$ stellar spectrum

$T(\lambda) \rightarrow$ instrumental transmission

$$f(\delta, \lambda) = 1 + \text{Re} \left(V(\lambda) e^{-2\pi\delta/\lambda} \right)$$

$\lambda \rightarrow$ photon detection

$$PSD(\sigma) = |FT_{\delta}[F(\delta)]|^2$$

$$= \left| FT_{\delta} \left[\frac{\int_0^{\infty} B(\lambda)T(\lambda)f(\delta, \lambda)\lambda d\lambda}{\int_0^{\infty} B(\lambda)T(\lambda)\lambda d\lambda} \right] \right|^2$$

linearity
of FT

$$= \frac{\left| \int_0^{\infty} B(\lambda)T(\lambda)\tilde{f}(\delta, \lambda)\lambda d\lambda \right|^2}{\left| \int_0^{\infty} B(\lambda)T(\lambda)\lambda d\lambda \right|^2}$$

frequency
selection

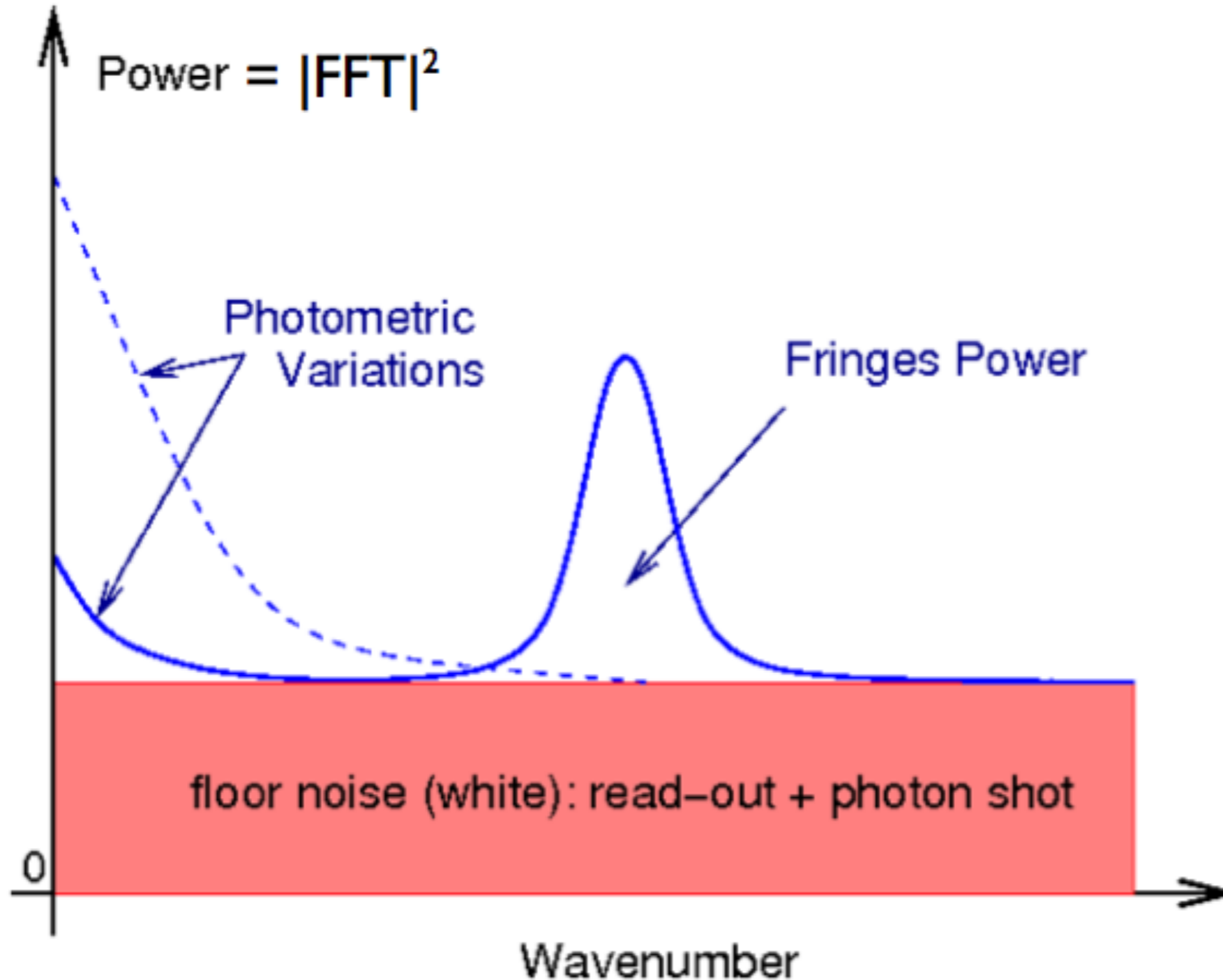
$$= \frac{\left| \int_0^{\infty} B(\lambda)T(\lambda)V(\lambda)\delta_{1/\sigma}(\lambda)\lambda d\lambda \right|^2}{\left| \int_0^{\infty} B(\lambda)T(\lambda)\lambda d\lambda \right|^2}$$

$$= \frac{[B(1/\sigma)T(1/\sigma)V(1/\sigma)1/\sigma]^2}{\left| \int_0^{\infty} B(\lambda)T(\lambda)\lambda d\lambda \right|^2}$$

Normalised, frequency
averaged
object's visibility

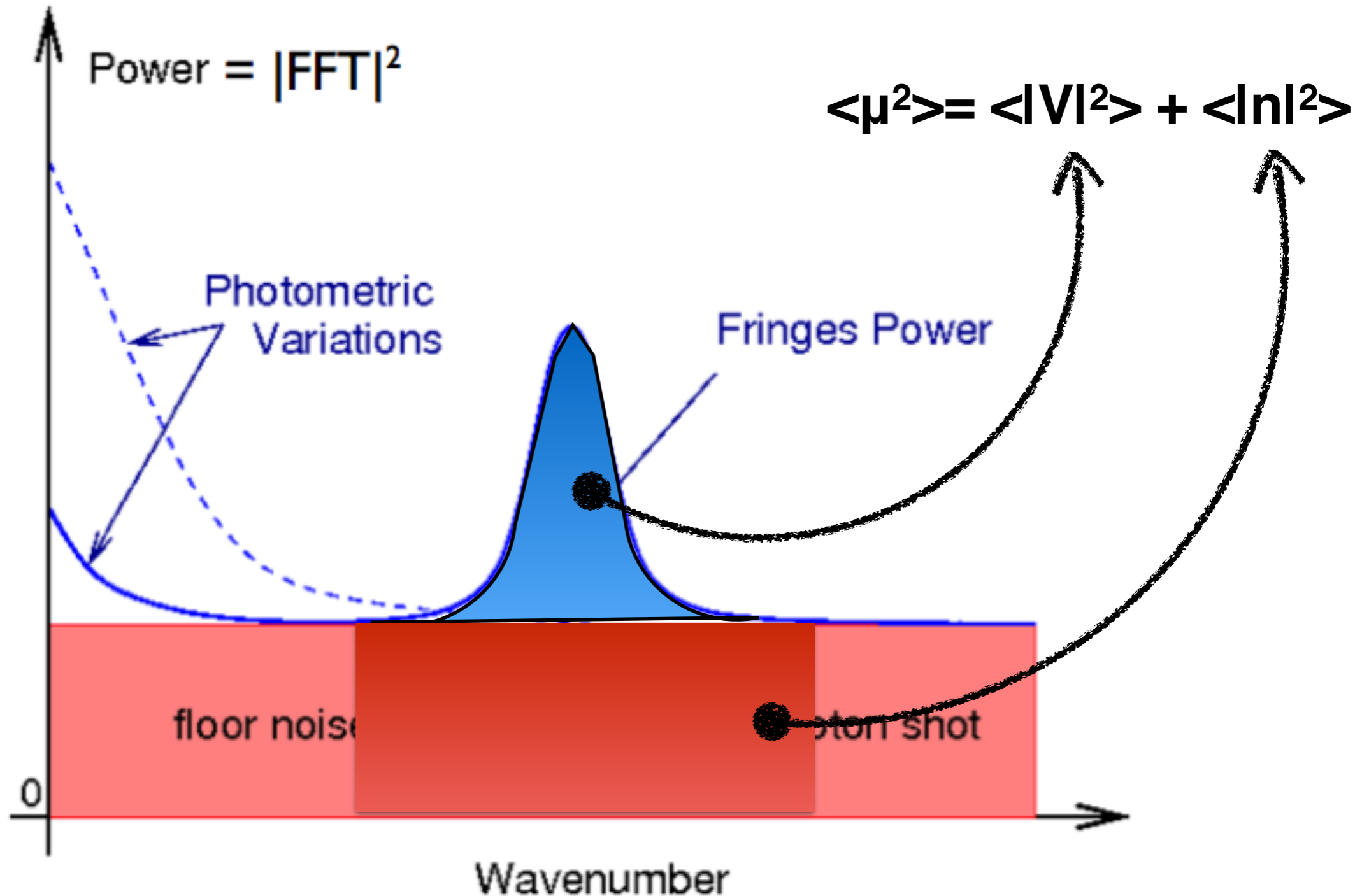


Real signal





Real signal

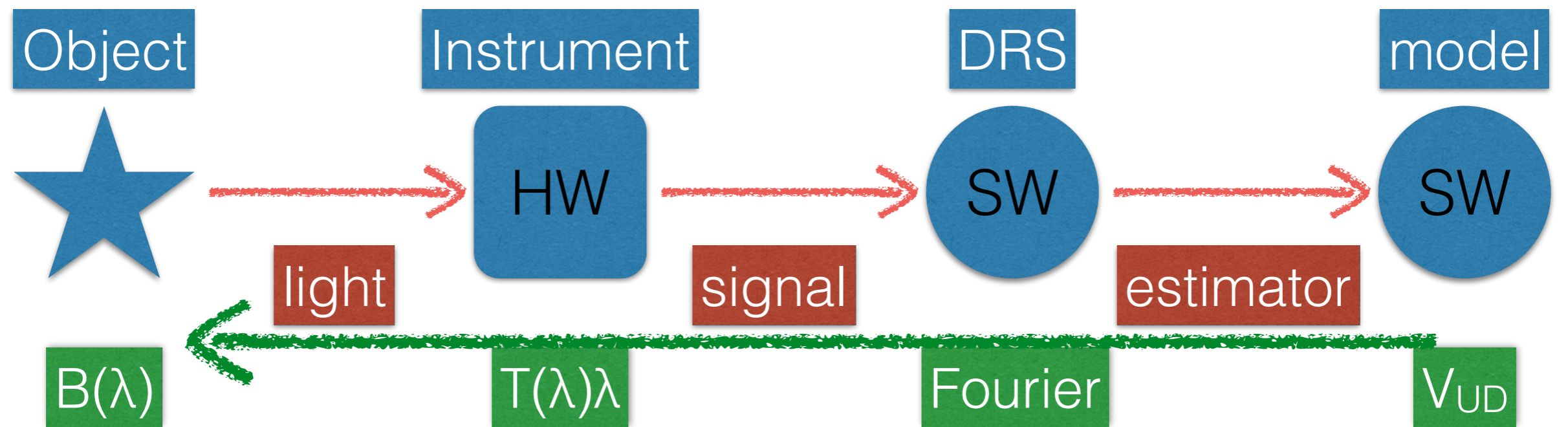




Fourier estimator

$$V_{\text{measured}}^2 \propto \int_{\sigma_{\text{min}}}^{\sigma_{\text{max}}} B(1/\sigma)^2 T(1/\sigma)^2 V(1/\sigma)^2 1/\sigma^2 d\sigma$$

- weighted average of the squared visibilities
- function of the object spectrum
- function of the instrumental transmission

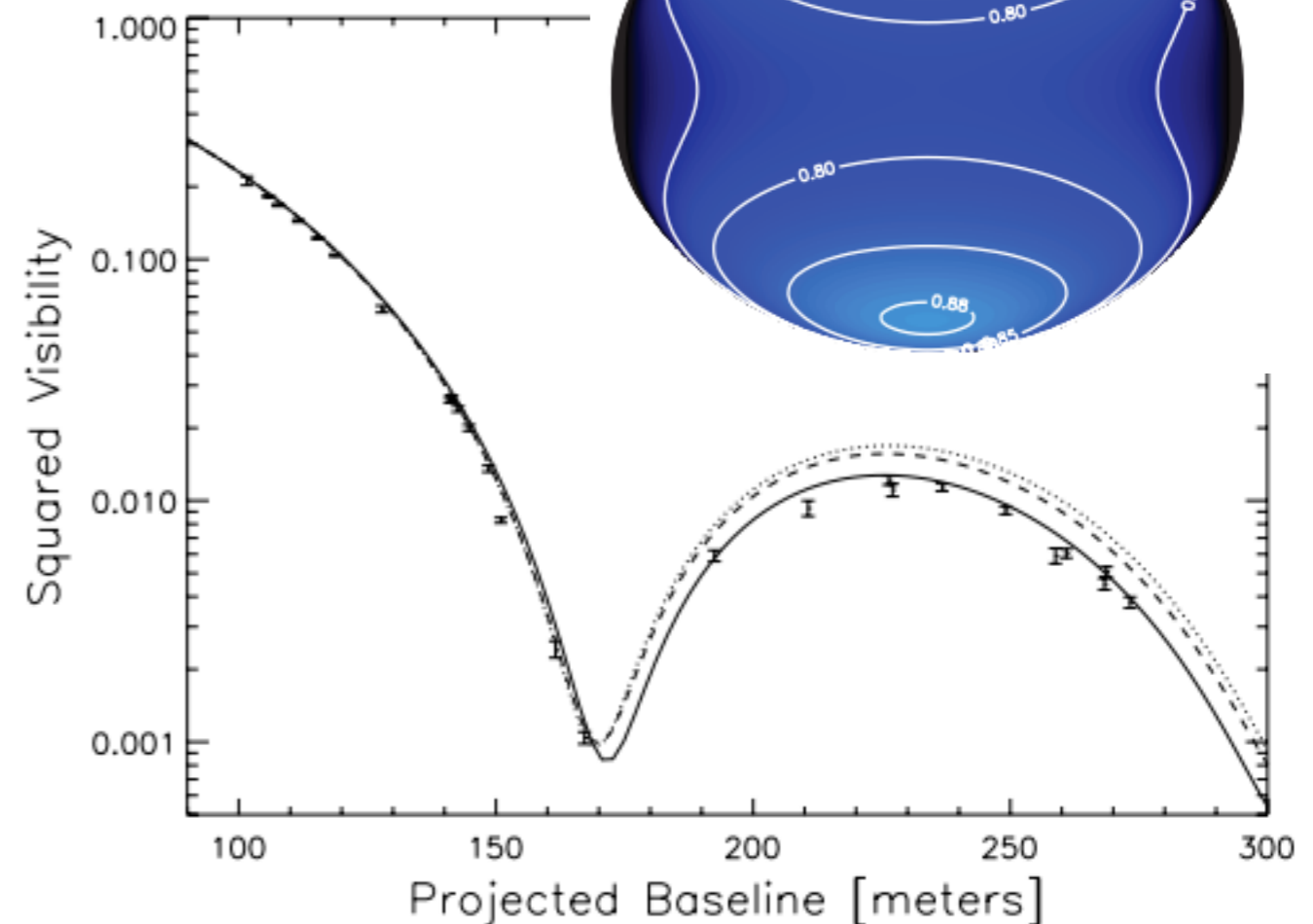
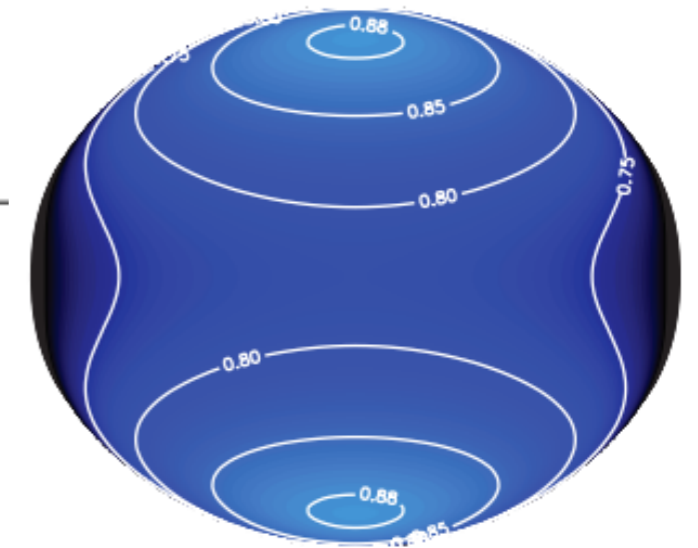
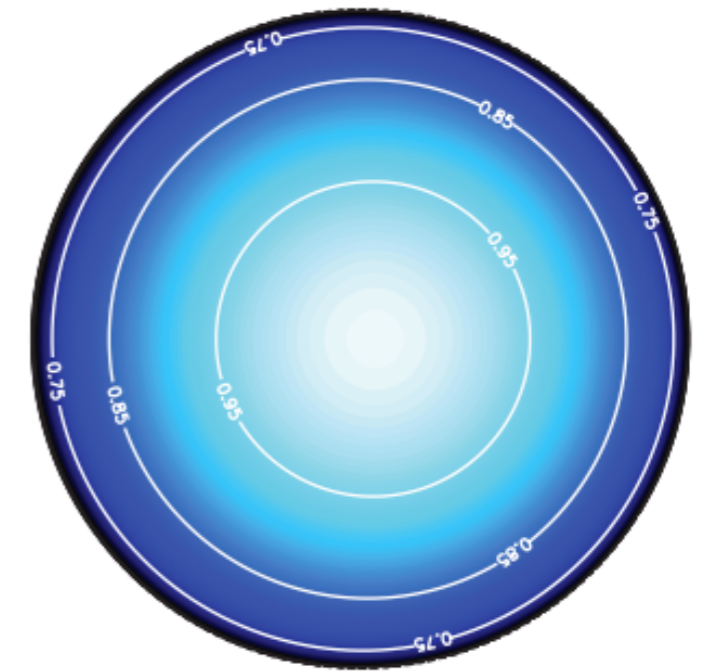


$$V_{\text{measured}}^2(\theta, B) \propto \int_{\sigma_{\text{min}}}^{\sigma_{\text{max}}} B(1/\sigma)^2 T(1/\sigma)^2 V_{UD}(B, \theta, 1/\sigma)^2 1/\sigma^2 d\sigma$$



Fast Rotating stars

- Aufdenberg+ 2006
- Observations of the star Vega with FLUOR@CHARA
- Accurate modelling allowed to prove that the star is a rapid rotator seen pole-on
- **astrophysical effect ~ bandwidth smearing effects**





Why it is important

- Interferometric observations lead to visibilities, closure phases (+ differential quantities)
- **Images** can be reconstructed...
- ... But the astrophysical **quantitative** results will always be derived from visibilities



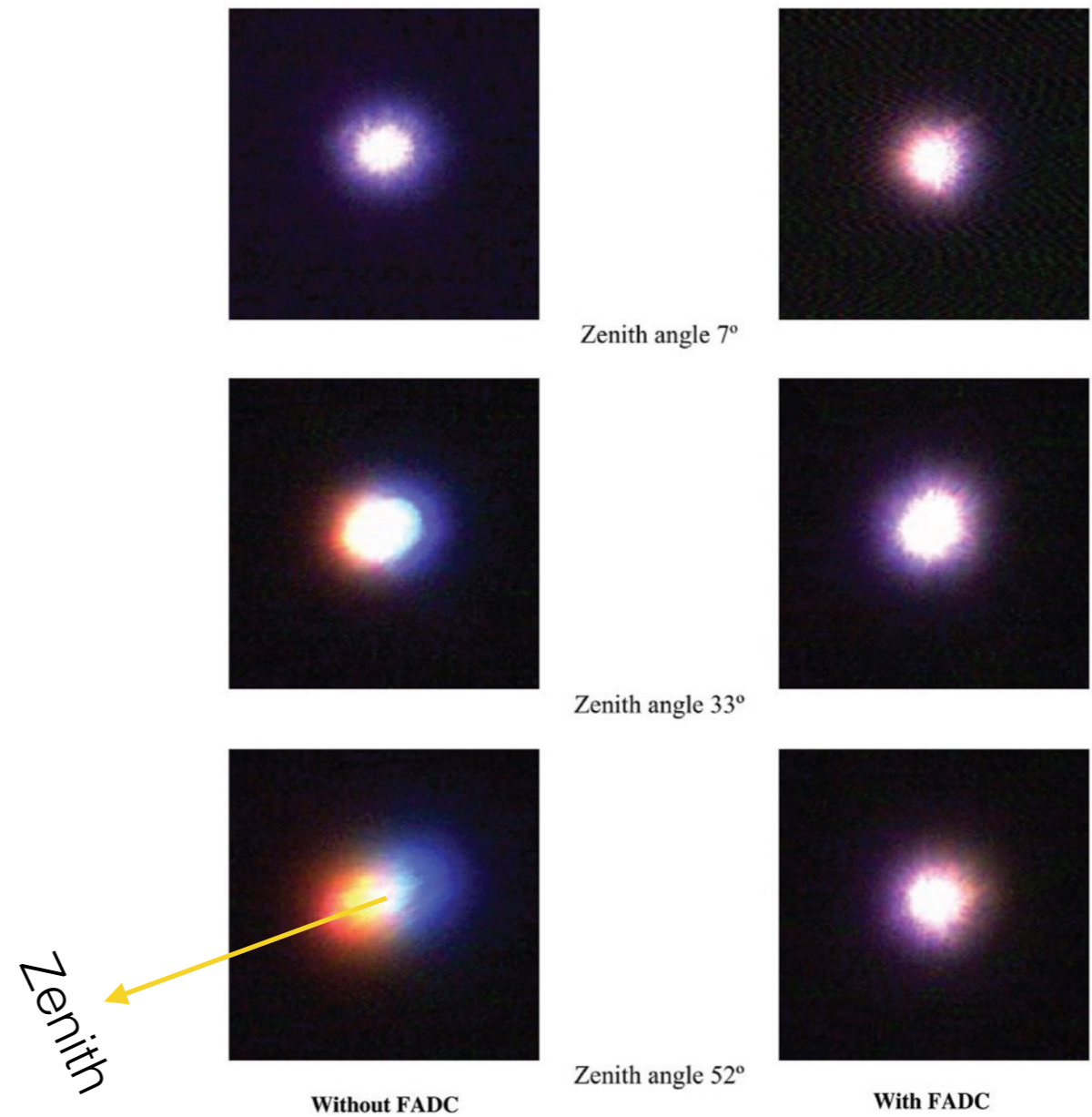
When instrumental effects mimic astrophysical signal

- We have seen “obvious” effects: model disagree with the observations
- Some effects are more difficult to spot!
- Some signal:
 - Astrophysical phenomenon?
 - instrumental effect?



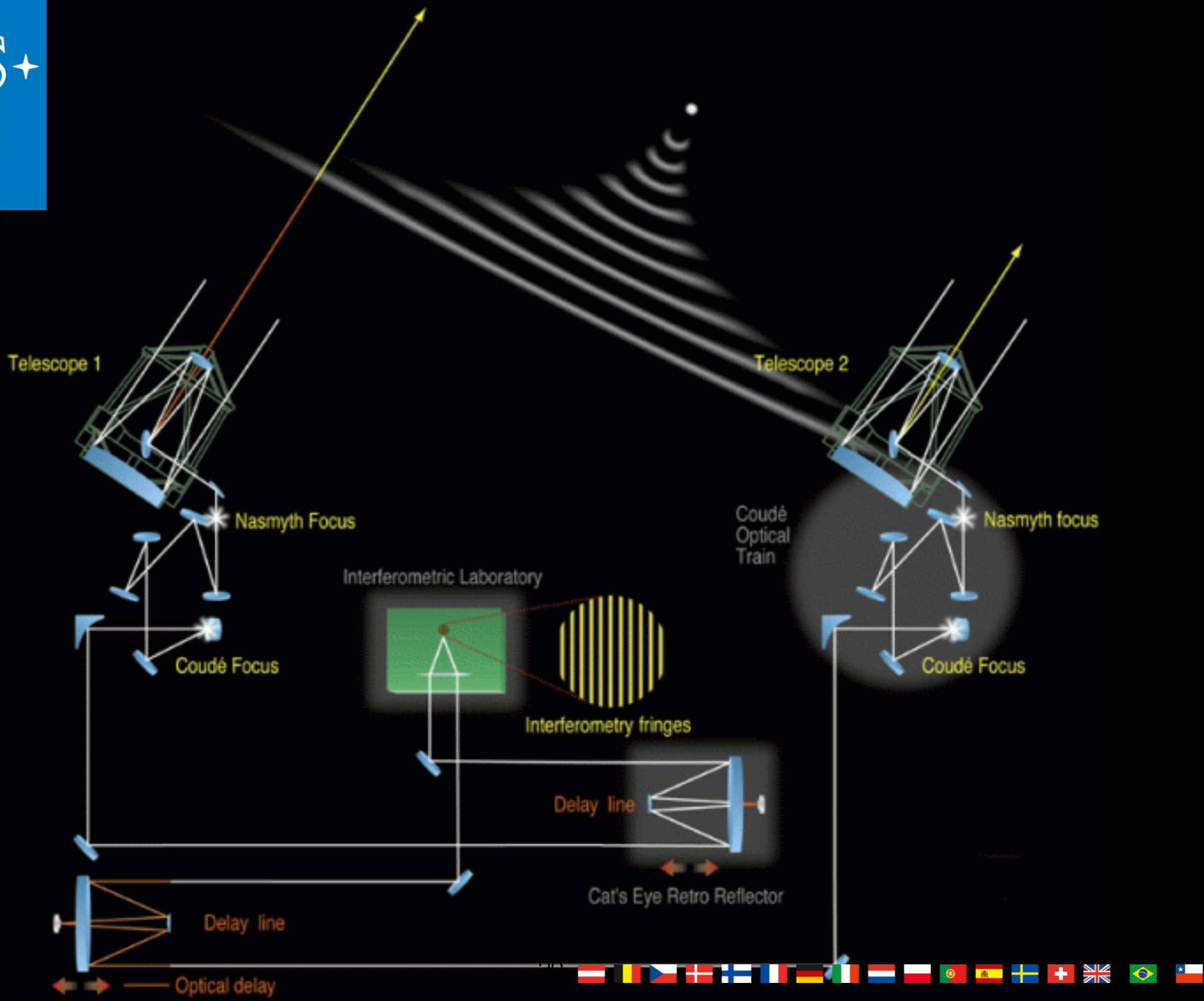
Atmospheric Dispersion

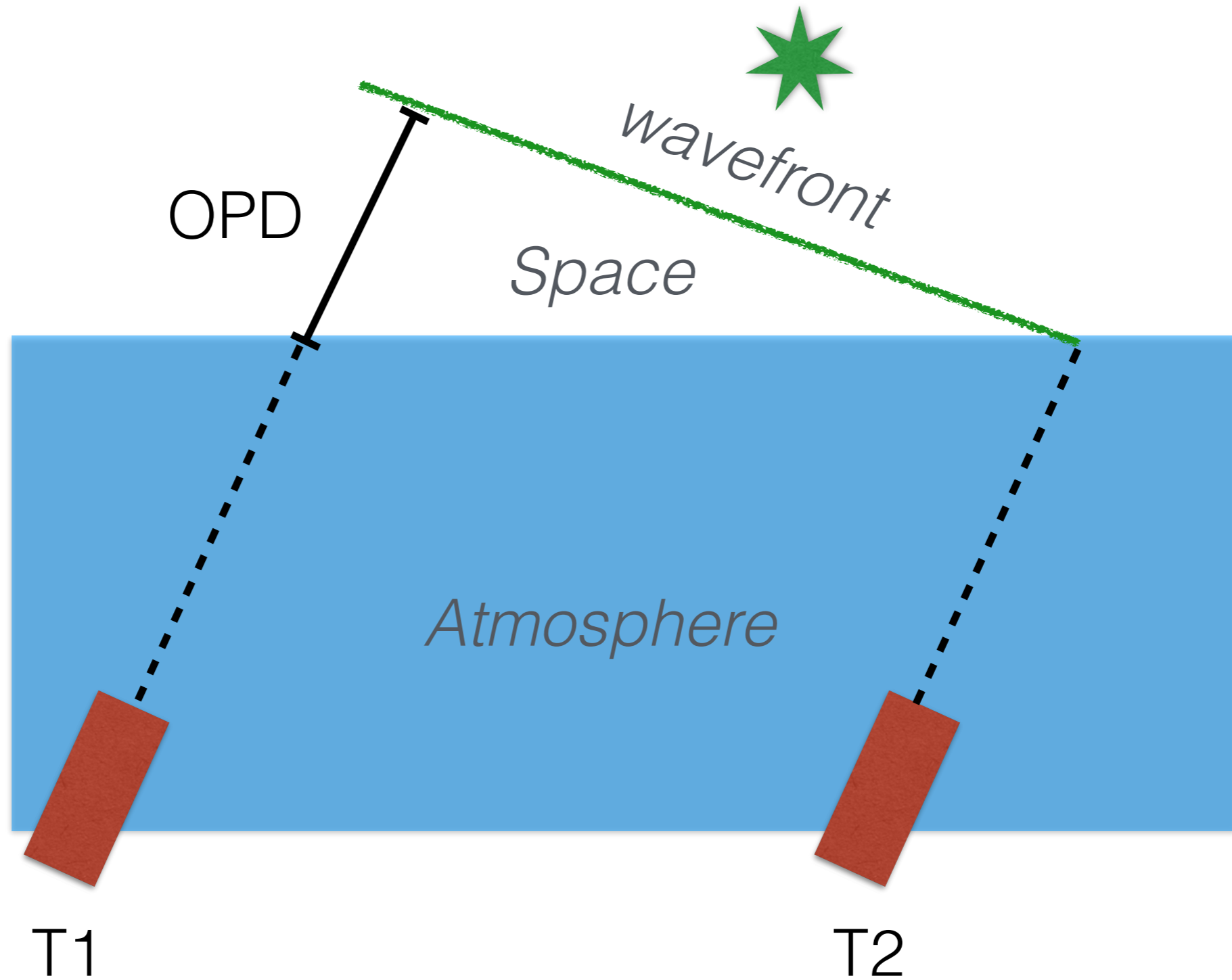
The refractive index of air is chromatic



Zheng+13





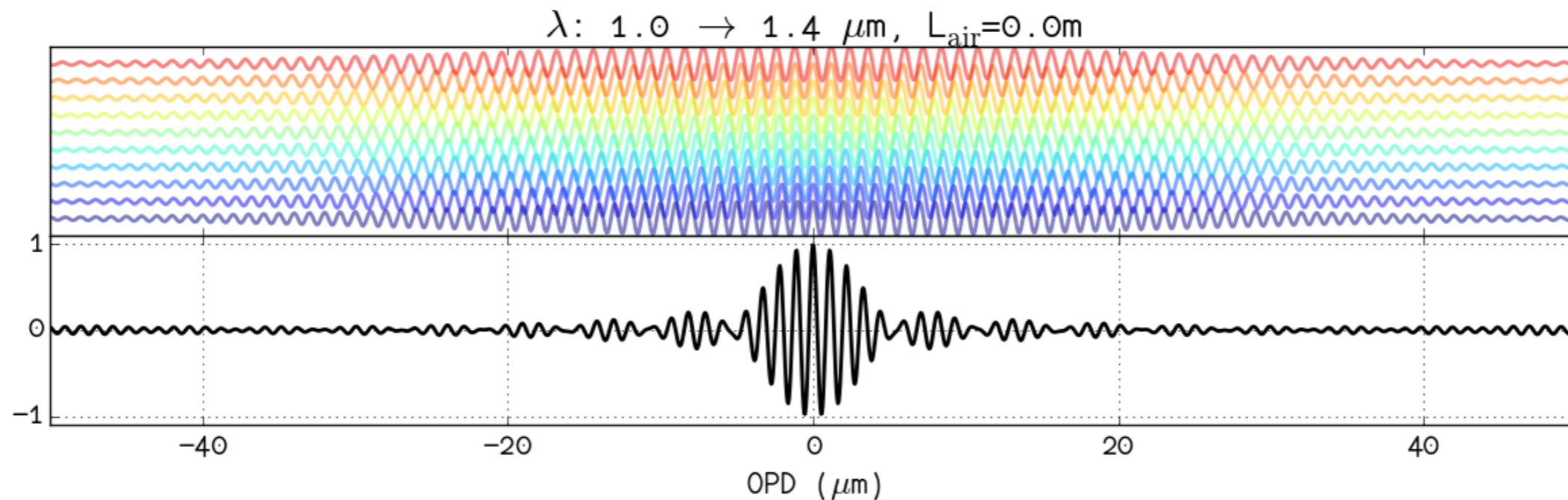


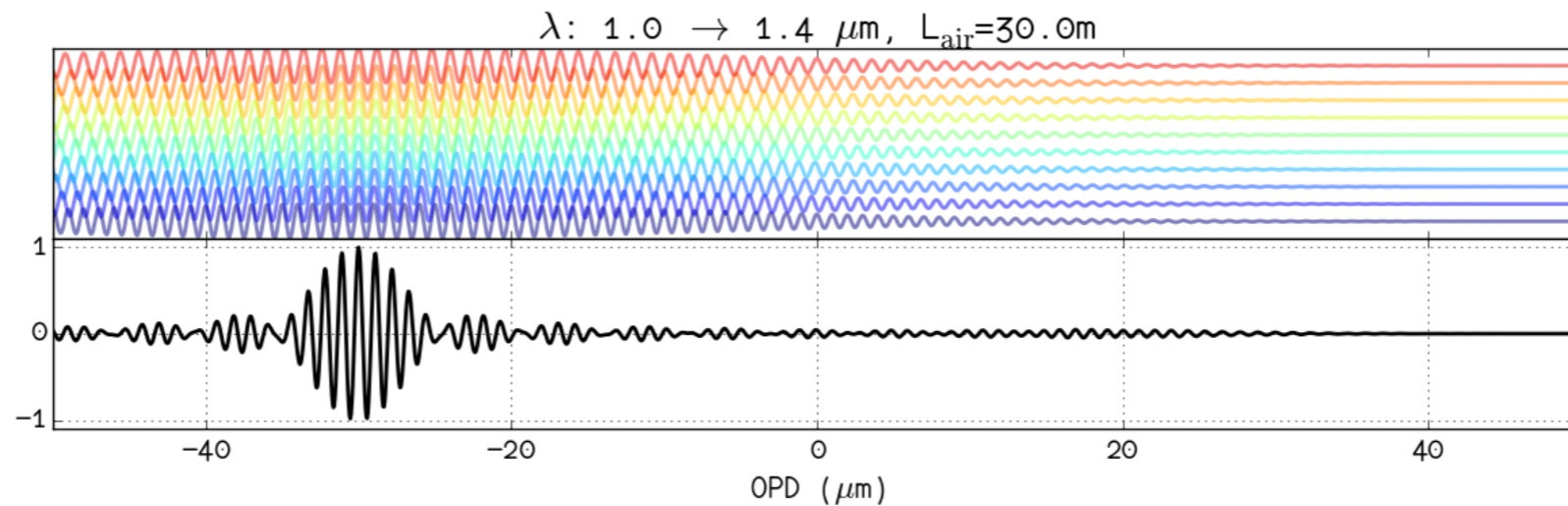
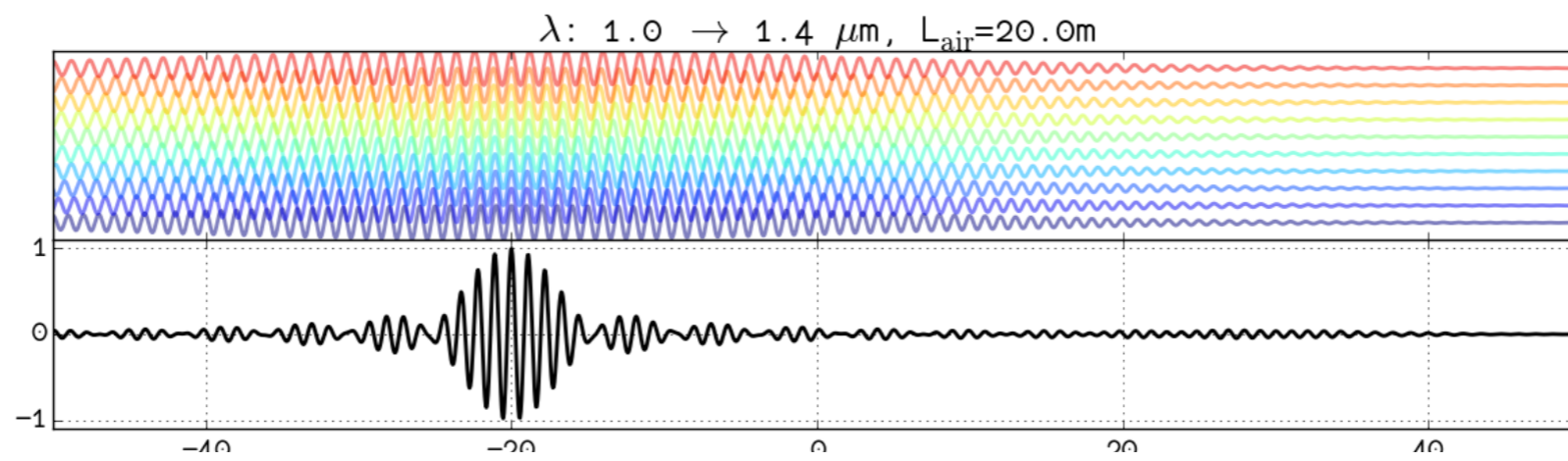
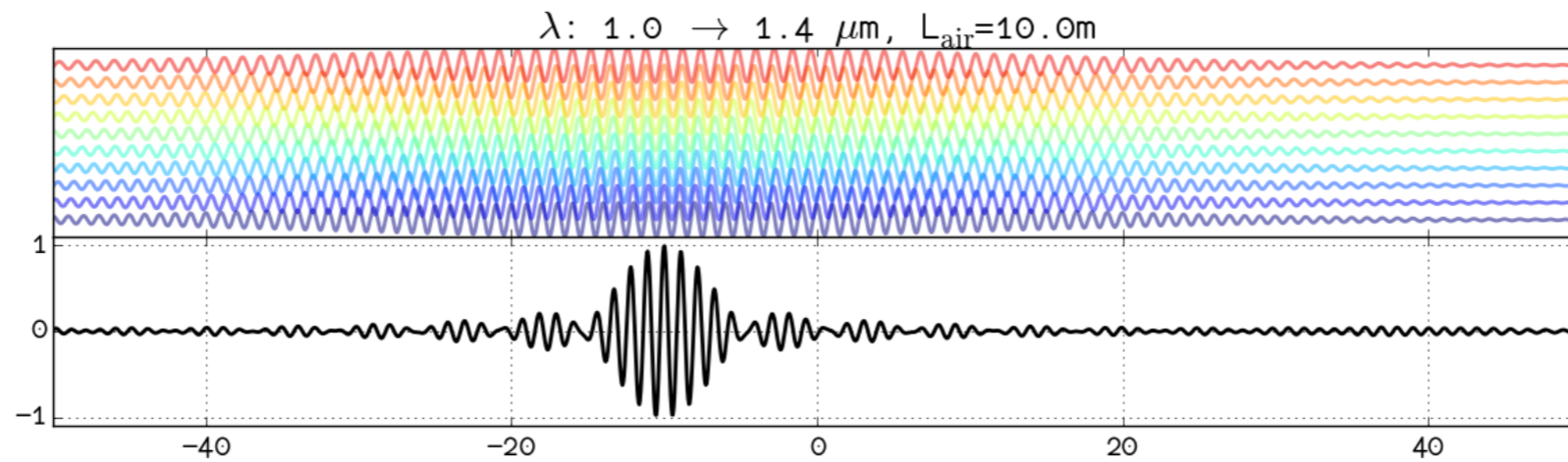
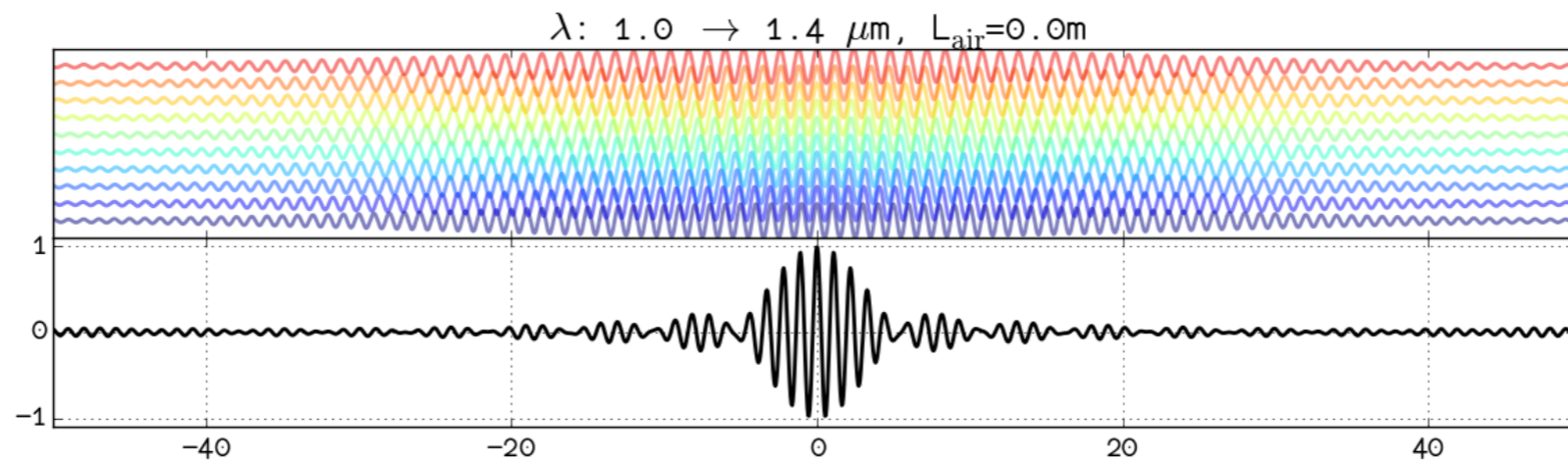
The OPD is in **vacuum**,
delay lines are in **air**



longitudinal dispersion

- The OPD in air is $l \times n(\lambda)$
- The OPD chromatic







Effect on differential phase

ideal fringes

$$F(x, \lambda) = 1 + \cos(2\pi x/\lambda + \phi_\lambda)$$

OPD modulation

$$x = L_{\text{vac.}} - [\Delta l_{i,j} = (l_j - l_i)] n_{\text{air}} = 0$$

polynomial expansion

$$F(x, \lambda) = 1 + \cos [2\pi(L_{\text{vac.}} - \Delta l_{i,j} n_0 + \delta x)/\lambda + \phi_\lambda + 2\pi\Delta l_{i,j} n_1 + 2\pi\Delta l_{i,j}(n_2\lambda + n_3\lambda^2 + \dots)]$$

actual fringes

$$F(x, \lambda) = 1 + \cos [2\pi\delta x n_0/\lambda + \phi_\lambda + 2\pi\Delta l_{i,j}(n_2\lambda + n_3\lambda^2 + \dots)]$$

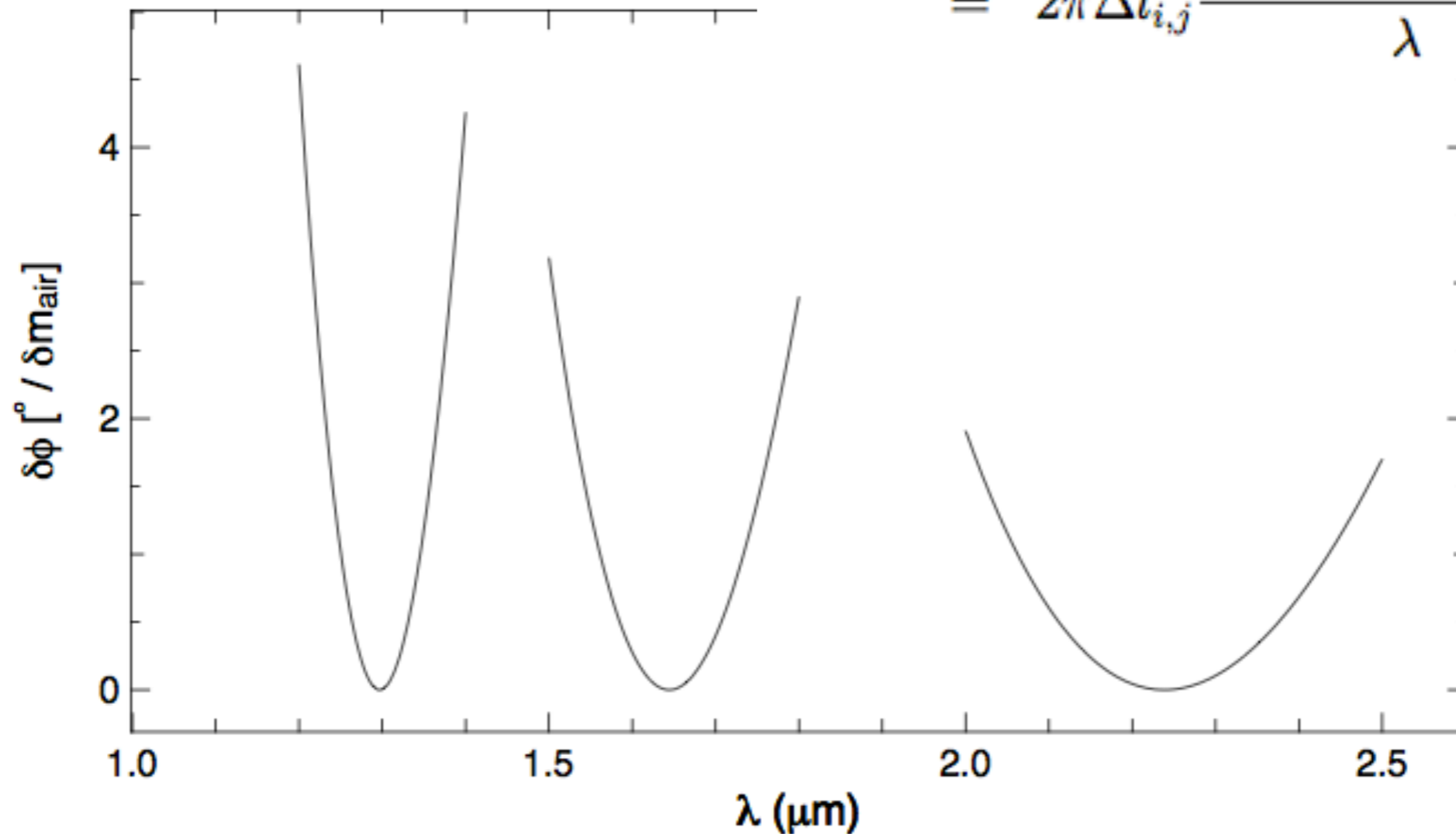
Chromatic phase

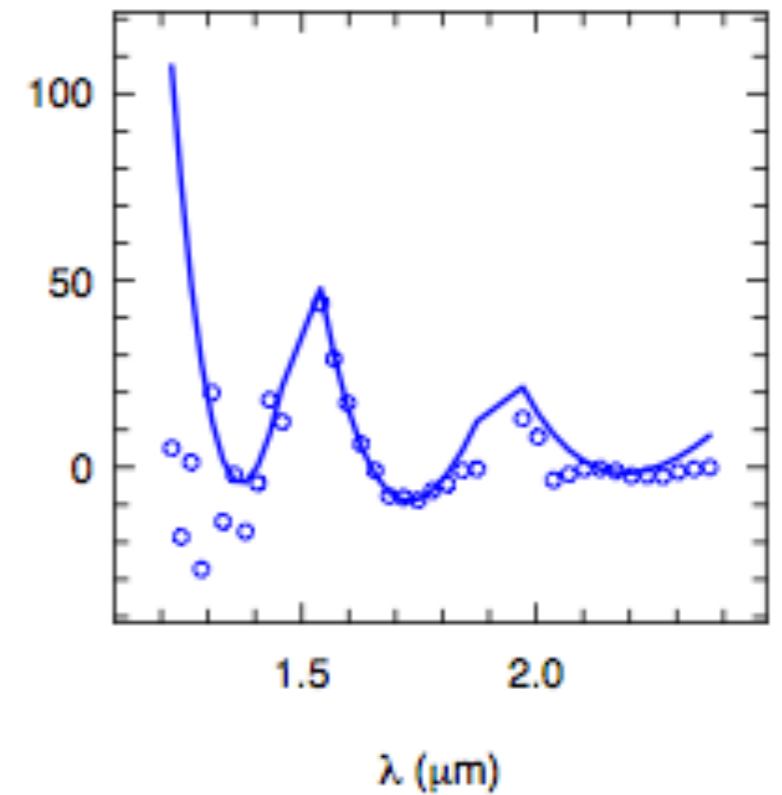
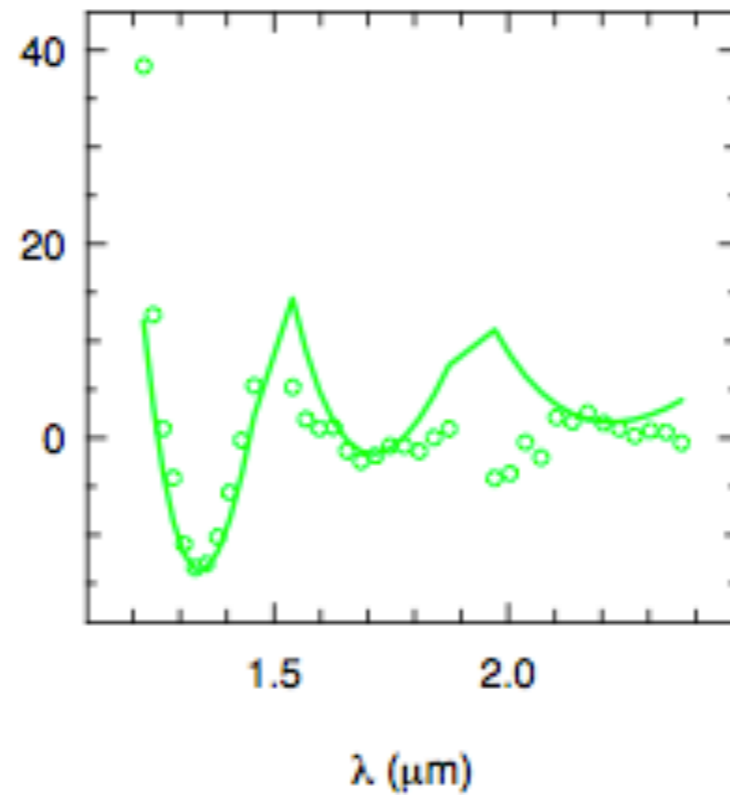
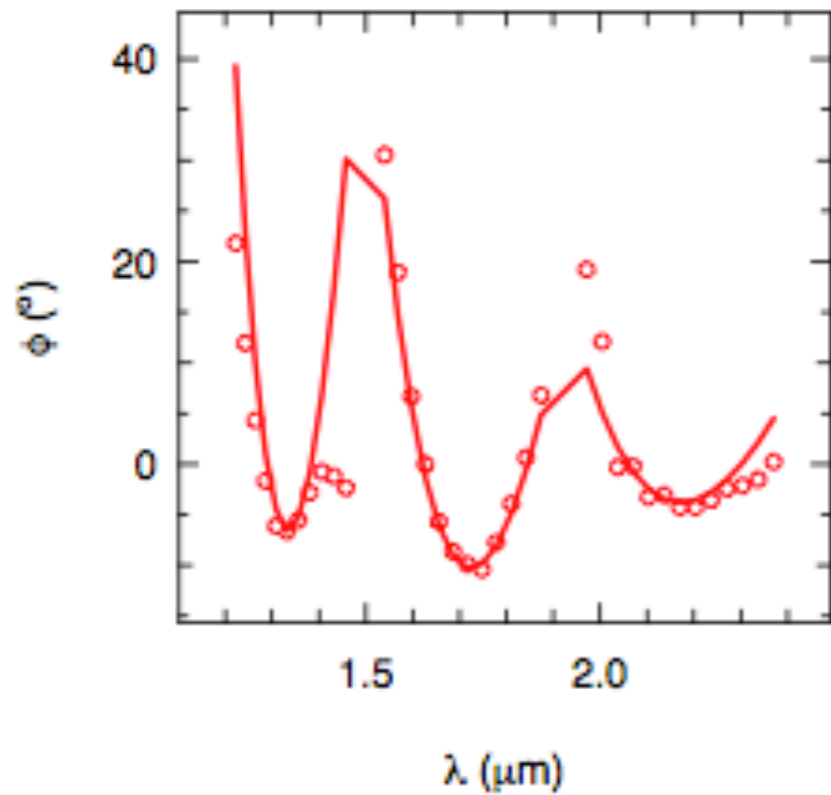
$$\begin{aligned} \phi_c(\lambda) &= 2\pi\Delta l_{i,j} (n_2\lambda + n_3\lambda^2 + \dots) \\ &= 2\pi\Delta l_{i,j} \frac{n_{\text{air}}(\lambda) - n_0 - n_1\lambda}{\lambda} \end{aligned}$$



Example: AMBER

$$\begin{aligned}\phi_c(\lambda) &= 2\pi\Delta l_{i,j} (n_2\lambda + n_3\lambda^2 + \dots) \\ &= 2\pi\Delta l_{i,j} \frac{n_{\text{air}}(\lambda) - n_0 - n_1\lambda}{\lambda}\end{aligned}$$

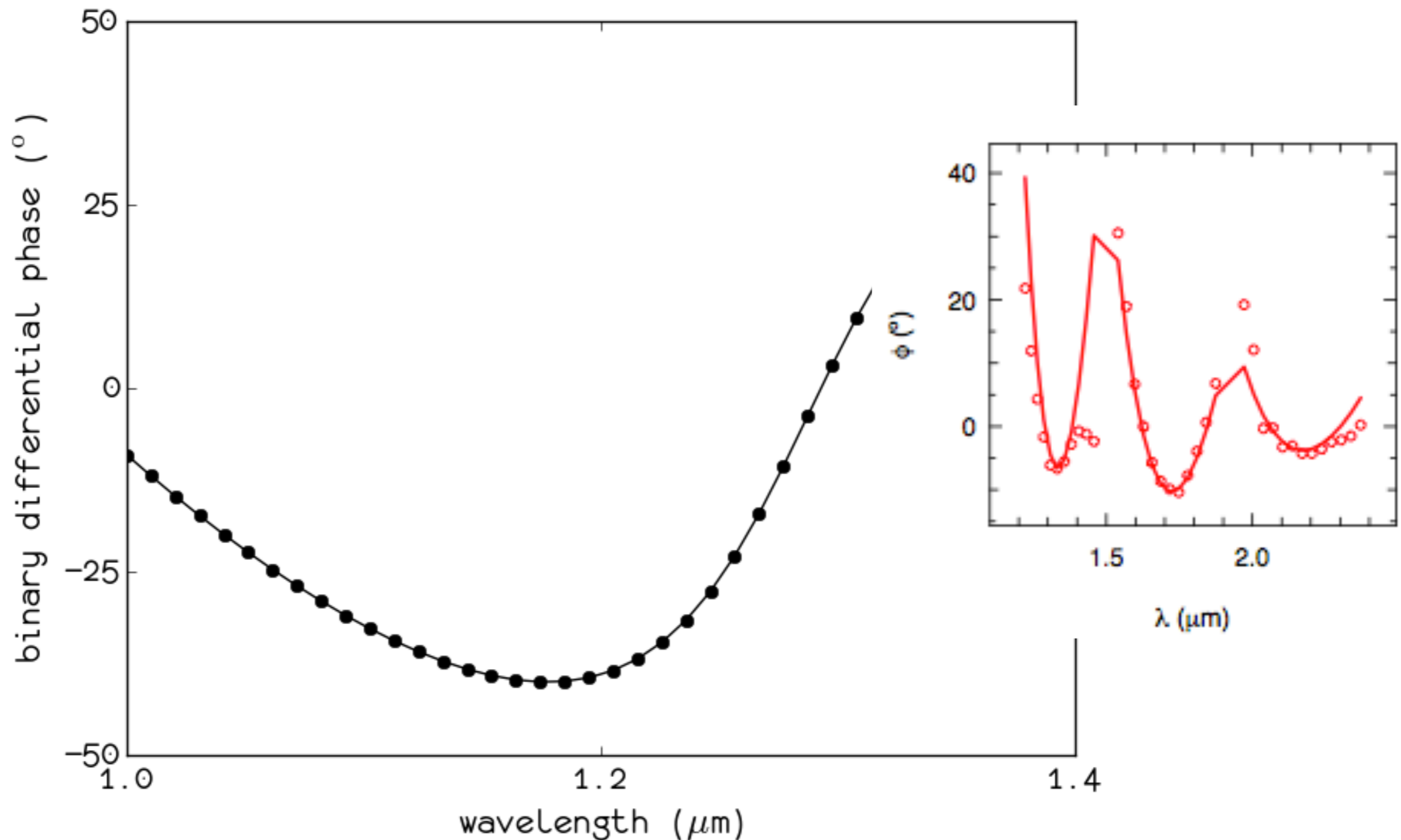




Observed differential phase
and model based on DL positions

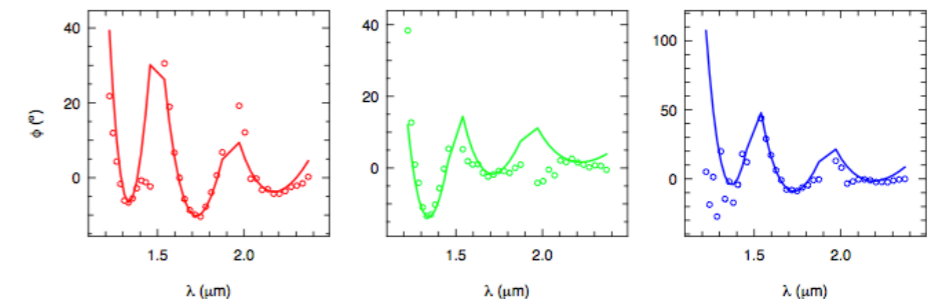
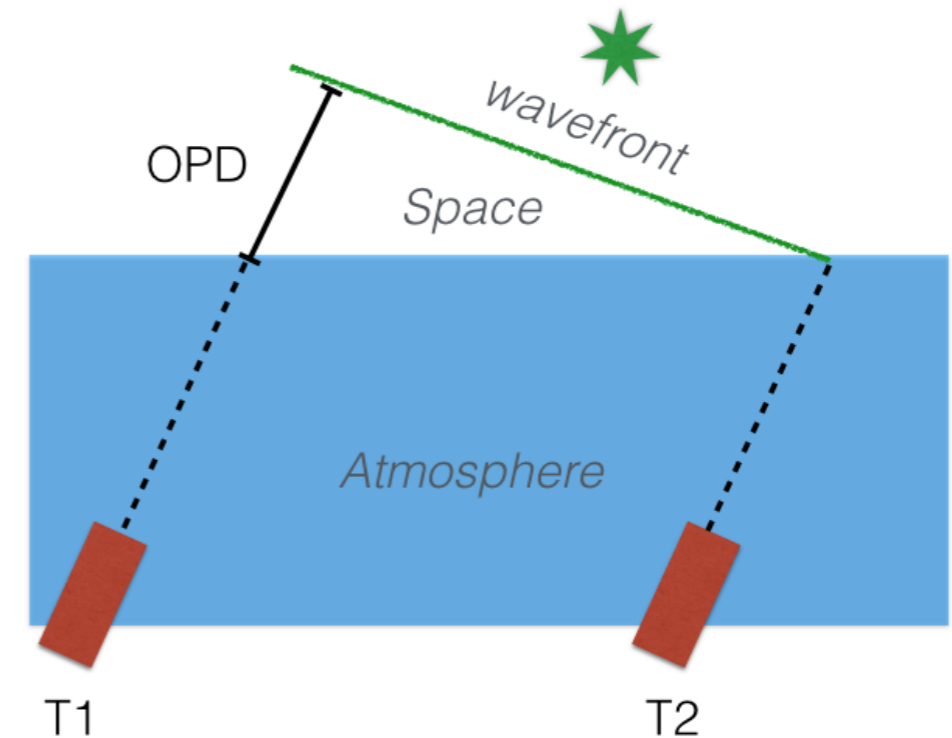


There is information in the differential phase!



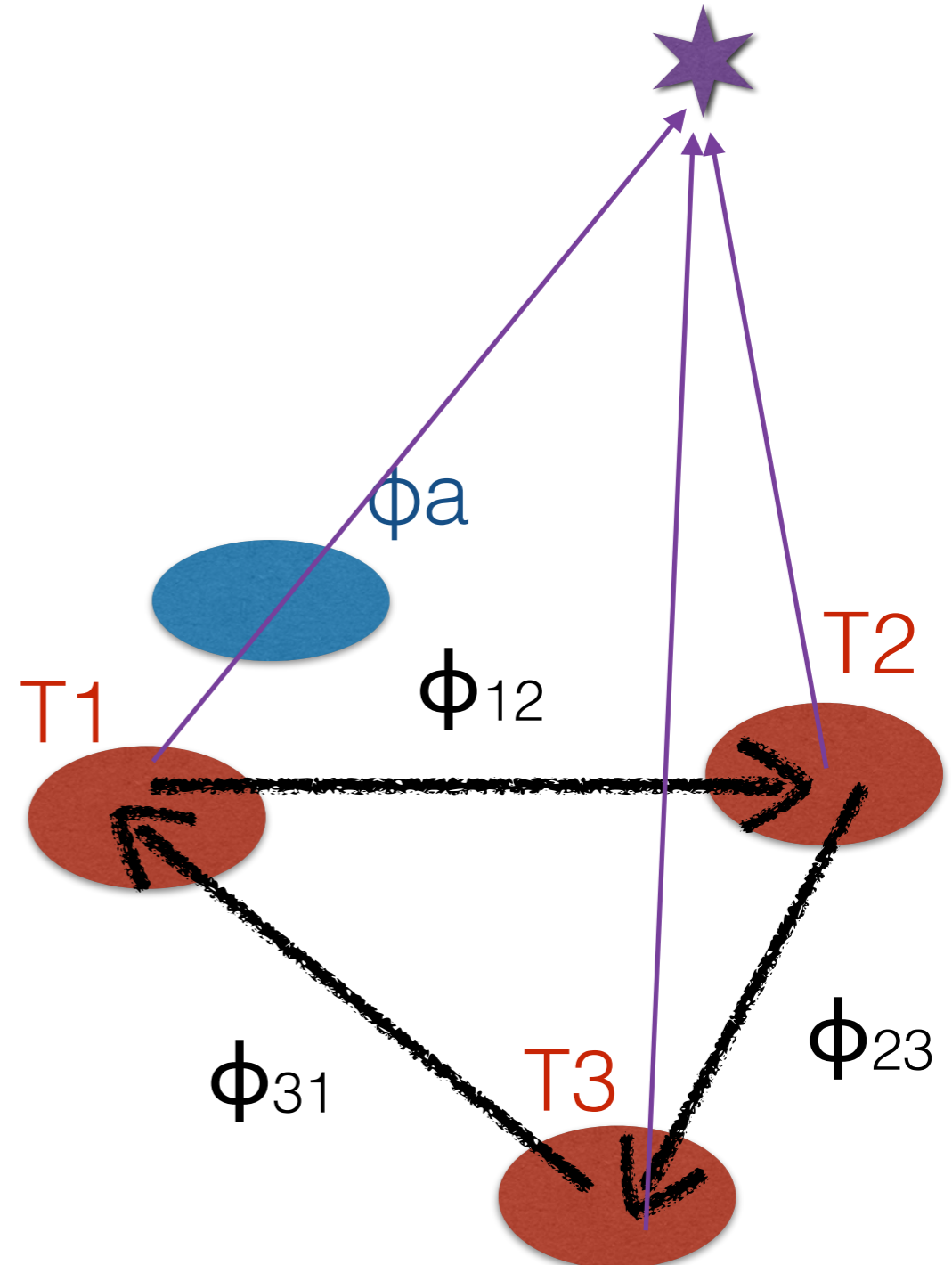
Use of differential phase?

- longitudinal air dispersion needs to be **accurately modeled** to extract the astrophysical signal
- **L_{air} is easy to estimate** (function of zenithal distance)
- n_{air} is a also function temperature, pressure, water content etc... **never perfect correction (\neq accurate)**



Closure Phase

- Measure phase sum in a **close triangle**
- $CP = (\phi_{12} + \phi_a) + (\phi_{23}) + (\phi_{31} - \phi_a) = \phi_{12} + \phi_{23} + \phi_{31}$
- CP is insensitive to longitudinal dispersion!





Differential Phase

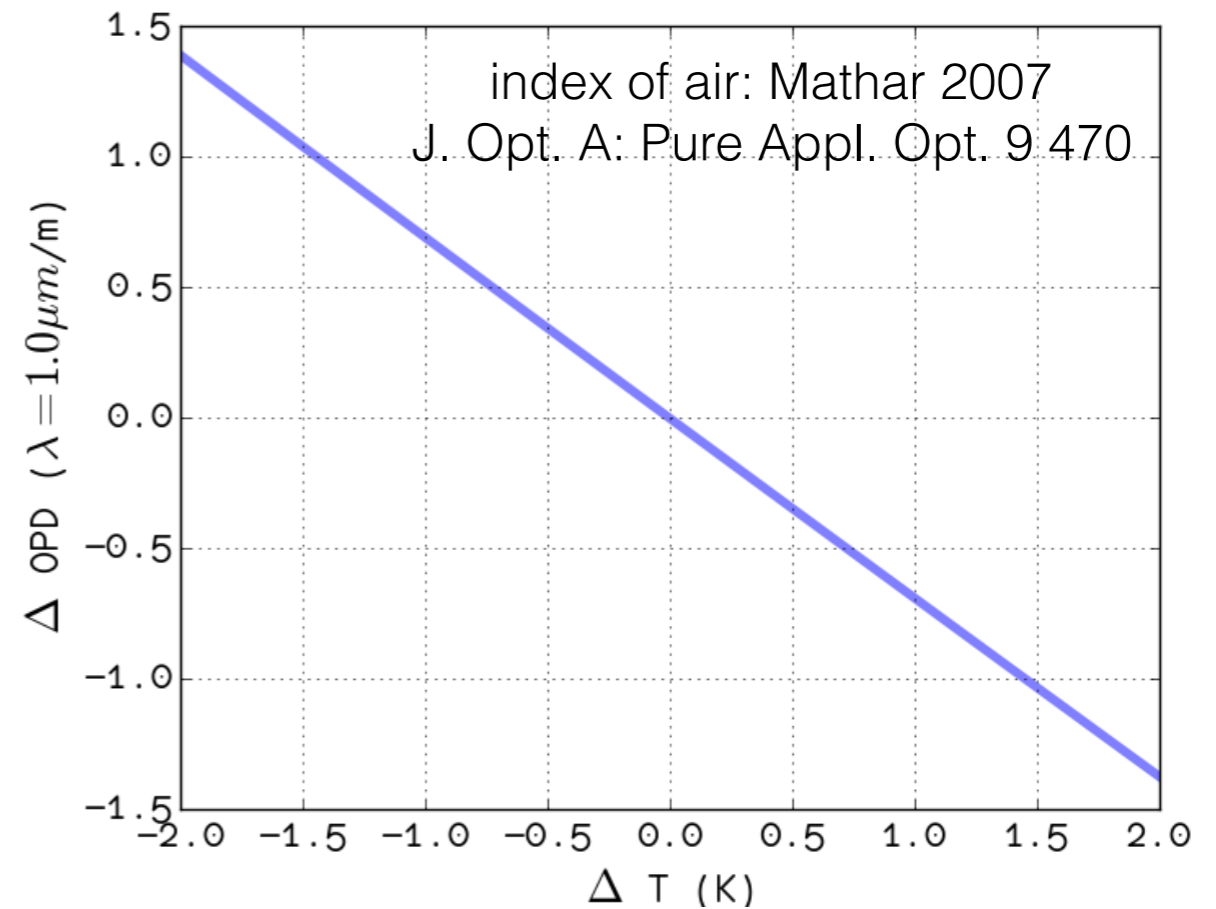
- Differential phase has a strong instrumental bias (air dispersion)
- Bias is very large (many 10°)
- We have seen 2 solutions:
 1. model the effect
 2. use a robust estimator (closure phase)
- *Alternate solution: correct with glass with refractive chromaticism inverse to air (hard to get accurate)*



Phase jitter correction

1m bubble of air with
1.5°K difference
produces a 1μm OPD
difference

same for 10m bubble
of air with 0.15°K
difference





How long can we integrate?

- Reminder: for photon shot and readout noises, the longer integration the better
- How about the turbulent piston?
- loss of contrast:

$$V_{\text{loss}}^2 = e^{-\phi^2} = e^{-\left(2\pi \frac{\sigma_{\text{OPD}}}{\lambda}\right)^2}$$



Interferometric SNR

- signal: coherent flux $\sim N_{\text{phot}} \times V_{\text{obs}}$
- Noises: read-out and photon noises

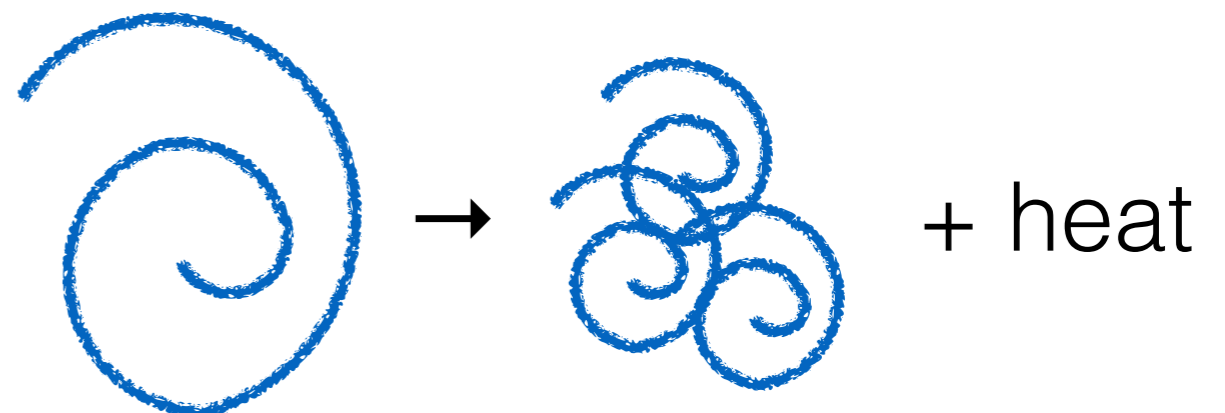
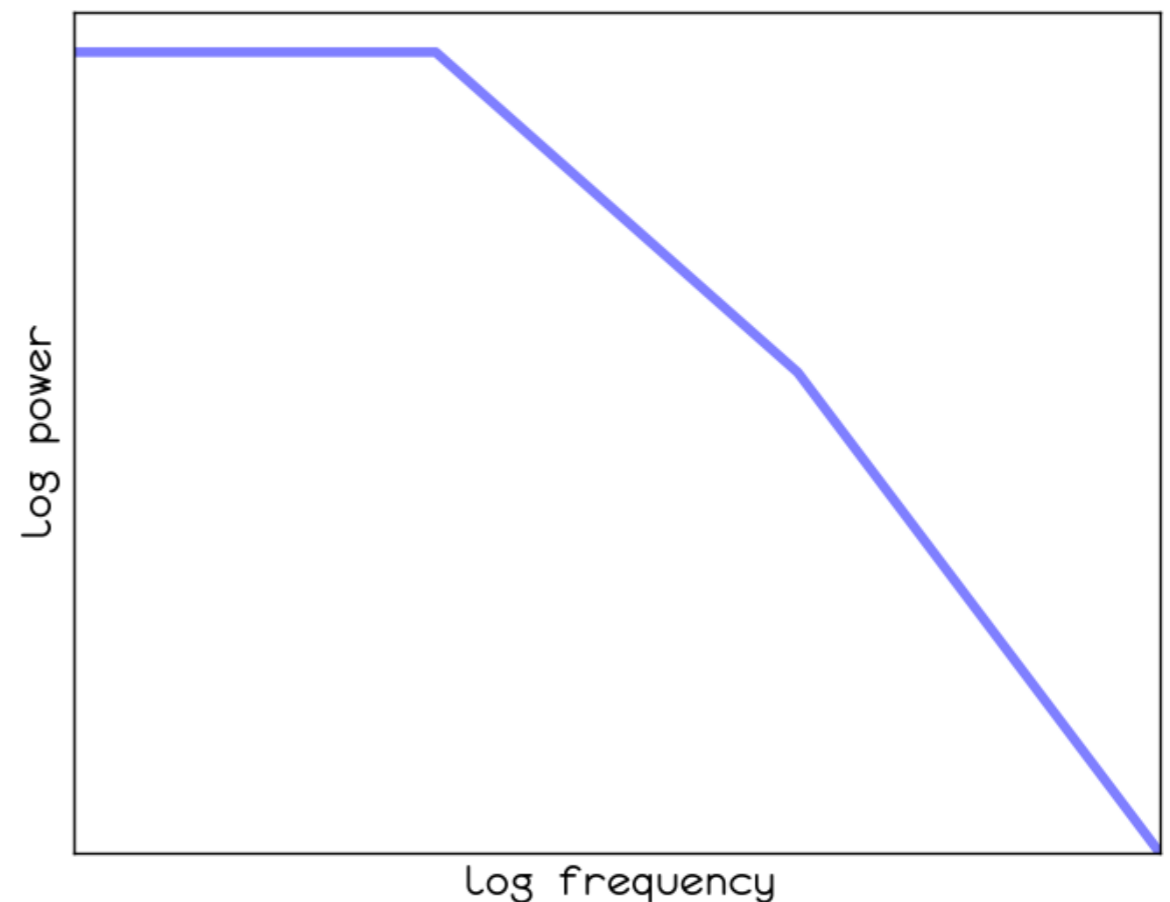
$$SNR \sim \frac{N_{\text{phot}}(t) e^{-\pi \frac{\sigma_{\text{opd}}^2(t)}{\lambda}}}{\sqrt{N_{\text{phot}}(t) + ron^2}}$$



Turbulent PSD

Typical of cascading energy phenomena

- energy injected a low frequency (wind, gravity waves)
- breaks down in smaller and smaller scales, losing each time more energy





variance?

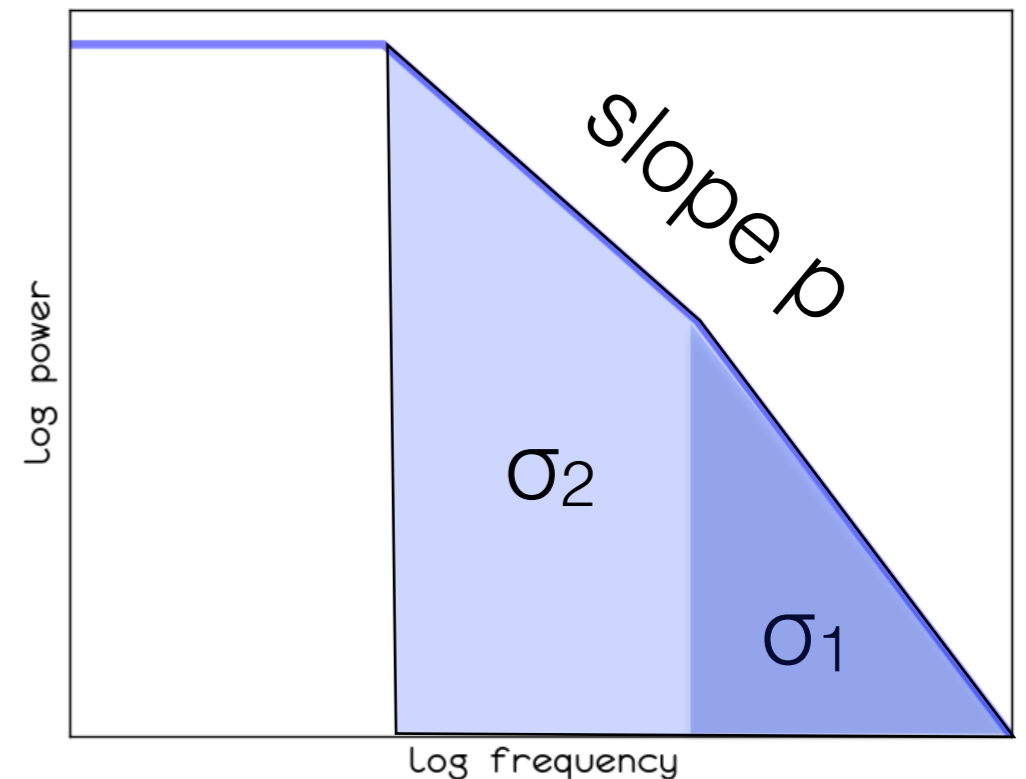
Parseval identity:

$$\int_{-\infty}^{\infty} |f(t)|^2 dt = \int_{-\infty}^{\infty} |TF(f)(\sigma)|^2 d\sigma$$

Variance \sim integral of PSD

$$\int_0^T |f(t)|^2 dt \sim \int_{1/T}^{\infty} |TF(f)(\sigma)|^2 d\sigma$$

variance grows as $T^{(-p-1)}$

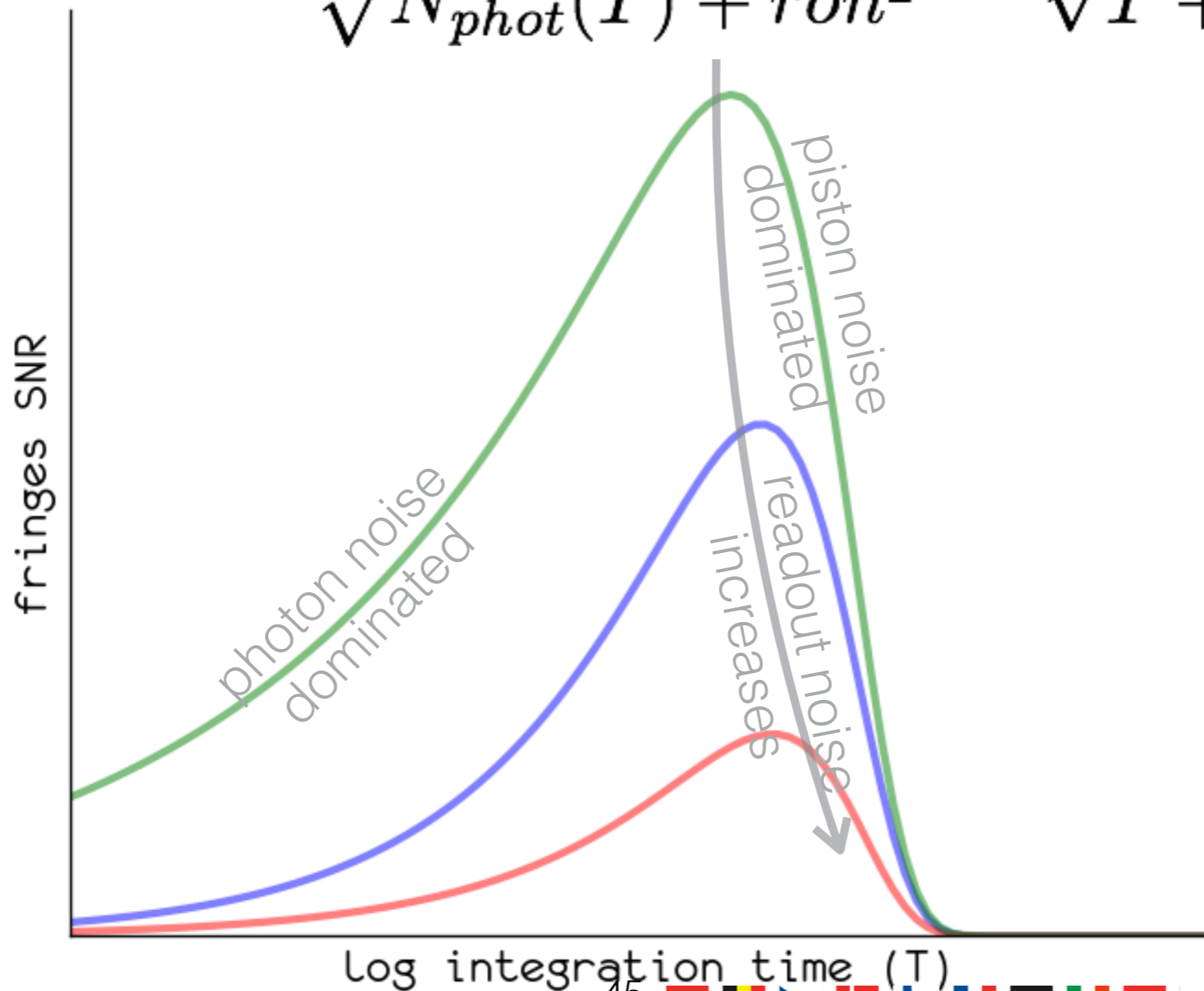


Kolmogorov $p \sim -11/3$ $\sigma_{OPD}^2(T) \sim T^{8/3}$



under turbulent atmosphere

$$SNR \sim \frac{N_{phot}(T) e^{-\pi \frac{\sigma_{opd}^2(T)}{\lambda}}}{\sqrt{N_{phot}(T) + ron^2}} \sim \frac{T e^{-T^{8/3}}}{\sqrt{T + ron^2}}$$

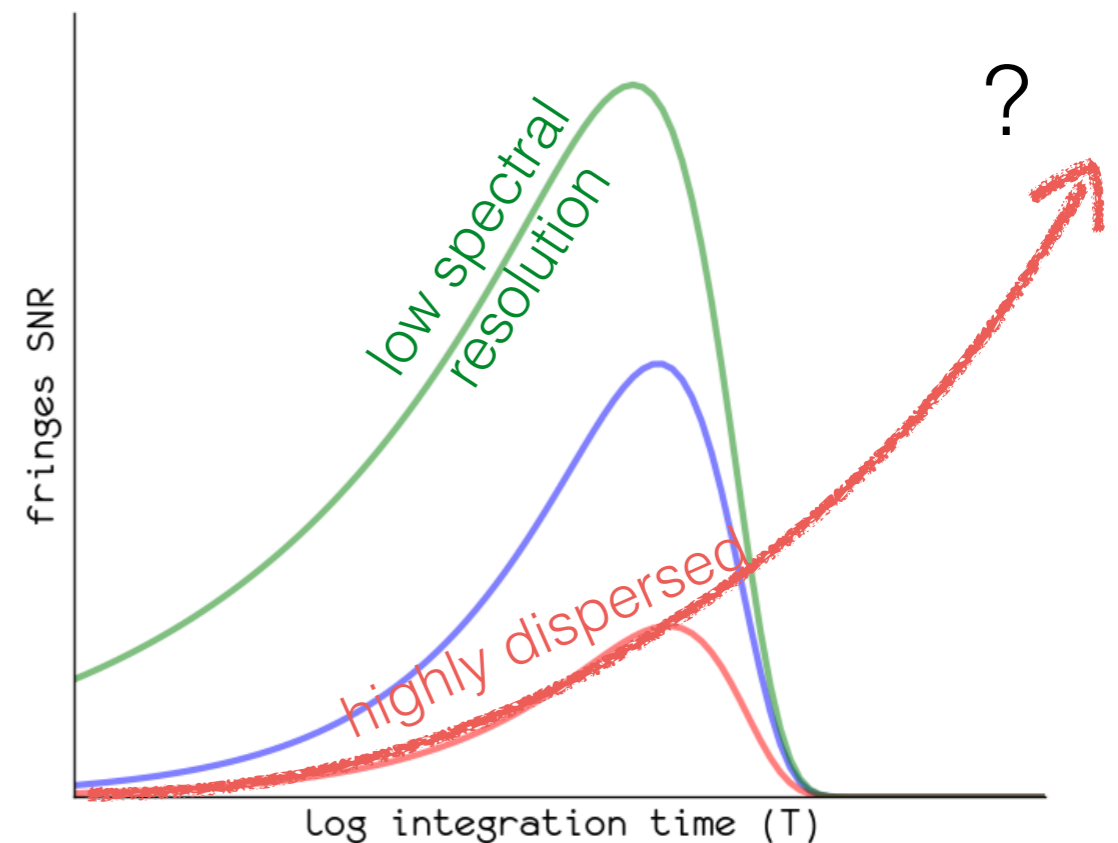




Fringe Tracking goal

Case of spectrally dispersed interferometer:

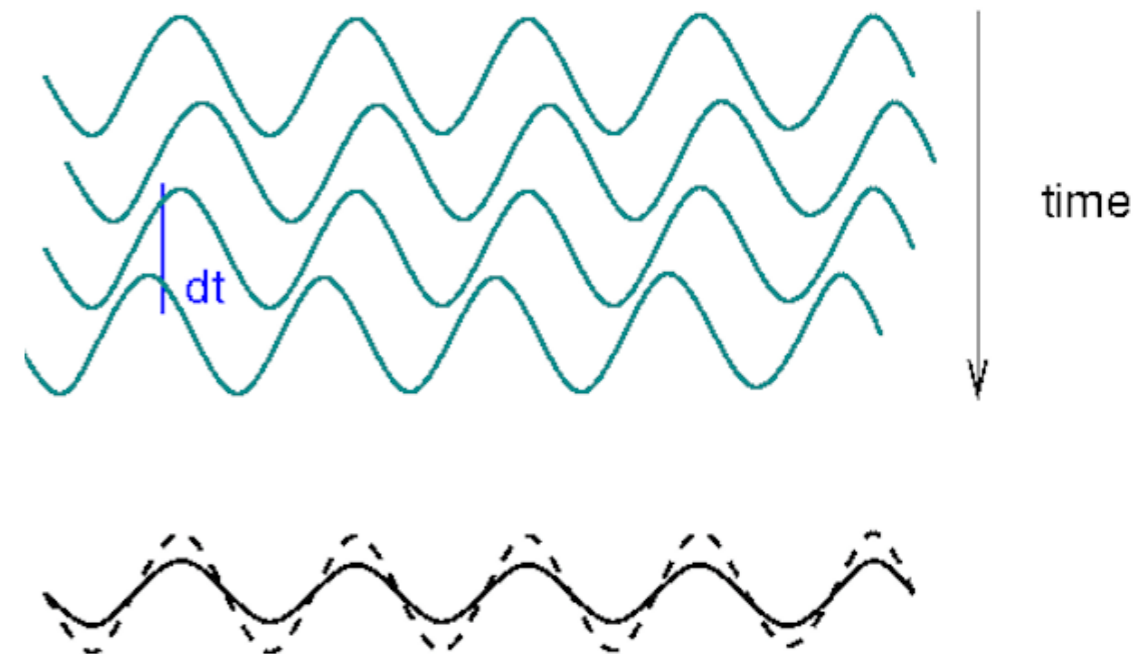
- Lack of sensitivity is a lack of photons, requiring long DIT
- everything being equal, low spectral resolution would have better SNR.





GRAVITY

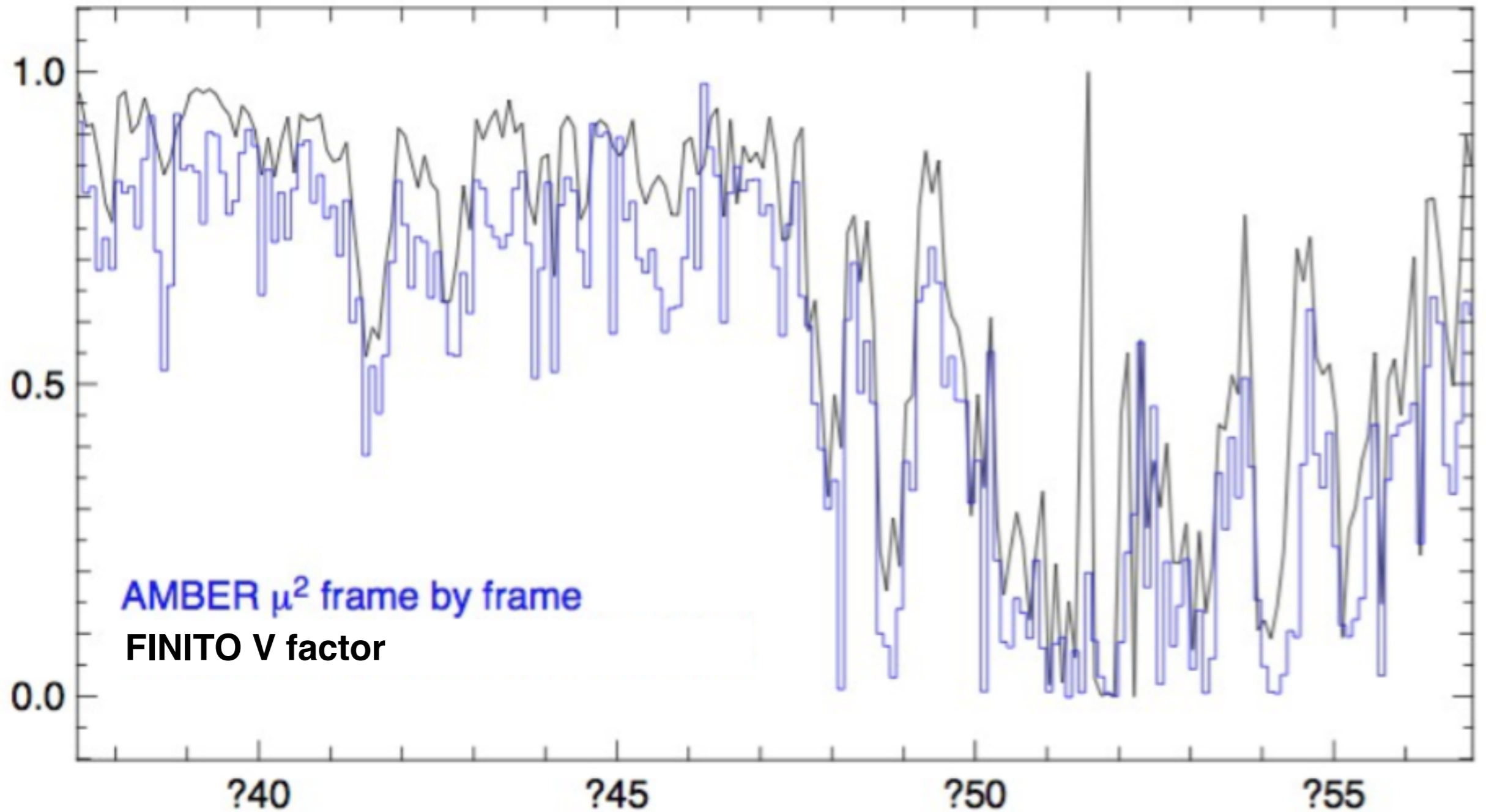
- FT measures fringes every 0.001s
- SC integrates for ~ 10 s
- FT has a tremendous amount of phase information during the SC integration
- post processing can assess the visibility loss due to FT residuals: **Gravity's V-Factor**





FINITO+AMBER

Fringe contrast per frame

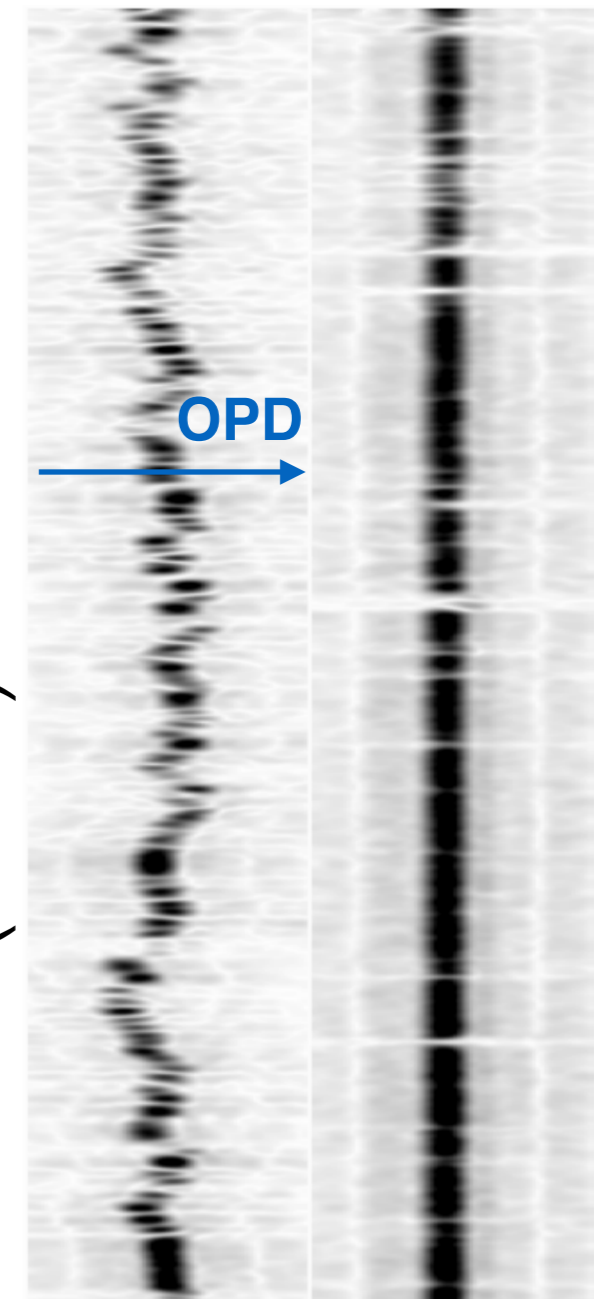




MIDI+FSU

- MIDI observed at $10\mu\text{m}$, PRIMA FSU tracked at $2.2\mu\text{m}$
- additional **post-processing correction** $2.2 \rightarrow 10\mu\text{m}$ assumes to estimate water vapor (Koresko+ 2006)
- **Gain in sensitivity of MIDI 2.5mag:**
 - Observing mode unchanged
 - FT telemetry data recorded
 - Post processing
- Paves the way for “Gravity for MATISSE”

without FT with FT





Recipe for Accuracy

Post processing and data modeling:

