Interferometric Constraints on Fundamental Stellar Parameters



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photo: S. Guisard

Monday 20 February 12











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- 2. Transportation of the energy from the center to the surface (radiation/convection)
- 3. Loss of energy from the surface by radiation into space
- The fundamental parameters: M, L, R, T_{eff}
- Also important: rotation, mass loss, metallicity, binarity,... but « second order »

- M is the source of the star's gravity, and the quantity of «fuel» it can burn
- L is the «power of the star», i.e. its energy production rate.
- R and T_{eff} are linked and essentially characterize the energy transportation mechanism inside the star
- Optical interferometry can constrain the radius R and the mass M

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 $L=4 \pi R^2 \sigma T_{eff}^4$

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L= 4 π R² σ T_{eff}⁴

• in Log, constant radius curves are:

 $Log L = 4 Log T_{eff} + 2 Log R + K$







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- Interferometric observations of VLMS: Lane et al. (2001, ApJ, 551, L81)
 Ségransan et al. (2003, A&A, 397, L5)
 Berger et al. (2006, ApJ, 644, 475)
 Boyajian et al. (2009, ApJ, 683, 424)
 Demory et al. (2009, ApJ, 505, 205)









- Ségransan et al. (2003): An accurate empirical mass-radius relation is an essential constraint on stellar interior structure, evolutionnary models and atmospheric physics. The interior structure is largely determined by the equation of state, whose derivation for very low mass stars, brown dwarfs, and planets involves the complex physics of strongly correlated and partially degenerated quantum plasma (Chabrier & Baraffe 2000).
- Demory et al. (2009): Radii of single inactive M dwarfs measured by interferometry are in excellent agreement with models from Baraffe et al. (1998). Thus, discrepancies pointed out by Torres & Ribas (2002) and Ribas (2003) only concern fast rotating stars, confirming the fact that rotation strongly affects the internal structure of those objects. Models are in good agreement with the observations, confirming a correct understanding of the underlying physics of low and very low mass stars.

- Remarkable *complementarity* of the radius and oscillation period:
 - for a pendulum: $P = 2 \pi \sqrt{(l/g)}$
 - for an oscillating star:

 $P \sim \sqrt{(R/(GM/R^2))} \sim \sqrt{(R^3/M)} \sim \sqrt{(density)^{-1}}$

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- with R and P we therefore constrain M
- Good match in measurement precision (nearby stars)

• More precisely, a prime asteroseismic observable is the large frequency spacing Δv

$$\overline{\Delta \nu_{\rm osc}} \sim 134.9 \, \sqrt{\frac{m/M_\odot}{(R_\star/R_\odot)^3}} \, (\mu {\rm Hz})$$

(Kjeldsen & Bedding 1995, A&A, 293, 87)

• R and Δv together give access to the mass, that is difficult to measure for single stars





Evolutionary modeling

• CESAM code (from P. Morel 1997, A&AS, 124, 597)

Constraints:

- L, T_{eff} , log g, [Fe/H] from spectroscopy and photometry
- Mass from the binary orbit: $M = 1.50 \pm 0.04 M_{\odot}$

(Girard et al. 2000, AJ, 119, 2428)

• Photospheric radius: $R = 2.03 \pm 0.02 R_{\odot} (1\%)$

Angular diameter θ_{LD} = 5.45 ± 0.05 mas

Limb darkening from Kurucz models

Parallax 285.93 ± 0.88 mas



8.0 C.0 model c model c	model a			
	Model m/M_{\odot} Y_i Y_s $\left(\frac{Z}{X}\right)_i$ $\left(\frac{Z}{X}\right)_s$ diffusion X_c	a 1.42 0.3012 0.2209 0.03140 0.02157 yes 0.00051	<i>b</i> 1.42 0.2580 0.2580 0.0218 0.0218 no 0.00000	c 1.50 0.345 0.202 0.0450 0.0220 yes 0.2180
Kervella et al. 2004, A&A, 413, 251	age (Myr) T_{eff} (K) $\log g$ $[Fe/H]_i$ $[Fe/H]_s$ $\log(L/L_{\odot})$ $\frac{R/R_{\odot}}{\Delta v_0}$ (μ Hz)	2 314 6524 3.960 +0.107 -0.055 0.8409 2.0649 54.7	2 710 6547 3.967 -0.051 -0.051 0.8405 2.0495 55.4	1 300 6553 3.994 +0.264 -0.043 0.8390 2.0420 56.4

Seismic frequencies

$$\overline{\Delta \nu_{\rm osc}} \sim 134.9 \, \sqrt{\frac{m/M_\odot}{(R_\star/R_\odot)^3}} \, (\mu {\rm Hz})$$

Procyon's large freq. spacing:

 $\Delta v_0 = 53.6 \pm 0.5 \ \mu \text{Hz}$

(Martic et al. 2004, A&A, 418, 295)

 $\Delta v_0 = 55.5 \pm 0.5 \ \mu Hz$

(Eggenberger et al. 2004, A&A, 422, 247)

 $\Delta v_0 = 55.66 \pm 0.15 \ \mu Hz$

(Leccia et al. 2006, MmSAI, 77, 462)

Lower mass favored:

- seismic large freq. spacing
- cooling time of the WD > 1.7 Gyr

Model	а	b	С
m/M_{\odot}	1.42	1.42	1.50
Yi	0.3012	0.2580	0.345
Y _s	0.2209	0.2580	0.202
$\left(\frac{Z}{X}\right)_i$	0.03140	0.0218	0.0450
$\left(\frac{Z}{X}\right)_{s}$	0.02157	0.0218	0.0220
diffusion	yes	no	yes
$X_{\rm c}$	0.00051	0.00000	0.2180
age (Myr)	2314	2710	1 300
$T_{\rm eff}$ (K)	6524	6547	6553
$\log g$	3.960	3.967	3.994
[Fe/H] _i	+0.107	-0.051	+0.264
[Fe/H] _s	-0.055	-0.051	-0.043
$\log(L/L_{\odot})$	0.8409	0.8405	0.8390
R/R_{\odot}	2.0649	2.0495	2.0420
$\overline{\Delta v_0} (\mu \text{Hz})$	54.7	55.4	56.4

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New HST astrometry gives a mass of 1.43 \pm 0.03 M_{\odot} (Gatewood & Han 2006, AJ, 131, 1015)

Classical field of long-baseline interferometry

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- Coeval stars with the same initial metallicities are excellent modeling subjects (e.g. α Cen A & B)

12 Boo



Boden et al. 2005, ApJ, 627, 464

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12 Boo



 $M(A) = 1.4160 \pm 0.0049 M_{\odot} (0.34\%)$ $M(B) = 1.3740 \pm 0.0045 M_{\odot} (0.33\%)$

Boden et al. 2005, ApJ, 627, 464

VLTI / PIONIER



Le Bouquin et al. 2011, A&A, in press, arXiv:1109.1918

What next ?



Cunha et al. 2007, Astron Astrophys Rev, 14, 217

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- Increasing the VLTI angular resolution (B=200m, shorter wavelength) would open up exciting new possibilities

