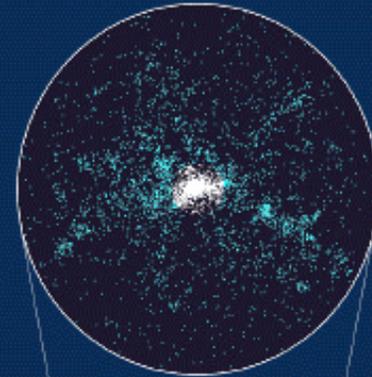


$\Lambda$ CDM

*The Agony and  
the Ecstasy:  
Galaxy Clusters in  
the Data Rich Era*



**August (Gus) Evrard**  
Arthur F. Thurnau Professor  
Departments of Physics and Astronomy  
Michigan Center for Theoretical Physics  
University of Michigan

Evrard et al 2002



Search All

Go

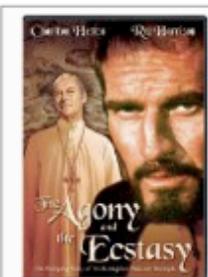
Register | Login | Help

The Internet Movie Database

Movies TV News Videos Community IMDbPro

IMDb - Now In Spanish

IMDb > The Agony and the Ecstasy (1965)



# The Agony and the Ecstasy (1965) [More at IMDbPro »](#)

## Overview

User Rating: ★★★★★☆ 6.9/10 [2,005 votes](#)

MOVIEmeter: Ⓢ Up 8% in popularity this week. See [why](#) on [IMDbPro](#).

Director: [Carol Reed](#)

Writers: [Irving Stone](#) (novel)  
[Philip Dunne](#) (screen story and screenplay)

Contact: View [company](#) contact information for The Agony and the Ecstasy on [IMDbPro](#).

Release Date: 7 October 1965 (USA) [more](#)

Genre: [Drama](#) | [History](#) [more](#)

Tagline: From the age of magnificence comes a new magnificence in motion pictures [more](#)

Plot: The biographical story of Michelangelo's troubles while painting the Sistine Chapel at the urging of Pope Julius II. [full summary](#) | [add synopsis](#)

Plot Keywords: Spoiler alert! Rollover or vote to view plot keyword: [more](#)

Awards: Nominated for 5 Oscars. Another 2 wins & 3 nominations [more](#)

NewsDesk: [Charlton Heston Dies at 84](#)  
(From IMDb News, 6 April 2008)

### Sponsored links

[Para Hombres Deportistas](#)  
Exclusivo Perfume Dior Para Hombres Modernos & Deportistas. Descúbrelo.  
[Dior.com/Homme](#)

advertisement

[Watch it at Amazon](#)

[Buy it at Amazon](#)

[Discuss in Boards](#)

[More at IMDb Pro](#)

[Add to My Movies](#)

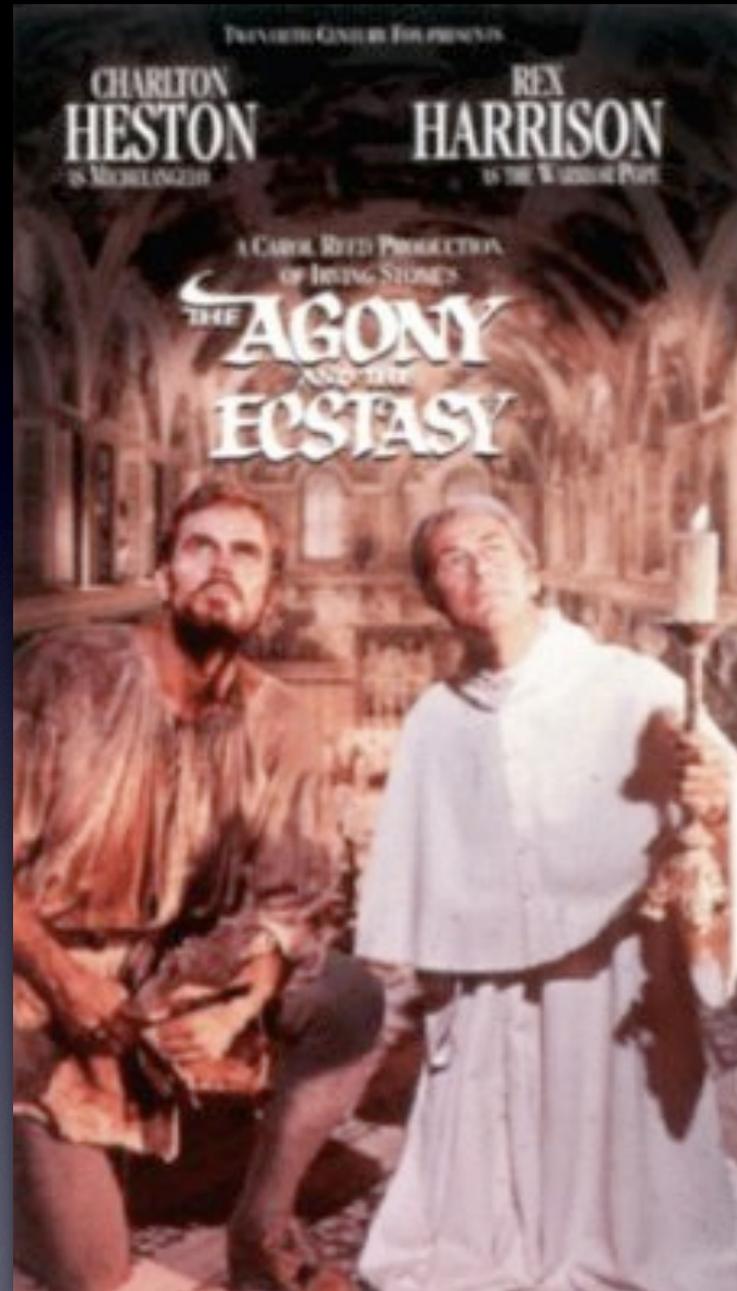
[Update Data](#)

### Quicklinks

main details

### Top Links

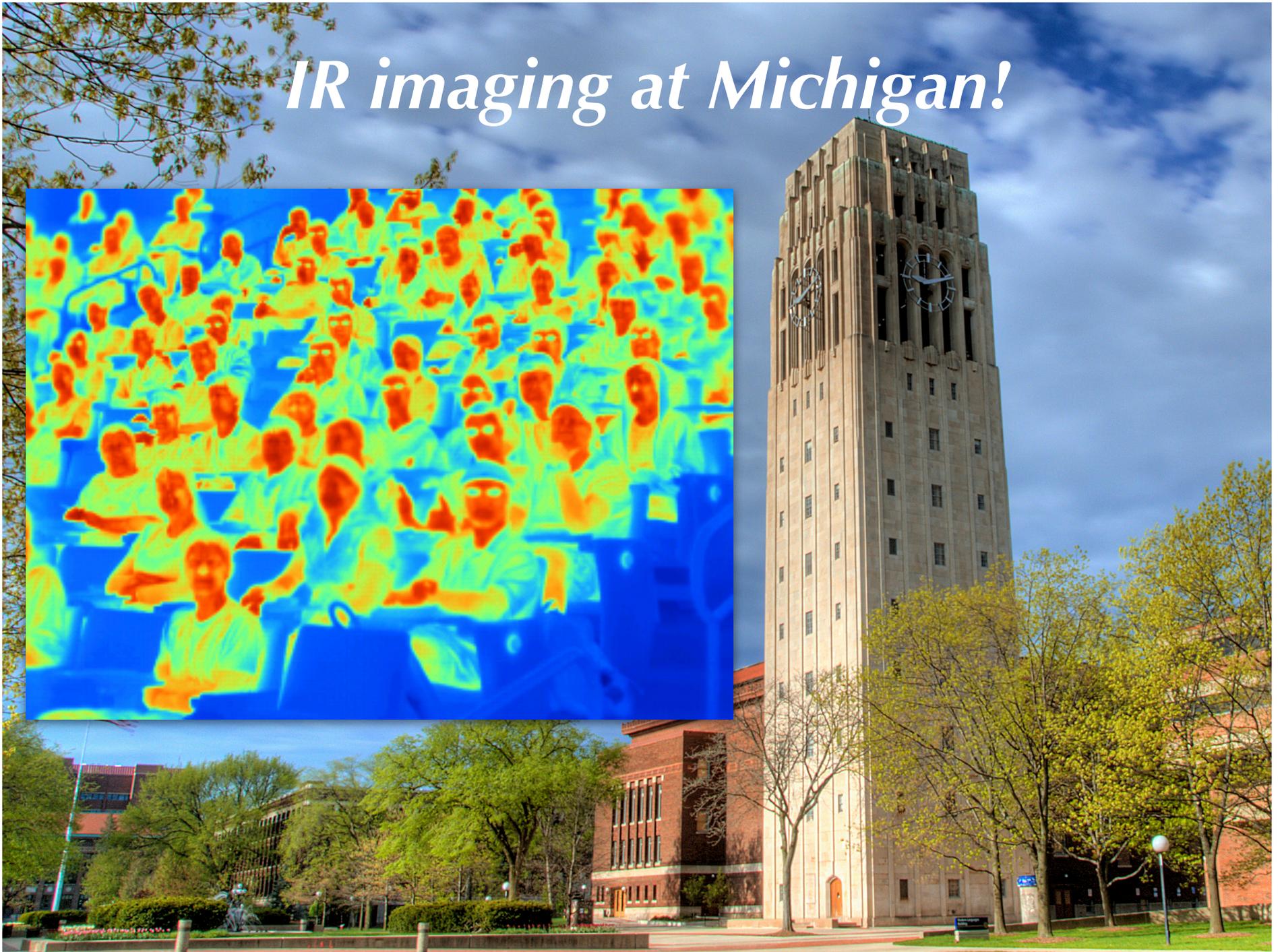
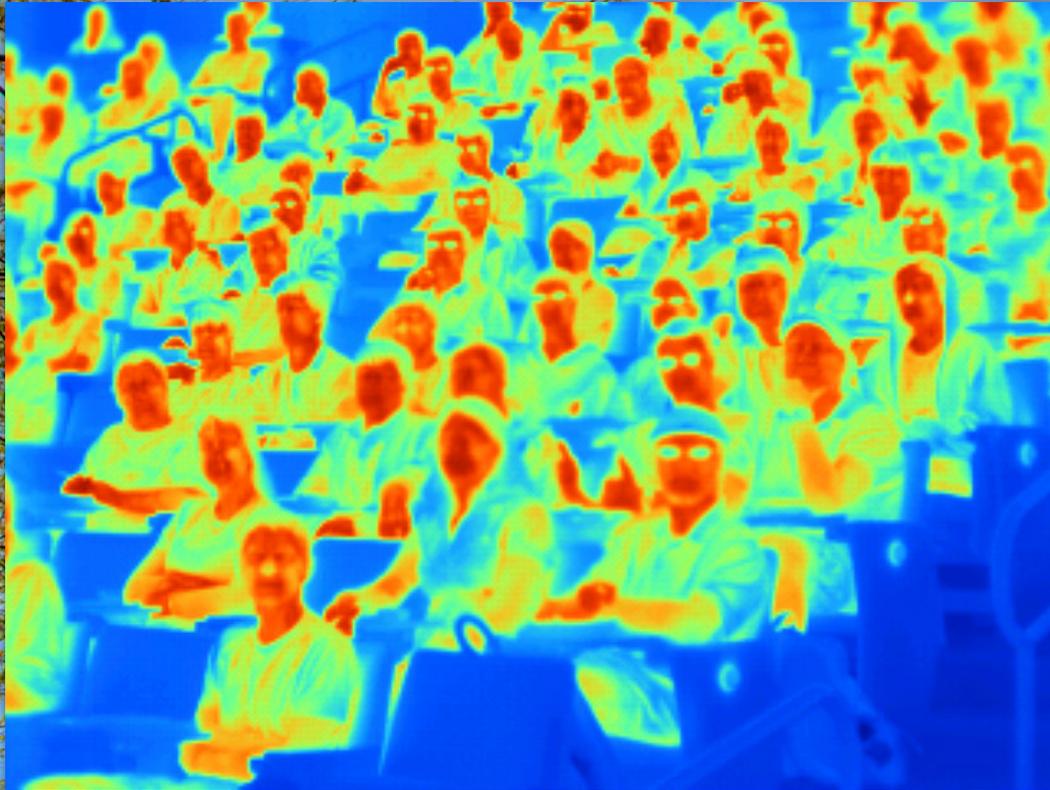
- trailers and videos
- [full cast and crew](#)
- [trivia](#)
- official sites
- [memorable quotes](#)



I will be finished  
when I am done!

Michaelangelo,  
when will you be  
finished ?!

*IR imaging at Michigan!*



## The ecstasy: new physics!

Use cluster signatures  
(and other probes) to  
convincingly  
demonstrate a  
non-trivial dark energy  
equation of state,  
 $w \neq -1$ .



## The agony: systematics!

1. 3D halo mass is not observable  
need accurate **form** of observable-mass  
scaling relations,  $p(M_{\text{obs}}, z_{\text{obs}} | M, z)$
2. We observe on the sky, not in real-space  
must model cluster-halo selection function including  
projection effects along Gpc sight-lines
3. Require theory to map counts to cosmological parameters  
N-body simulation calibrations need to include baryons
4. Cluster redshifts ( $\geq 10k$  of them!) are needed  
require mapping from photometric to spectroscopic  $z$ 's



## Some definitions

### 1. Halo -

a self-bound, quasi-equilibrium cosmic structure comprised of multiple interacting fluids (dark matter, multi-phase baryons, and radiation)

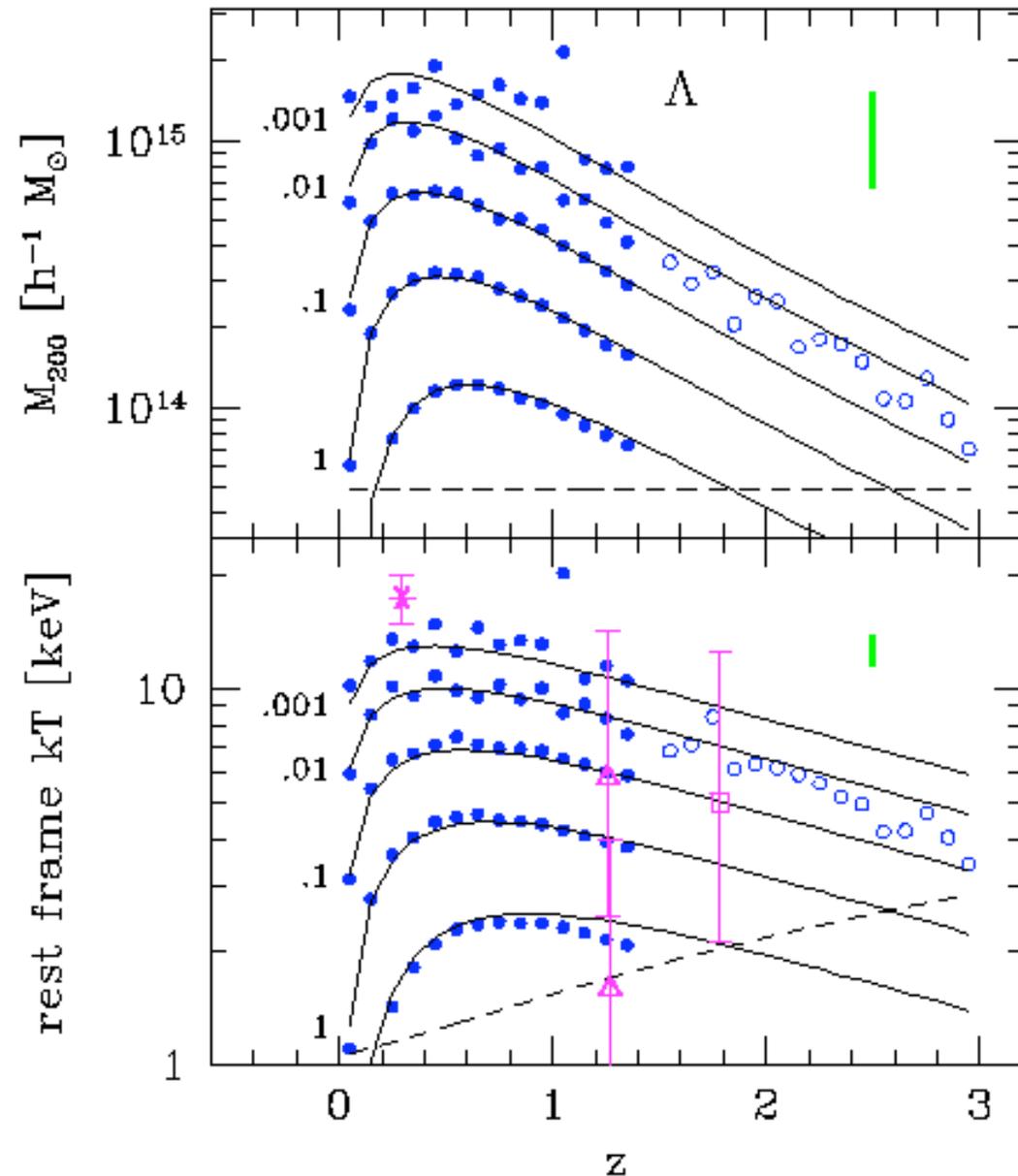
### 2. Cluster -

an observable manifestation of a massive halo containing multiple, bright galaxies

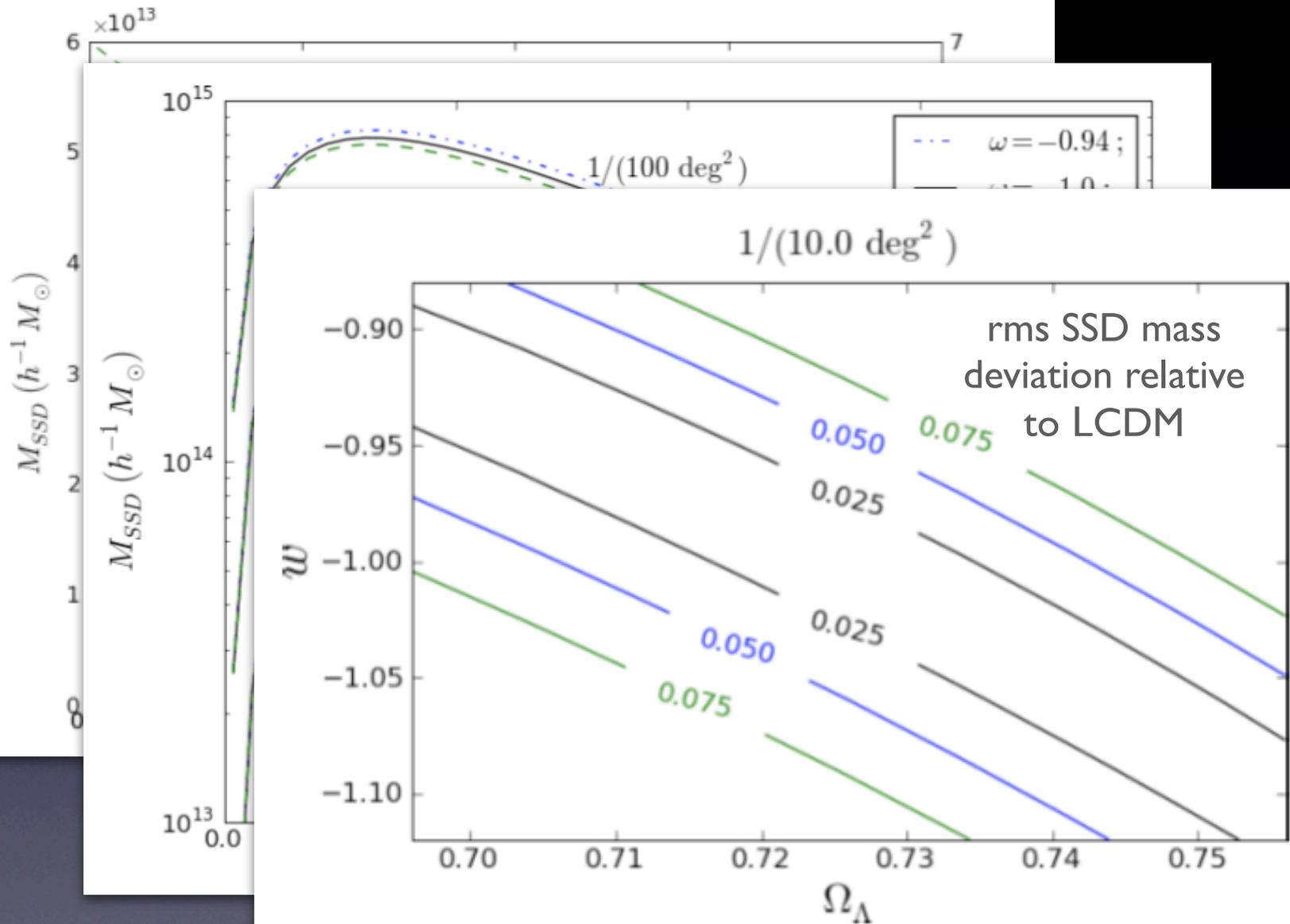
sky surface density  
characteristic mass  
and temperature

values obtained by  
ranking halos in thin  
redshift shells,  
identifying scales  
reached at fixed  $dN/dz$   
(# / sq deg / unit z)

Evrard et al 2002



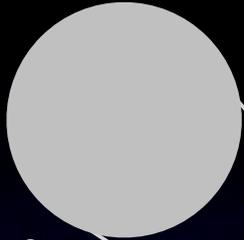




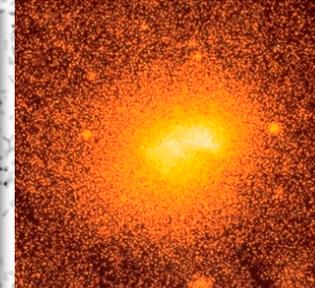
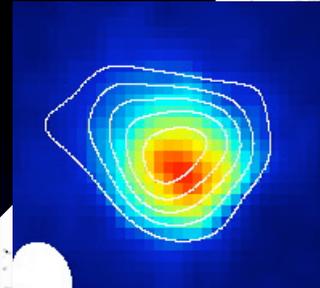
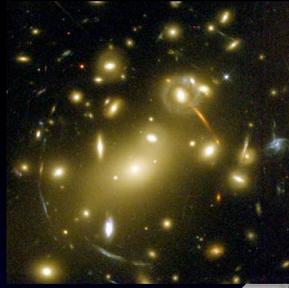
halos = *amplified noise peaks*

$a=0.4$   
( $z=1.5$ )

optical/**lensing** sub-mm **X-ray**

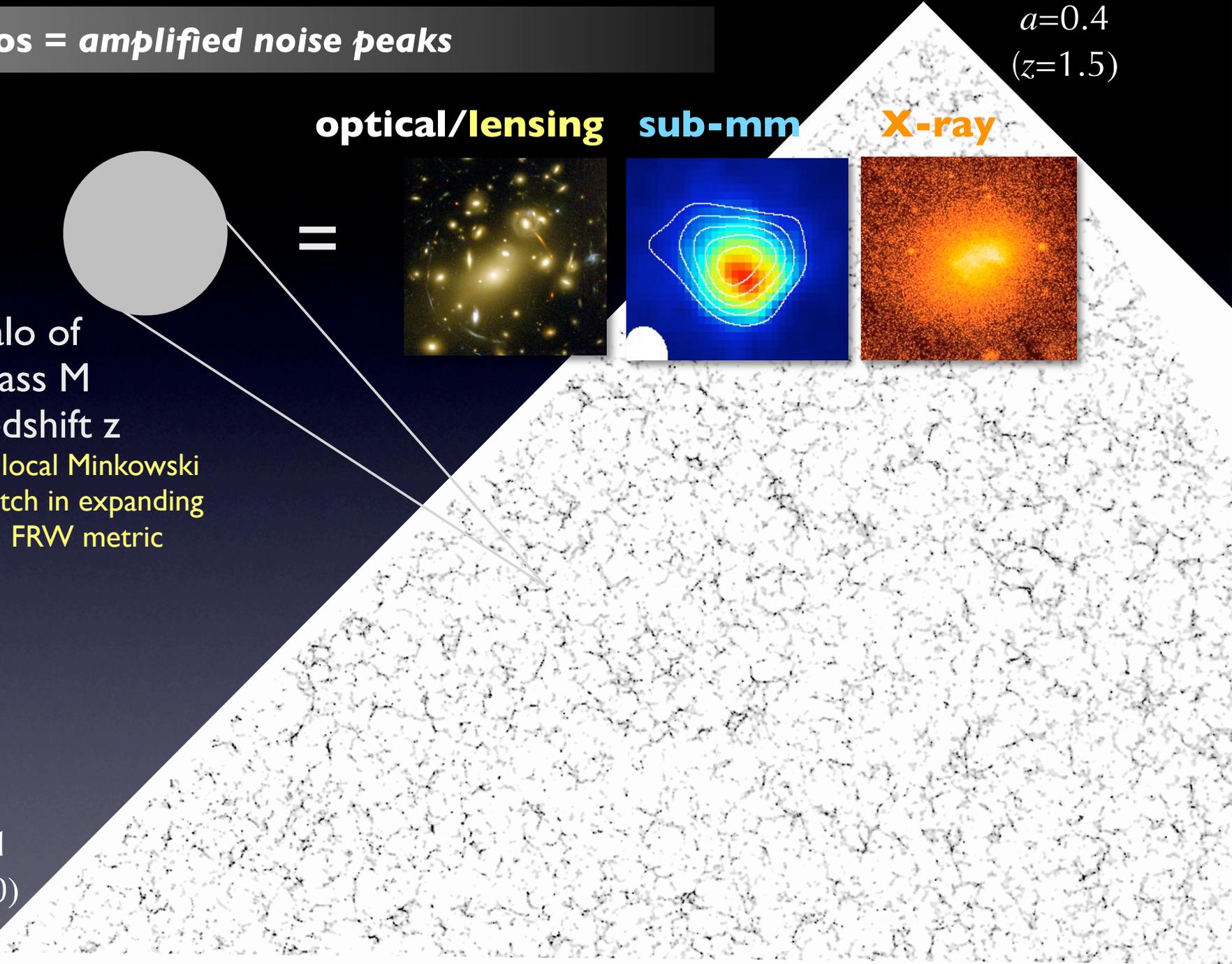


=

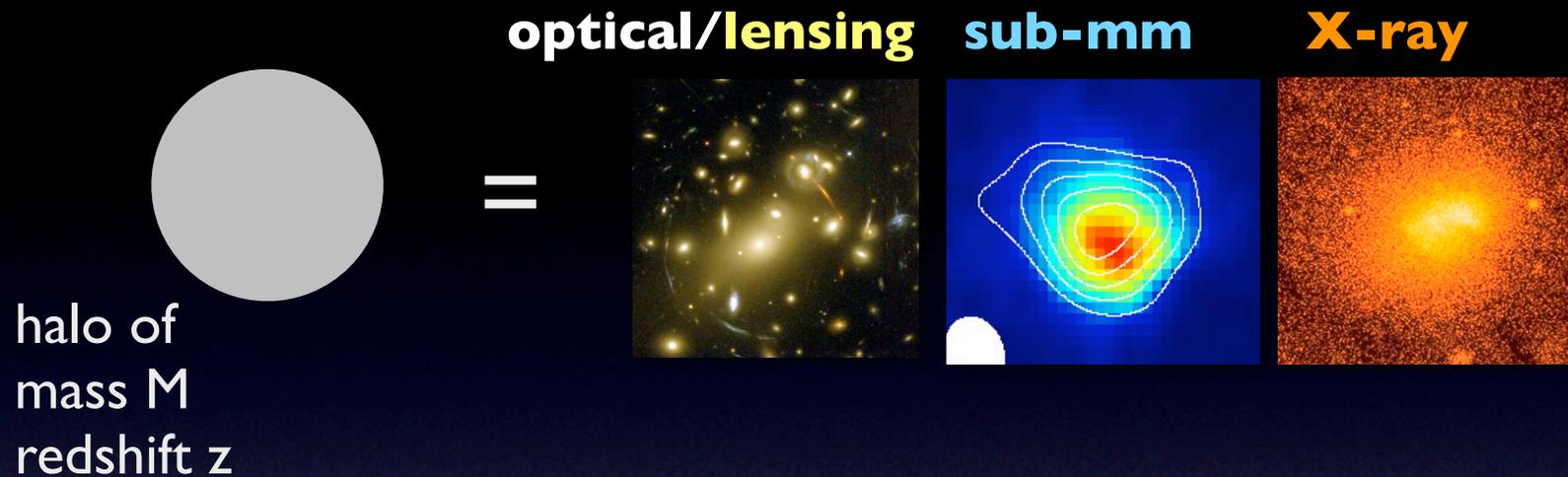


halo of  
mass  $M$   
redshift  $z$   
= local Minkowski  
patch in expanding  
FRW metric

$a=1$   
( $z=0$ )



## massive halo phenomenology: observable signal likelihoods



## “Astrophysics for Dummies”

1. Dimensional analysis  $\Rightarrow$  mean relations are power-laws
2. Central Limit Theorem  $\Rightarrow$  deviations are log-normal

## power-law mean + log-normal covariance model for signals

For  $i^{\text{th}}$  signal, mean behavior of  $s_i = \ln(S_i)$  has slope  $m_i$  in  $\ln M$ .

For  $N$  such signals,

$$\bar{\mathbf{s}}(\mu, \mathbf{z}) = \mathbf{m}(\mathbf{z})\mu + \mathbf{b}(\mathbf{z})$$

$$\mu = \ln M$$

and the halo signal likelihood is

$$p(\mathbf{s} | \mu, \mathbf{z}) = \frac{1}{(2\pi)^{N/2} |\Psi|^{1/2}} \exp\left[-\frac{1}{2} (\mathbf{s} - \bar{\mathbf{s}})' \Psi^{-1} (\mathbf{s} - \bar{\mathbf{s}})\right]$$

with covariance

$$\Psi_{ij} = \left\langle (s_i - \bar{s}_i(\mu, \mathbf{z})) (s_j - \bar{s}_j(\mu, \mathbf{z})) \right\rangle$$

## local approach to the signal space density

Take a (locally) power-law mass function

$$n(\mu) = A \exp(-\alpha\mu) \quad ; \quad \alpha = \alpha(\mu, z)$$

and convolve it with the signal–mass relation to find the  
*signal space density*

$$n(\mathbf{s}) = \frac{A\Sigma}{(2\pi)^{(N-1)/2} |\Psi|^{1/2}} \exp \left[ -\frac{1}{2} \left( \mathbf{s}'\Psi^{-1}\mathbf{s} - \frac{\bar{\mu}^2(\mathbf{s})}{\Sigma^2} \right) \right]$$

with mean mass

$$\bar{\mu}(\mathbf{s}) = \frac{\mathbf{m}'\Psi^{-1}\mathbf{s}}{\mathbf{m}'\Psi^{-1}\mathbf{m}} - \alpha\Sigma^2$$

and mass variance

$$\Sigma^2 = \left( \mathbf{m}'\Psi^{-1}\mathbf{m} \right)^{-1}$$

## signal selection effects

Select a halo sample based on some signal,  $s_1$ .

Then the {mass,  $s_2$ } likelihood is Gaussian with covariance

$$\tilde{\Psi} = \begin{bmatrix} \sigma_{21}^2 & \tilde{r}\sigma_{21}\sigma_{\mu 1} \\ \tilde{r}\sigma_{21}\sigma_{\mu 1} & \sigma_{\mu 1}^2 \end{bmatrix}$$

$$\sigma_{21}^2 = m_2^2 (\sigma_{\mu 1}^2 + \sigma_{\mu 2}^2 - 2r\sigma_{\mu 1}\sigma_{\mu 2})$$

$$\sigma_{\mu i} = \sigma_i / m_i \text{ mass scatter}$$

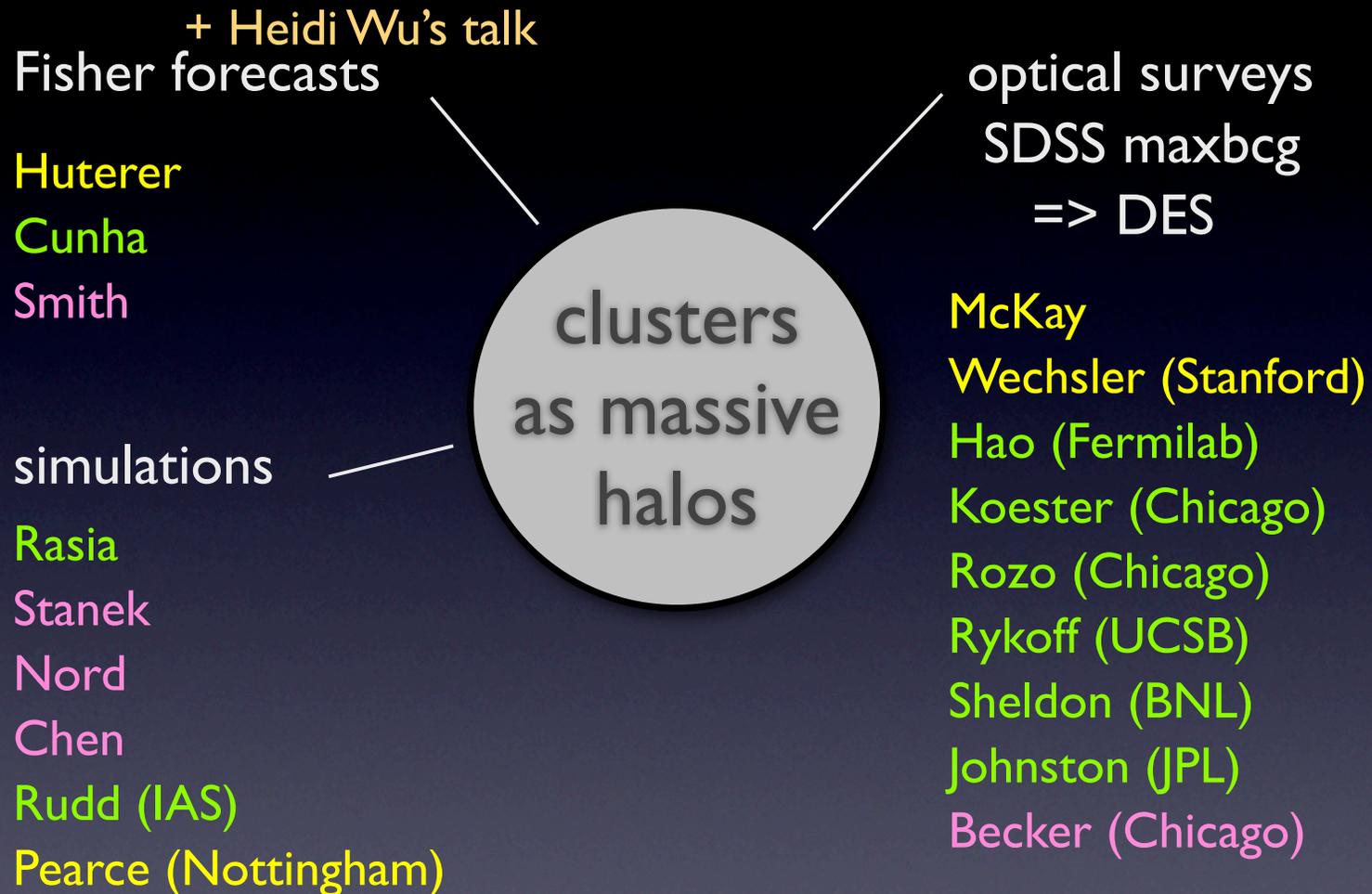
$$\tilde{r} = \frac{(\sigma_{\mu 1} / \sigma_{\mu 2} - r)}{\sqrt{1 - r^2 + (\sigma_{\mu 1} / \sigma_{\mu 2} - r)^2}}$$

and the  $s_2$ -mass scaling for  $s_1$ -binned samples may be biased

$$\bar{s}_2(s_1) = m_2 \left( \bar{\mu}(s_1) + \alpha(\bar{\mu}, z) r \sigma_{\mu 1} \sigma_{\mu 2} \right)$$

$$d\bar{s}_2 / d\bar{\mu} = m_2 \left( 1 + (r \sigma_{\mu 1} \sigma_{\mu 2}) \partial \alpha(\bar{\mu}, z) / \partial \mu \right)$$

## collaborators and topics



Fall 2009: Jeff McMahon (SPT) joins Physics

Winter 2010: Chris Miller (SDSS C4) joins Astronomy

# **SDSS maxbcg analysis**



## SDSS maxbcg cluster sample studies

~13,000 clusters,  $\geq 10$  galaxies

$0.1 < z < 0.3$

based on excess counts of  $g-r$   
red sequence galaxies

Koester et al 2007a,b

follow-up studies:

★ stacked weak lensing masses

Johnston et al 2007; Sheldon et al 2007

★ velocity dispersion–richness

Becker et al 2007

★ X-ray luminosity–richness

Rykoff et al 2008a

★ X-ray luminosity–lensing mass

Rykoff et al 2008a

★ improved richness estimator

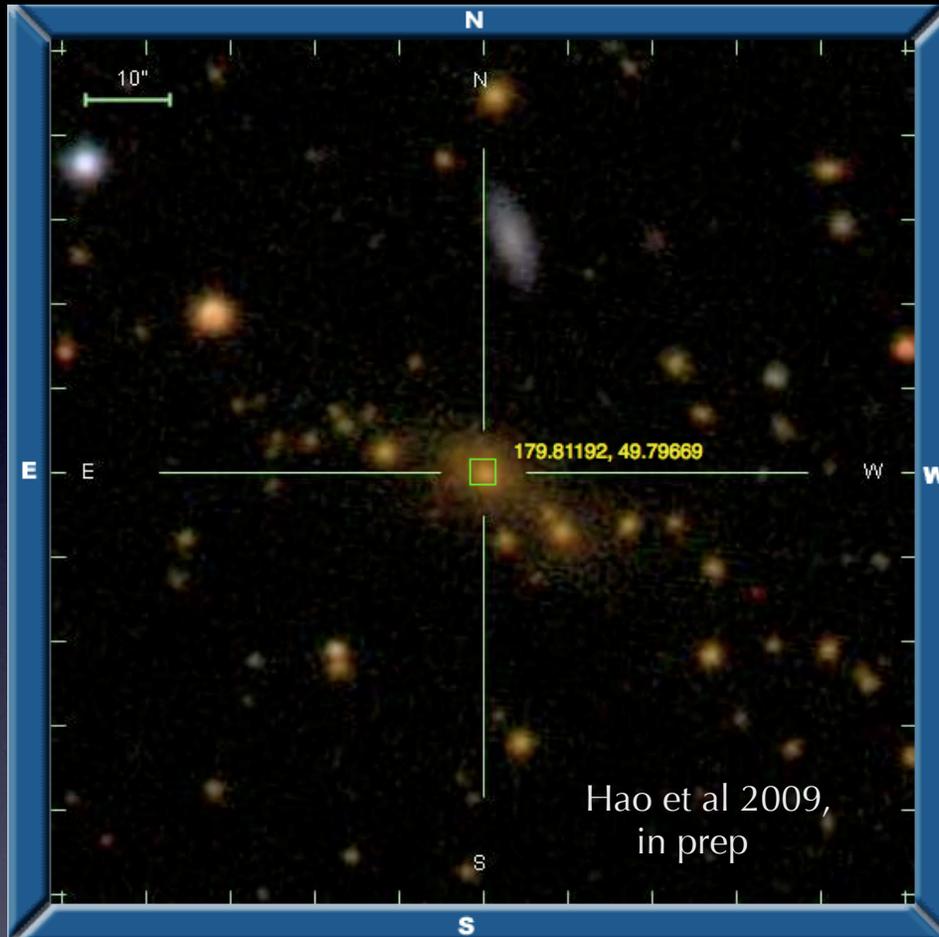
Koza et al 2008a

★ scatter in mass–richness

Koza et al 2008b

★ cosmological constraints

Koza et al 2009



extension to  $r-i$  pushes to higher  $z$

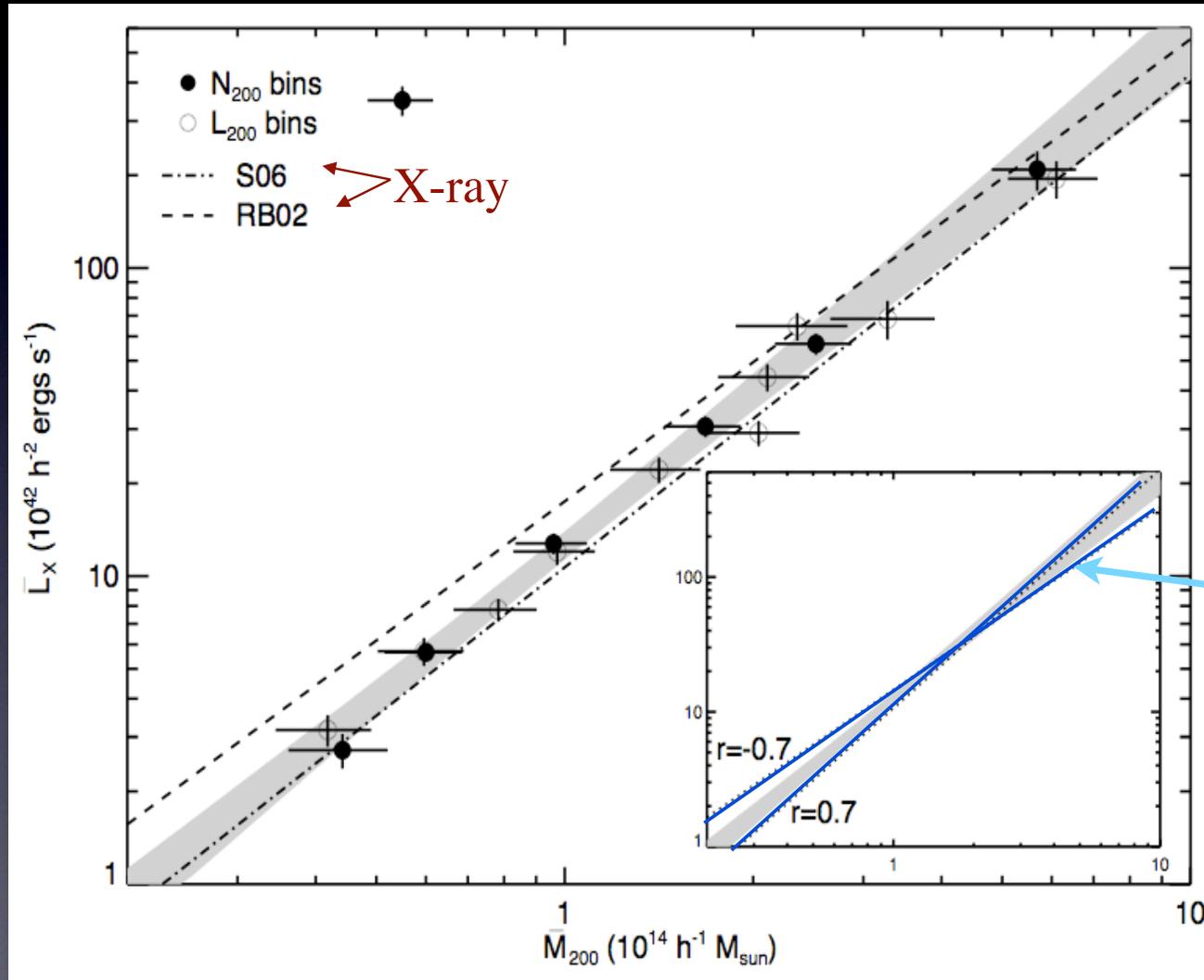
$N_{\text{gals}}=79, z=0.35$

## mean Lx-Mass scaling

17000 clusters,  $N_{\text{gal}} \geq 9$

$M_{200}$  from weak lensing,  $L_x$  from RASS (stacked  $N_{\text{gal}}$  bins)

Johston et al 2007  
Rykoff et al 2008b



Good agreement  
between X-ray and  
optically selected  
samples

slope =  $1.6 \pm 0.1$

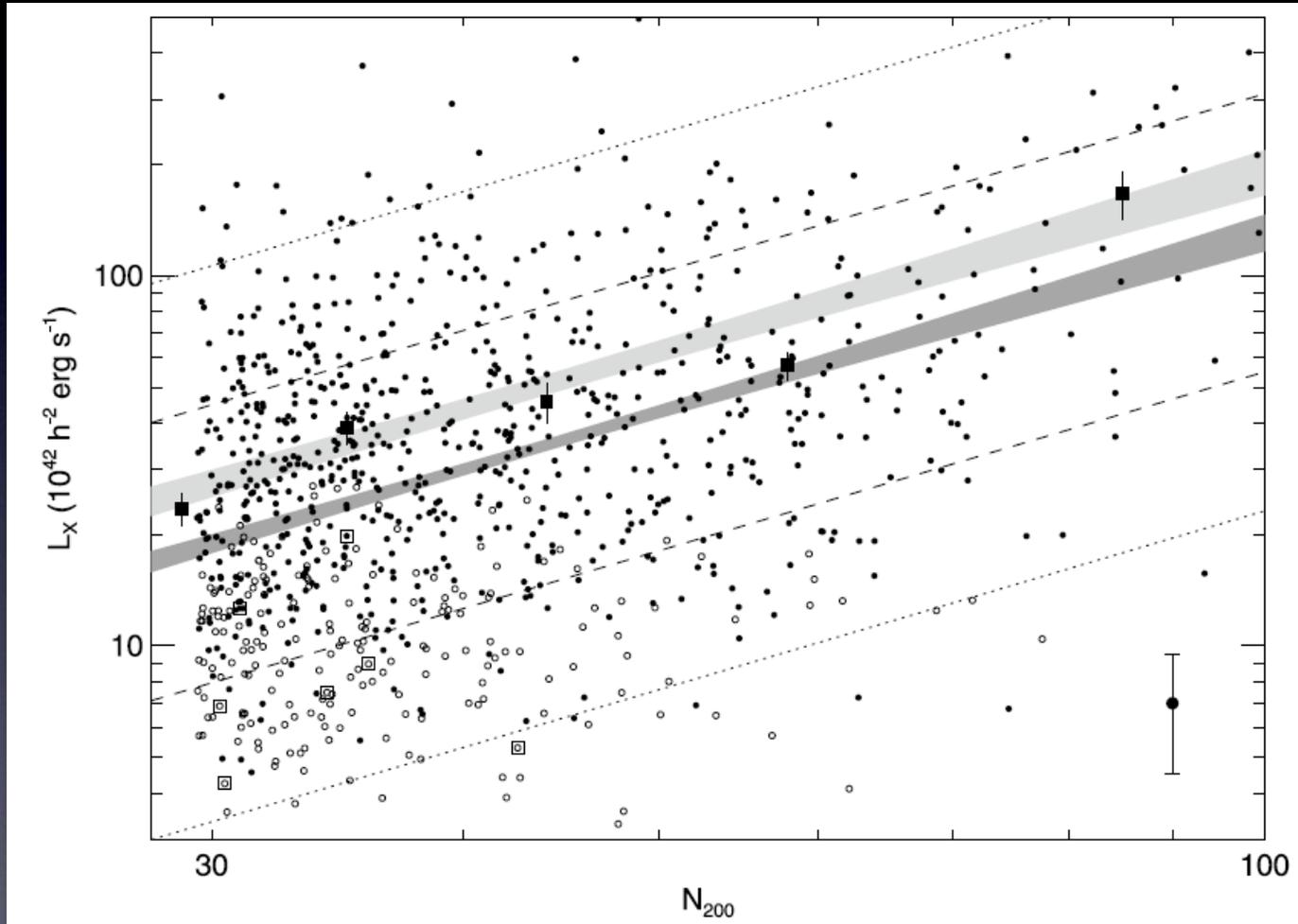
potential *tilt* due to  
optical-X-ray  
correlation and  
running of MF slope

# Lx variance at fixed optical richness

Rykoff et al 2008a

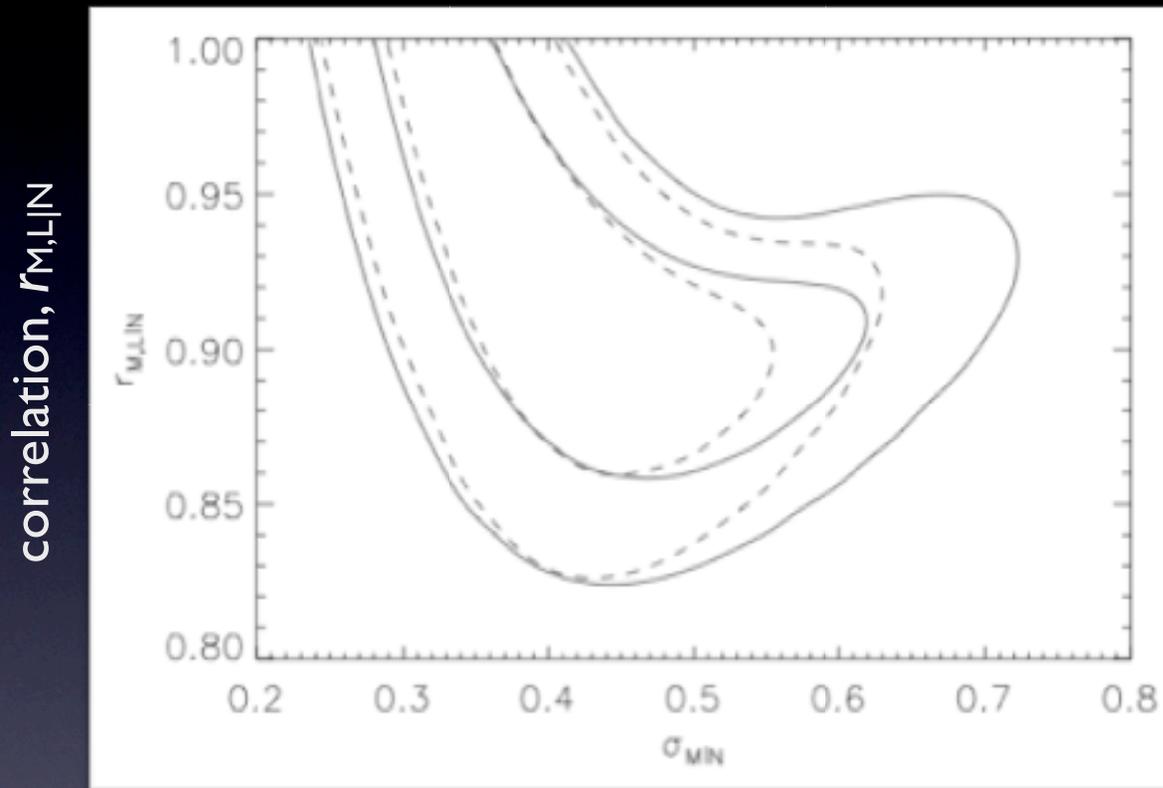
variance in Lx at fixed Ngal

$$\sigma_{\ln L_X | N_{gal}} = 0.83 \pm 0.03$$



# first measurement of property covariance for clusters

Rozo et al 2008b



scatter in  $\ln(\text{mass})$  at fixed  $N_{gal}$

From SDSS-RASS:

- $dn(N_{200})/dN_{200}$
- $L_X - N_{200}$  scaling  
slope, norm, scatter
- $M_{200} - N_{200}$  scaling  
slope, norm

missing:

$M_{200} - N_{200}$  scatter

$M_{200}, L_X | N_{200}$  correlation

Extra information:

400d survey

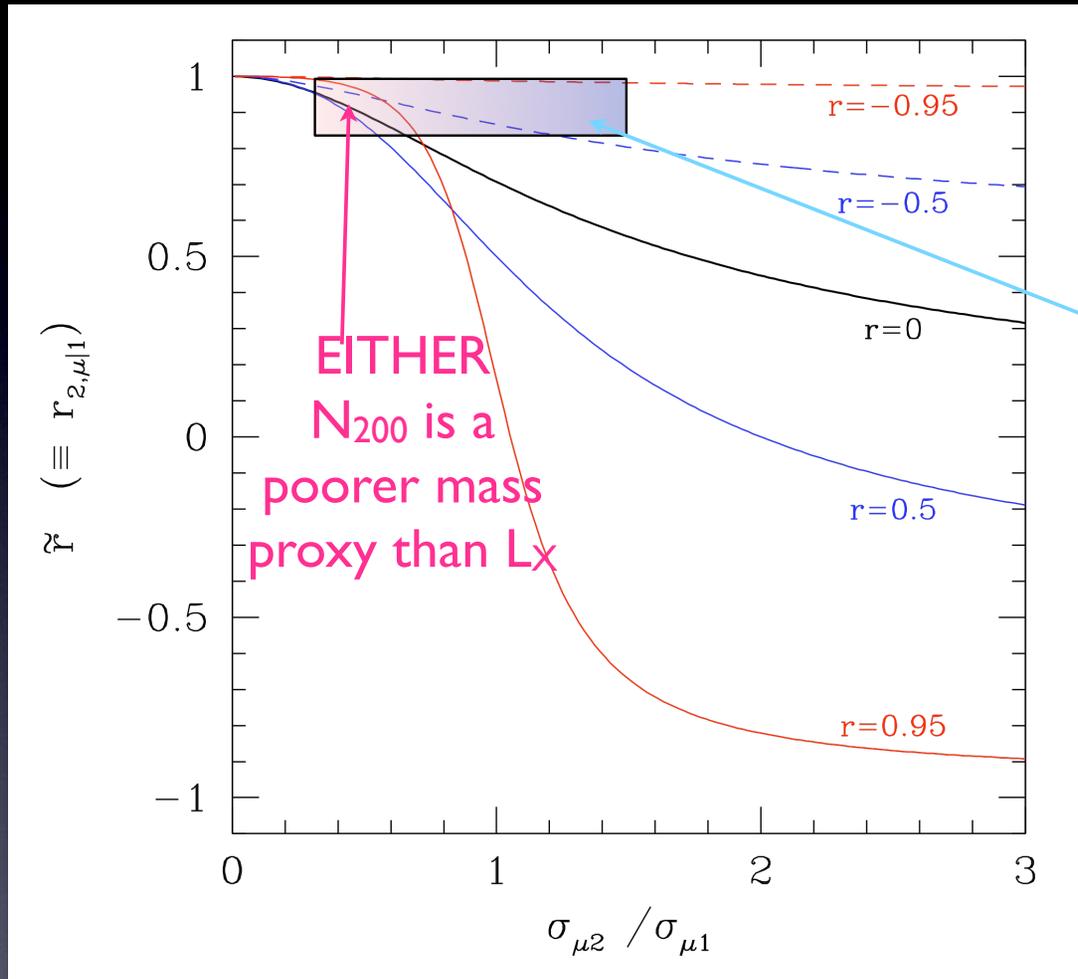
$L_X - M_{500}$  scaling

slope, norm, scatter

Vikhlinin et al 2008

# what does a large covariance in mass and $L_x$ mean?

correlation,  $r_{M,L|N}$



ratio of rms mass variance ( $L_x / N_{gal}$ )

## COSMOLOGICAL CONSTRAINTS FROM THE SDSS MAXBCG CLUSTER CATALOG

EDUARDO ROZO<sup>1</sup>, RISA H. WECHSLER<sup>2</sup>, ELI S. RYKOFF<sup>3</sup>, JAMES T. ANNIS<sup>4</sup>, MATTHEW R. BECKER<sup>5,6</sup>, AUGUST E. EVRARD<sup>7,8,9</sup>, JOSHUA A. FRIEMAN<sup>4,6,10</sup>, SARAH M. HANSEN<sup>11</sup>, JIANGANG HAO<sup>7</sup>, DAVID E. JOHNSTON<sup>12</sup>, BENJAMIN P. KOESTER<sup>6,10</sup>, TIMOTHY A. MCKAY<sup>7,8,9</sup>, ERIN S. SHELDON<sup>13</sup>, DAVID H. WEINBERG<sup>1,14</sup>

*Draft version February 21, 2009*

### ABSTRACT

We use the abundance and weak lensing mass measurements of the SDSS maxBCG cluster catalog to simultaneously constrain cosmology and the richness–mass relation of the clusters. Assuming a flat  $\Lambda$ CDM cosmology, we find  $\sigma_8(\Omega_m/0.25)^{0.41} = 0.832 \pm 0.033$  after marginalization over all systematics. In common with previous studies, our error budget is dominated by systematic uncertainties, the primary two being the absolute mass scale of the weak lensing masses of the maxBCG clusters, and uncertainty in the scatter of the richness–mass relation. Our constraints are fully consistent with the WMAP five-year data, and in a joint analysis we find  $\sigma_8 = 0.807 \pm 0.020$  and  $\Omega_m = 0.265 \pm 0.016$ , an improvement of nearly a factor of two relative to WMAP5 alone. Our results are also in excellent agreement with and comparable in precision to the latest cosmological constraints from X-ray cluster abundances. The remarkable consistency among these results demonstrates that cluster abundance constraints are not only tight but also robust, and highlight the power of optically-selected cluster samples to produce precision constraints on cosmological parameters.

# cosmological constraints from maxbcg counts and lensing

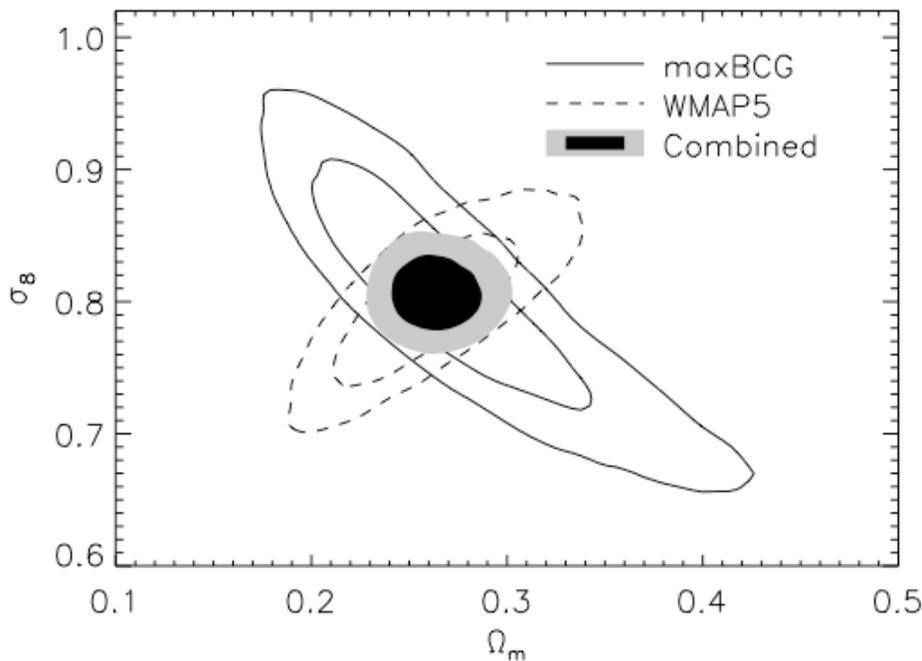
Rozo et al 2009

TABLE 3

Parameter <sup>a</sup>	Prior <sup>b</sup>	Importance <sup>c</sup>
$\sigma_8$	[0.4, 1.2]	unrestrictive
$\Omega_m$	[0.05, 0.95]	unrestrictive
$\langle \ln N_{200}   M_1 \rangle$	flat	unrestrictive
$\langle \ln N_{200}   M_2 \rangle$	flat	unrestrictive
$\sigma_{N_{200}   M}$	[0.1, 1.5]	unrestrictive
$\beta$	$1.00 \pm 0.06$ ; [0.5, 1.5]	restrictive

TABLE 4  
BEST FIT MODEL

Parameter <sup>a</sup>	maxBCG	maxBCG+WMAP5 <sup>b</sup>
$\sigma_8$	$0.804 \pm 0.073$	$0.807 \pm 0.020$
$\Omega_m$	$0.281 \pm 0.066$	$0.269 \pm 0.018$
$\langle \ln N_{200}   M_1 \rangle$	$2.47 \pm 0.10$	$2.48 \pm 0.10$
$\langle \ln N_{200}   M_2 \rangle$	$4.21 \pm 0.19$	$4.21 \pm 0.13$
$\sigma_{N_{200}   M}$	$0.357 \pm 0.073$	$0.348 \pm 0.071$
$\beta$	$1.016 \pm 0.060$	$1.013 \pm 0.059$

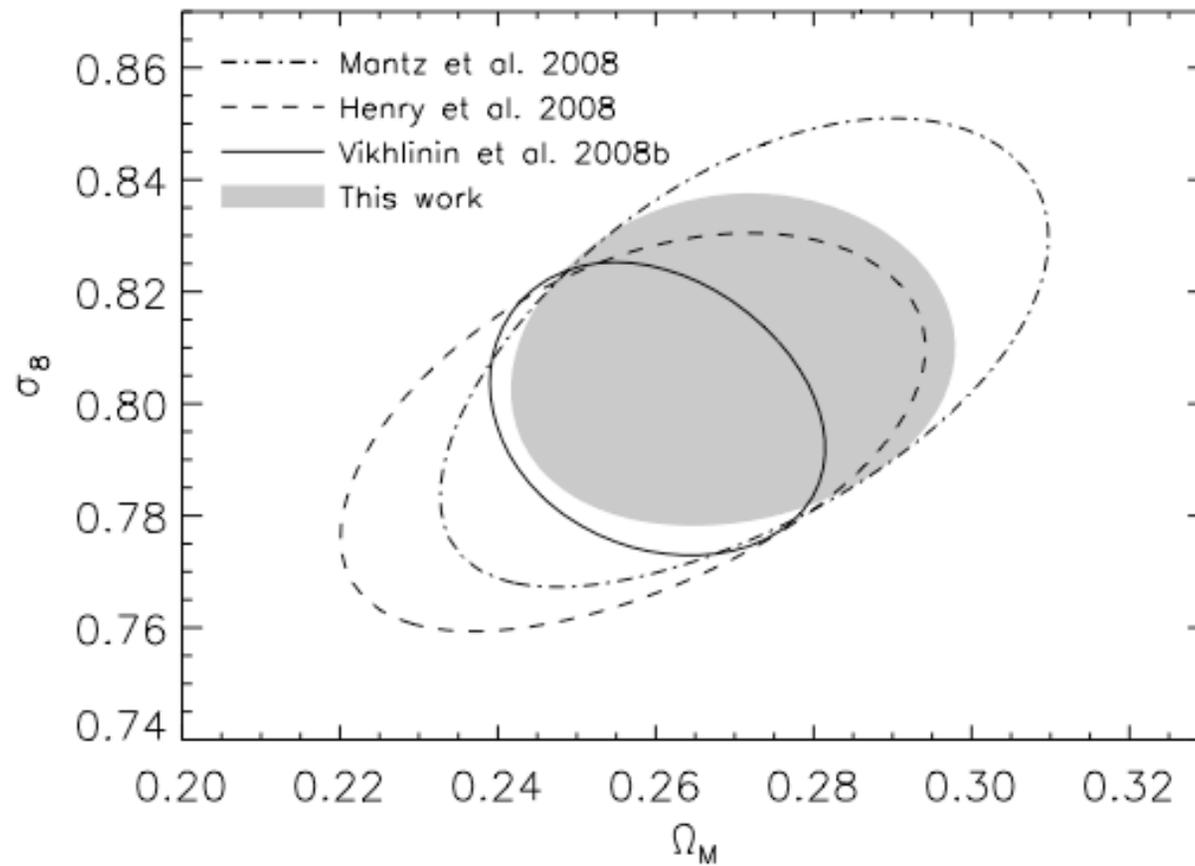


<sup>a</sup>The masses  $M_1$  and  $M_2$  are set to  $1.3 \times 10^{14} M_\odot$  and  $1.3 \times 10^{15} M_\odot$  respectively.

<sup>b</sup>These values are obtained by including the WMAP5 prior  $\sigma_8(\Omega_m/0.25)^{-0.312} = 0.790 \pm 0.024$ . See Section 4.3 for details.

## comparison with X-ray constraints

Rozo et al 2009





## SDSS analysis summary

- ★ red sequence finders identify ~same population as X-rays
- ★ stacked lensing masses + X-ray  $\Rightarrow L_x \sim M^{(1.6 \pm 0.1)}$
- ★ scatter in mass-richness  $\sim 0.45 \pm 0.09$
- ★ beginning to explore covariance in  $L_x$ - $N_{gal}$
  
- ★ cosmological constraints

TABLE 4  
BEST FIT MODEL

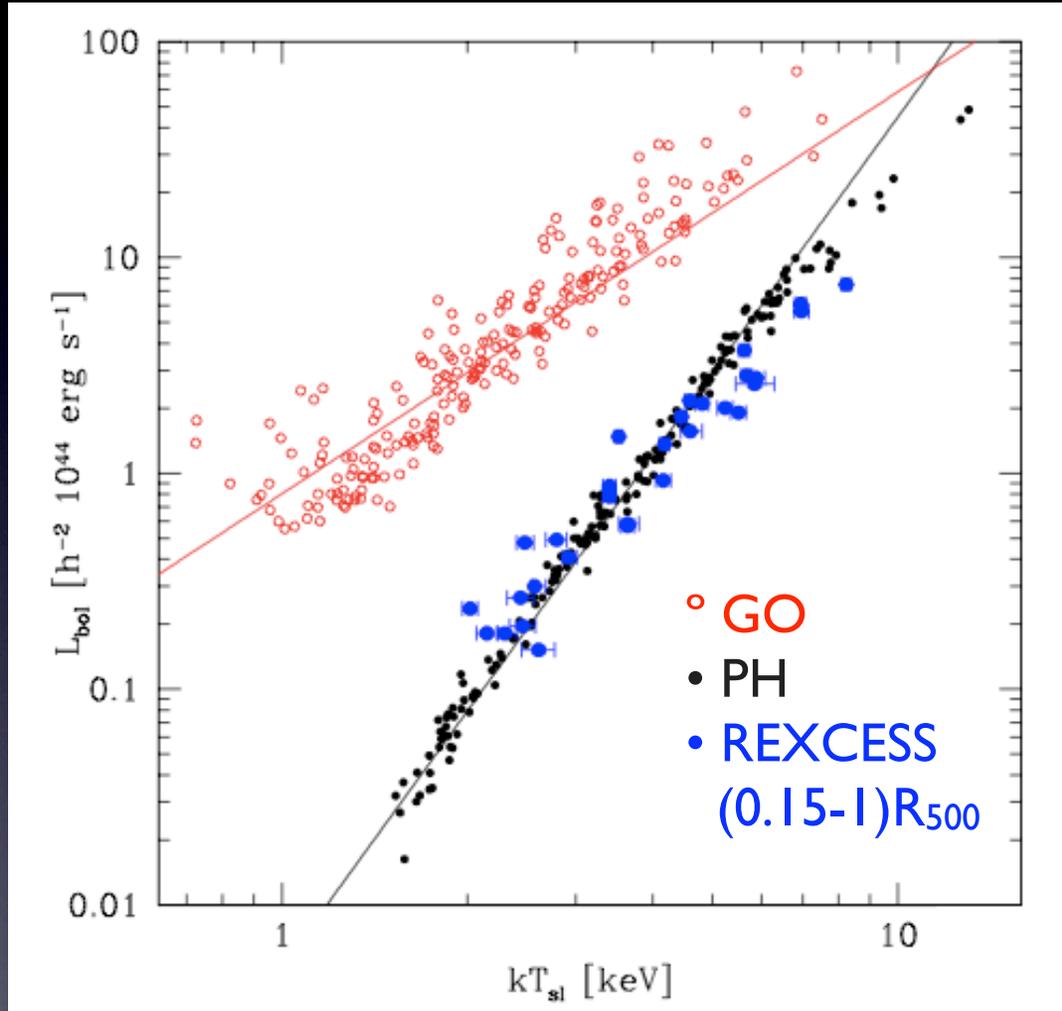
Parameter <sup>a</sup>	maxBCG	maxBCG+WMAP5 <sup>b</sup>
$\sigma_8$	$0.804 \pm 0.073$	$0.807 \pm 0.020$
$\Omega_m$	$0.281 \pm 0.066$	$0.269 \pm 0.018$
$\langle \ln N_{200}   M_1 \rangle$	$2.47 \pm 0.10$	$2.48 \pm 0.10$
$\langle \ln N_{200}   M_2 \rangle$	$4.21 \pm 0.19$	$4.21 \pm 0.13$
$\sigma_{N_{200}   M}$	$0.357 \pm 0.073$	$0.348 \pm 0.071$
$\beta$	$1.016 \pm 0.060$	$1.013 \pm 0.059$

<sup>a</sup>The masses  $M_1$  and  $M_2$  are set to  $1.3 \times 10^{14} M_\odot$  and  $1.3 \times 10^{15} M_\odot$  respectively.

<sup>b</sup>These values are obtained by including the WAMP5 prior  $\sigma_8(\Omega_m/0.25)^{-0.312} = 0.790 \pm 0.024$ . See Section 4.3 for details.

# **Millennium Gas Simulations (MGS)**

# Millennium Gas Simulations



GADGET-2 resimulations  
of Millennium Sim volume

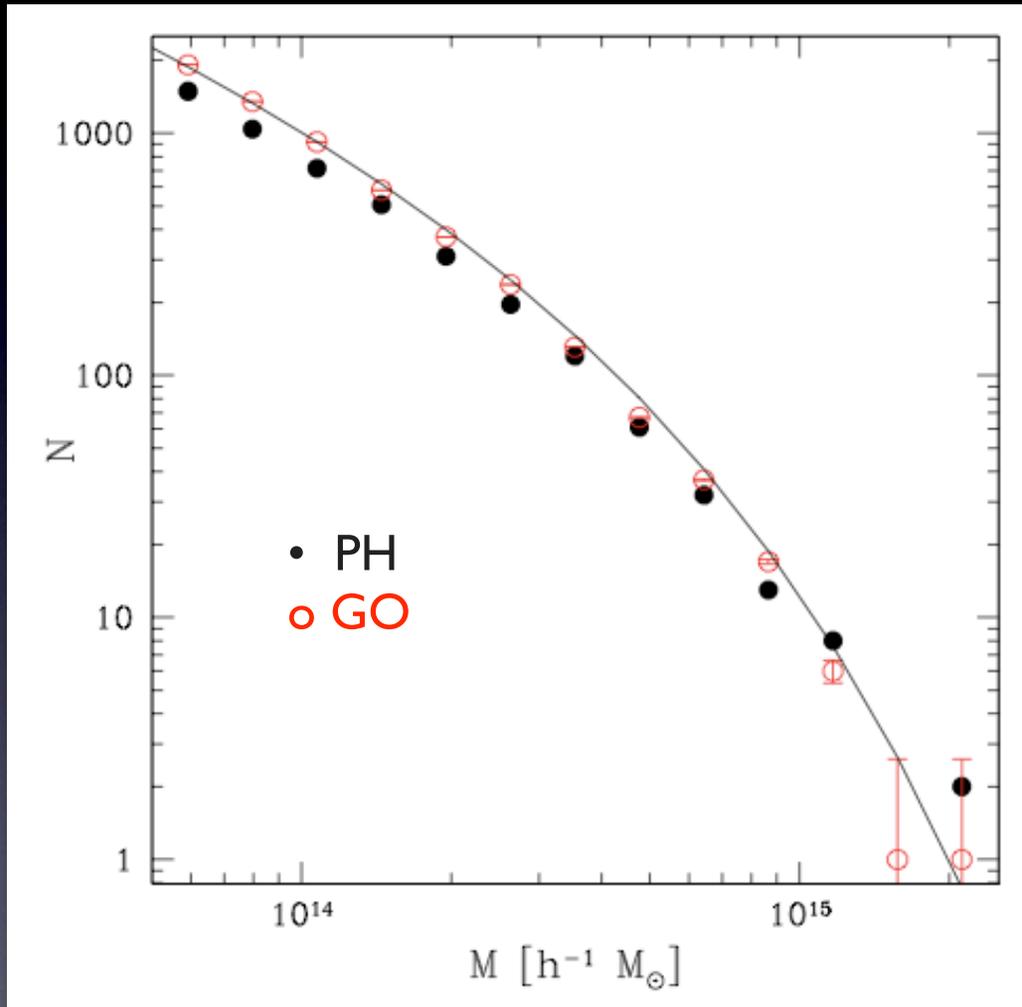
- 500 Mpc/h
- $1e9$  gas+DM particles
- $m_p(\text{DM}) \sim 1.4e10 \text{ Msun}$
- 25 kpc/h softening
- same cosmology as MS

physical treatments:

**GO:** gravity only

**PH:** preheated gas  
200 keV-cm<sup>2</sup> @z=4

## MGS massive halo yield



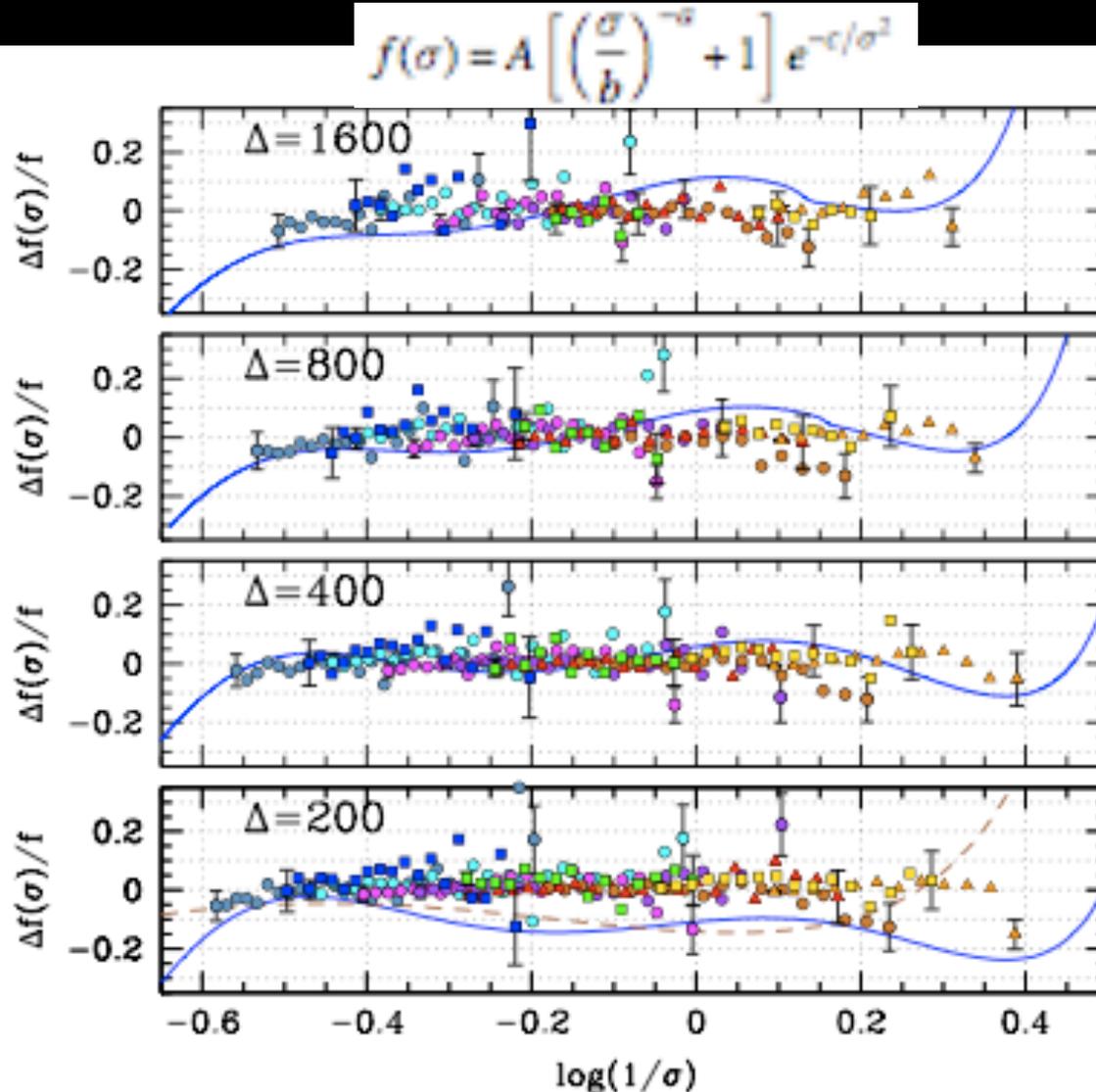
halos at  $z=0$  with  
 $M_{200} \geq 5e13 M_{\text{sun}}/h$ :

4474 (PH)

5612 (GO)

# halo space density from large N-body simulations

Tinker et al (2008)



22 N-body  
simulations with  $N \geq 512^3$

– 5% statistical  
accuracy in counts

– similarity not  
exact in time (need  
 $z$ -factors)

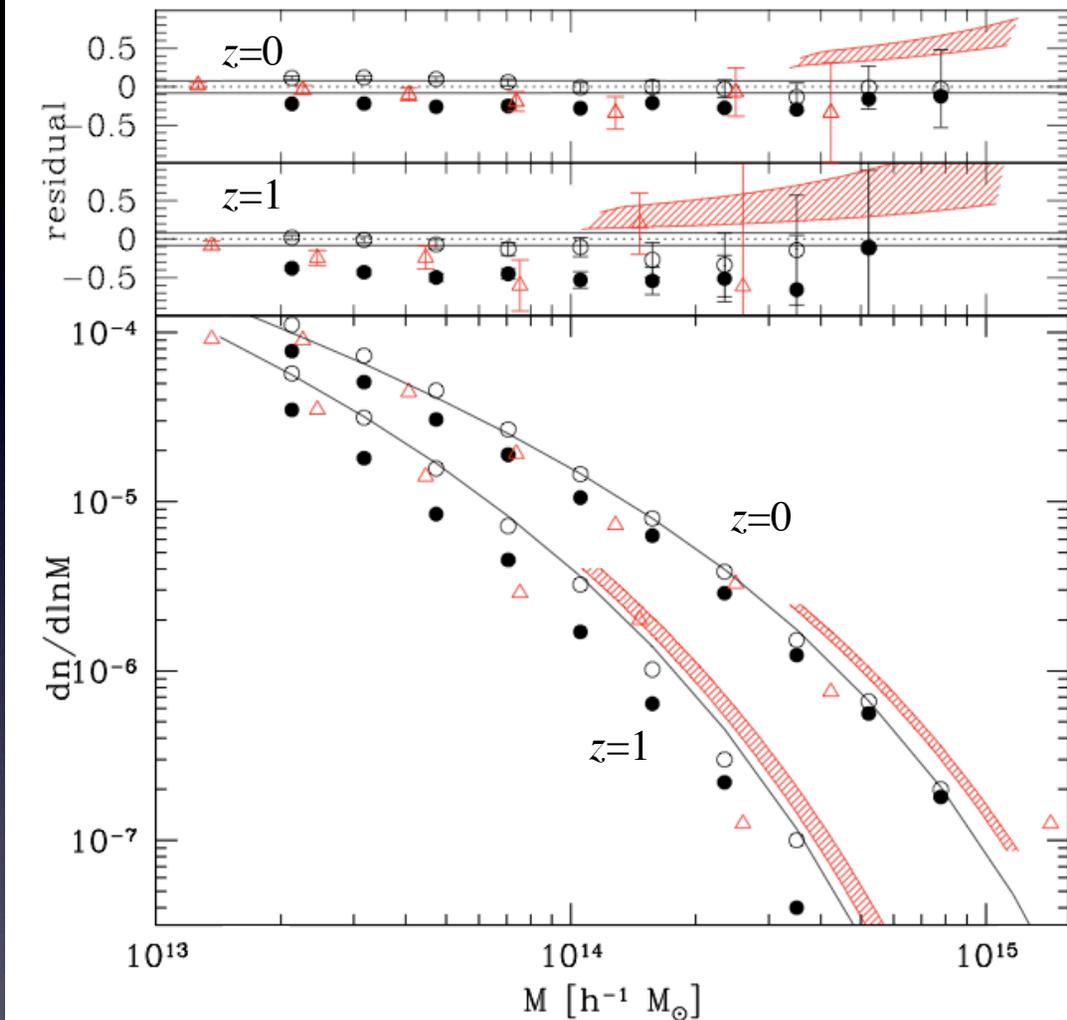
see also:  
Sheth & Tormen 1999  
Reed et al 2000  
Jenkins et al 2001  
Evrard et al 2002  
Hu & Kravtsov 2003  
Warren et al 2006

## sensitivity of halo space density to baryon physics

Stanek et al 2009

– complex baryon physics shifts halo total mass ( $M_{500}$ )

– **maximal** effects are  $>5\%$  statistical error of Tinker et al (2008)



2 pairs of simulations

○ MGS-gravity only

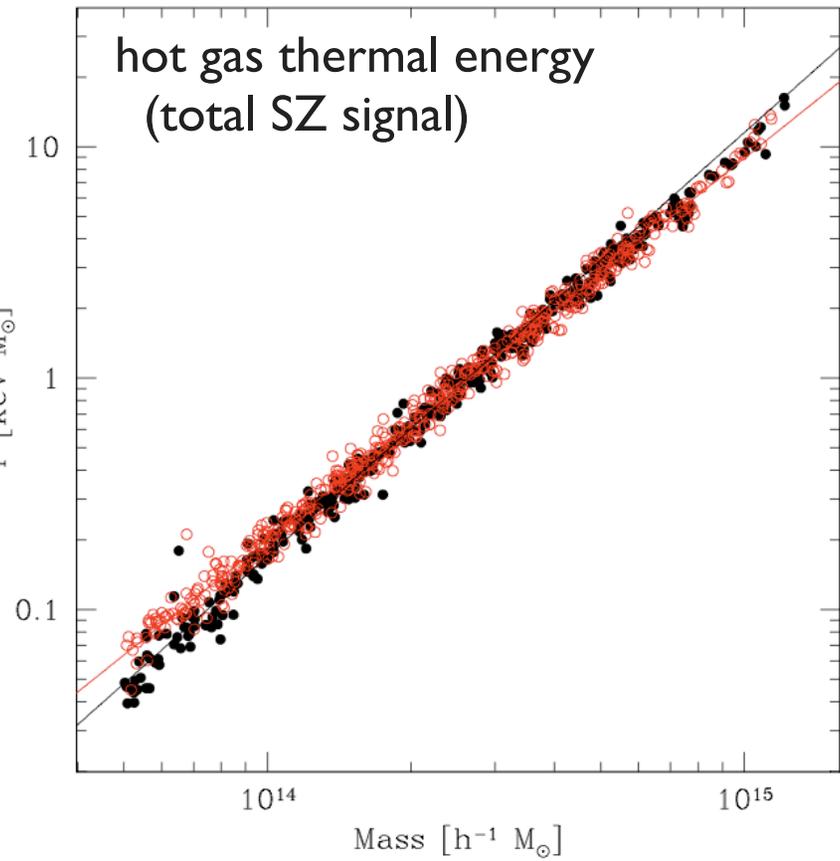
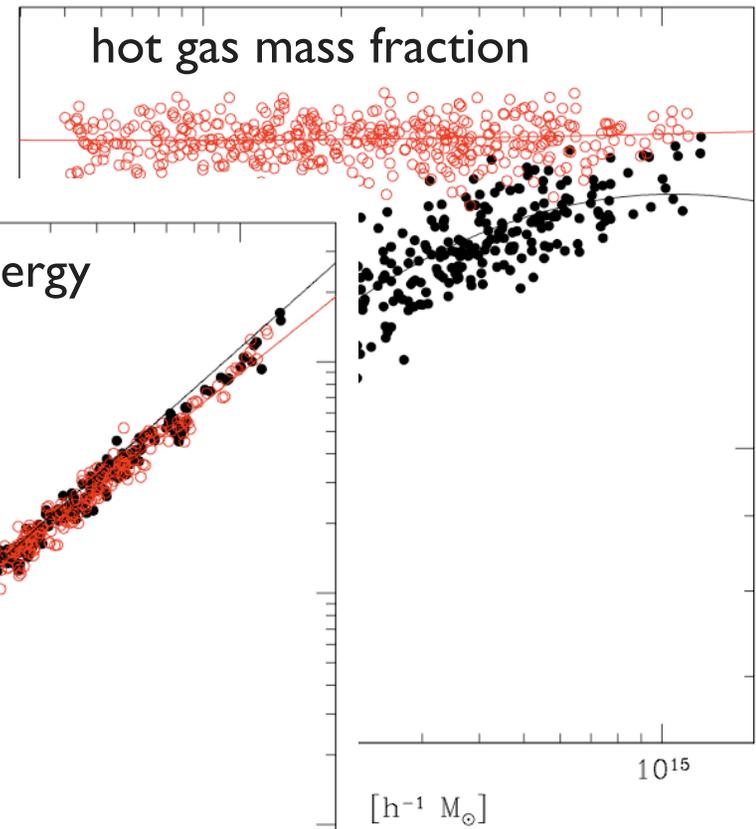
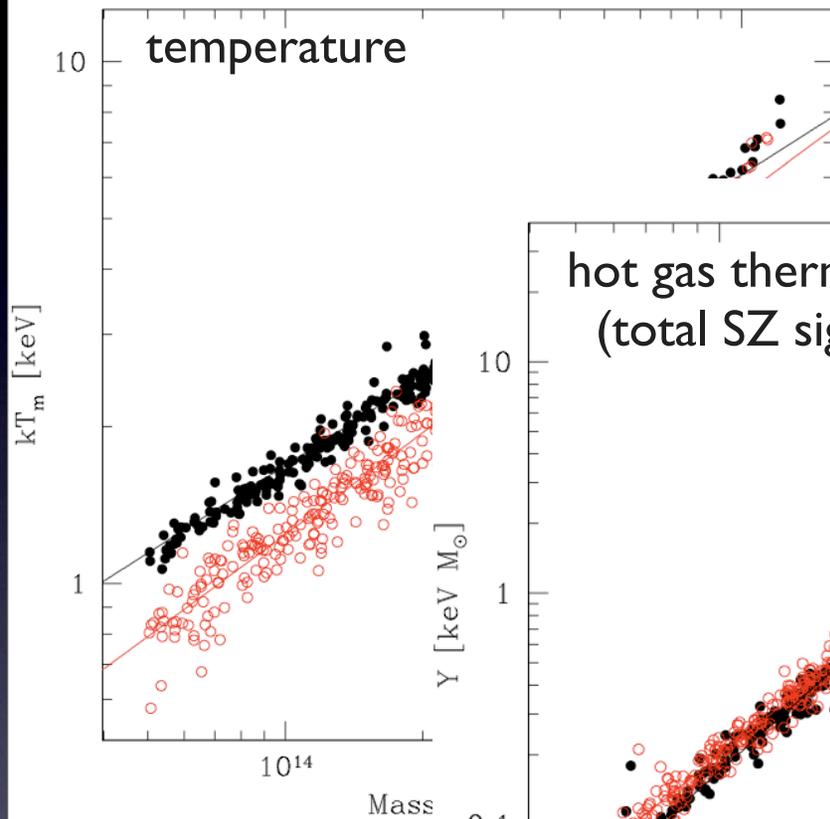
● MGS-preheat

△ ART-gravity only

▨ ART-cool/star/feedback

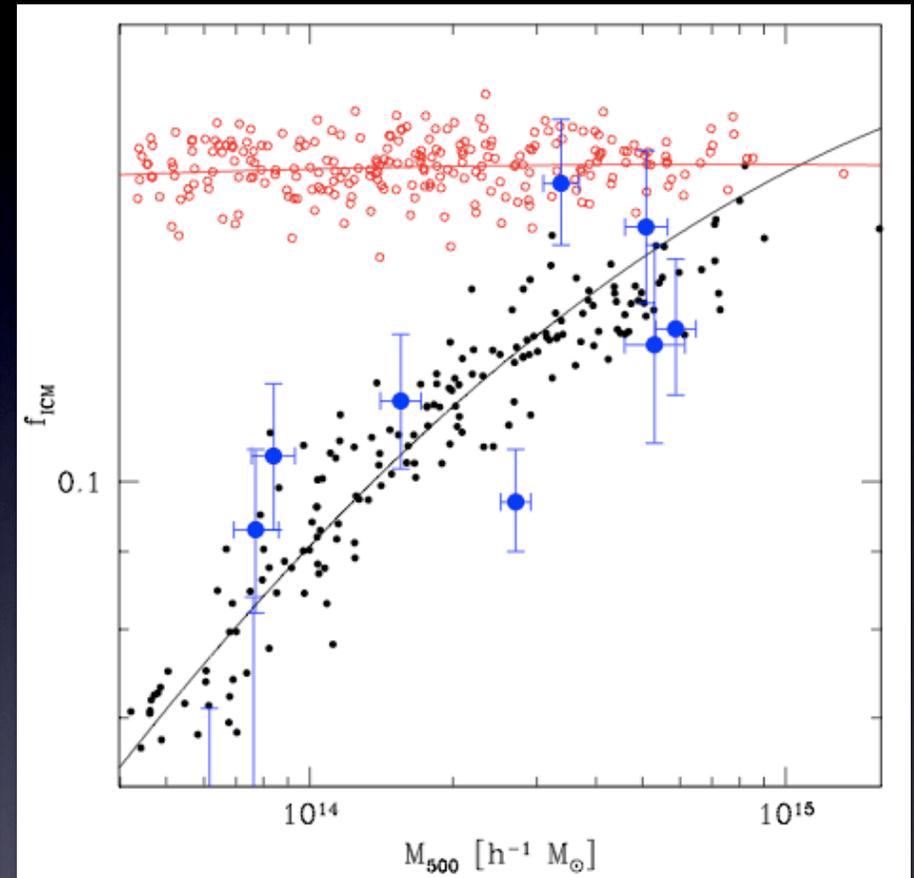
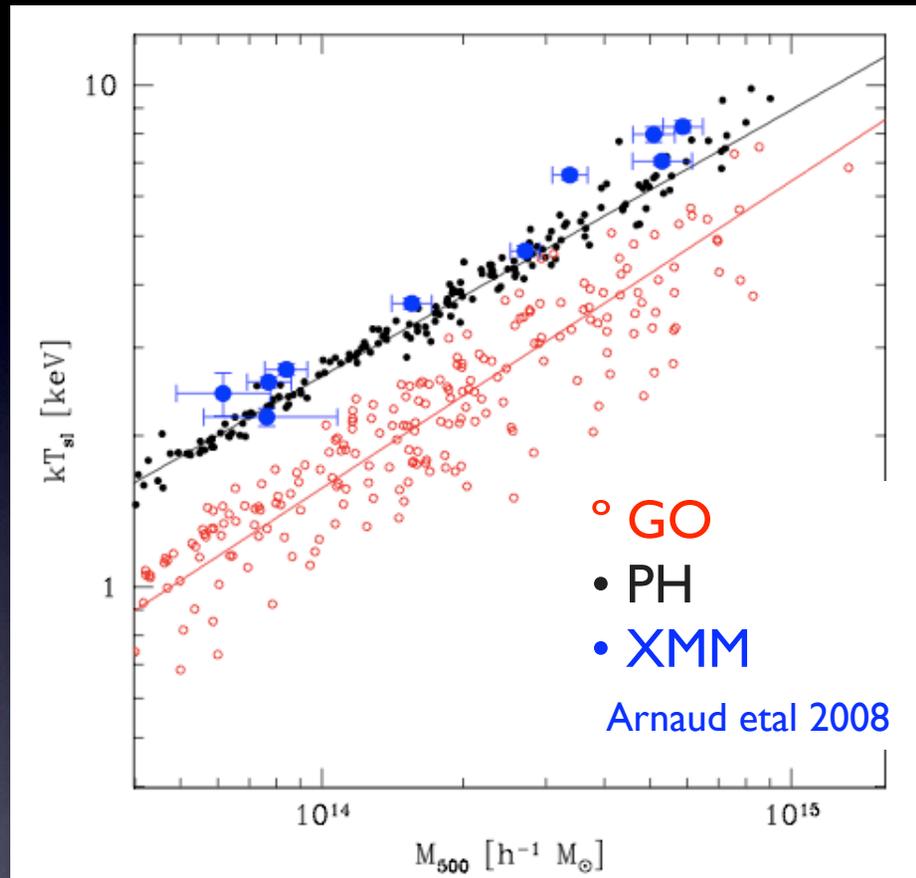
# MGS: z=0 scaling relations w/ different physical treatments

Staneck et al, 0910.1599



# MGS: comparison to observations

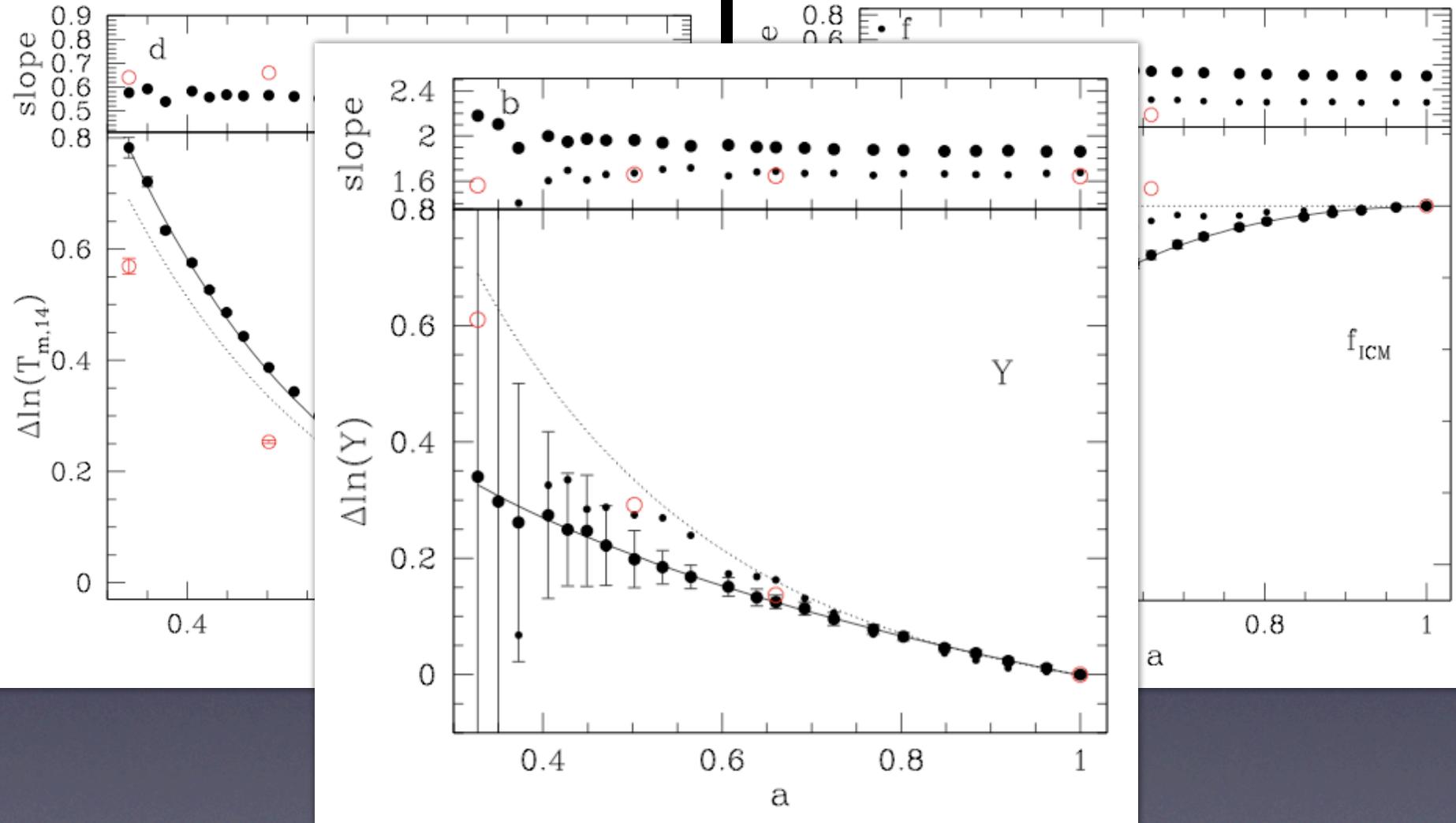
Staneck et al, 0910.1599





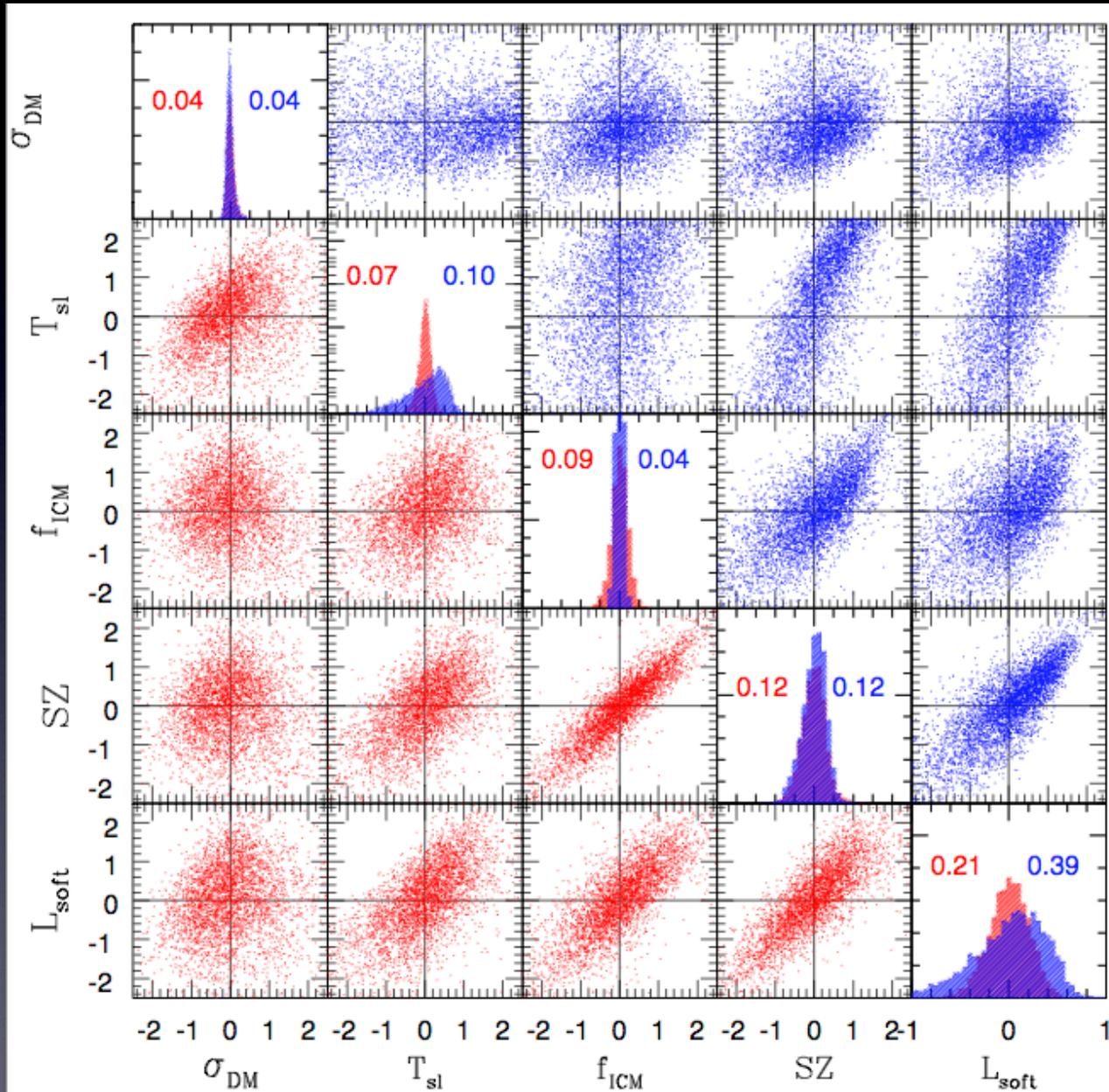
# MGS: evolution of scaling relation slope and intercept

Stanek et al, 0910.1599



# MGS: covariance of multiple signals at fixed halo mass

Staneck et al, 0910.1599



preheating  
gravity only

## effective mass scatter using pairs of signals

$$\Sigma^{-2} = (1 - r^2)^{-1} (\sigma_{\mu 1}^{-2} + \sigma_{\mu 2}^{-2} - 2r\sigma_{\mu 1}^{-1}\sigma_{\mu 2}^{-1}).$$

TABLE 6  
MASS SCATTER AT REDSHIFT ZERO <sup>a</sup>

Cluster Property	$\sigma_{DM}$	$T_{sl}$	$f_{ICM}$	$Y$	$L$	PH	GO
$\sigma_{DM}$	—	0.12	0.12	0.075	0.12	0.12	0.12
$T_{sl}$	0.10	—	0.35	0.050	0.26	0.12	0.38
$f_{ICM}$	0.11	0.12	—	0.054	0.21	0.28	0.12
$Y$	0.062	0.069	0.041	—	0.056	0.069	0.075
$L$	0.090	0.10	0.093	0.066	—	0.11	0.26

<sup>a</sup> The redshift zero mass scatter for each pair of signals, with the results from the PH simulation in the lower, left-hand half, and the results from the GO simulation in the upper, right-hand half, as in Figure 11. The mass scatter for the individual signal is listed on the right-hand side of the table.

## MGS summary

- ★ preheating offers good match to observed core-excised X-ray emission properties
- ★ halo mass is affected by baryon physics at  $\sim 10\%$  level  $\Rightarrow$   
number density at fixed mass shifts by  $\sim 20-30\%$ ; **need**  
**more large volume simulations with gas physics!**
- ★ preheating causes scale-dependent deviations from self-similar evolution in  $Y$ ,  $L_x$ ,  $T$  and  $f_{\text{cM}}$  (**few % at  $10^{14.5} M_{\text{sun}}/h$** )
- ★ covariance in signal pairs generally positive and stable in  $z$
- ★ pairing of  $f_{\text{cM}}$  and  $Y$  may offer sensitive mass selection (4%)

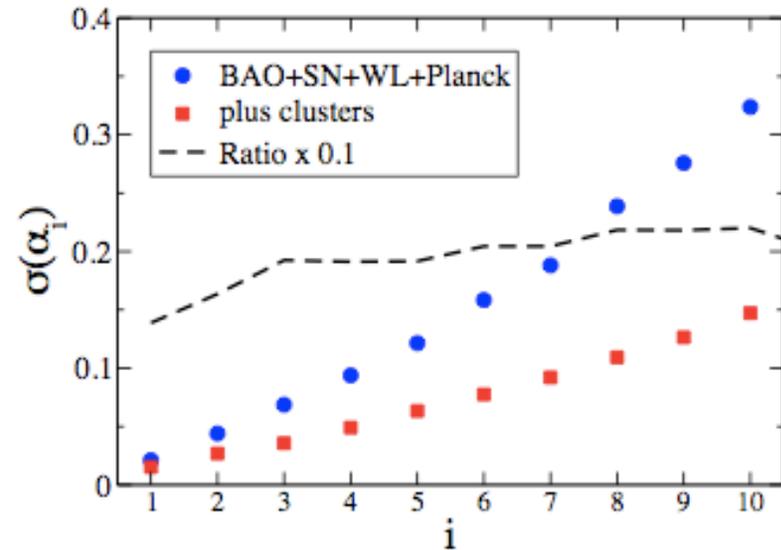
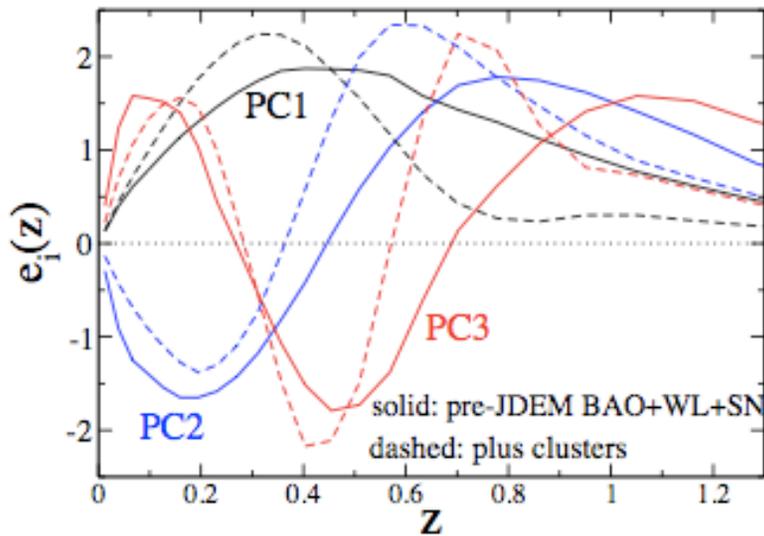
# Fisher forecasts

# Fisher analysis of value of cluster counts and clustering

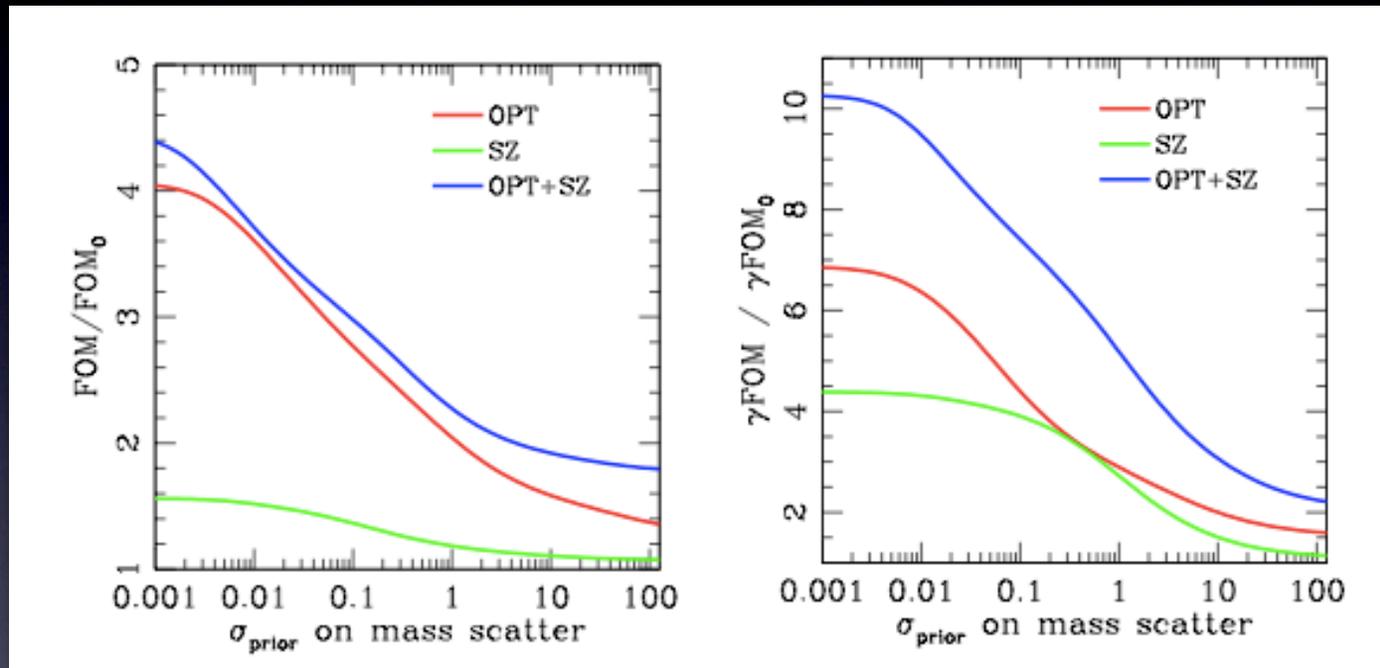
Cunha, Huterer  
Frieman, 0904.1589

$$\ln M^{\text{bias}}(M_{\text{obs}}, z) = \ln M_0^{\text{bias}} + a_1 \ln(1+z) + a_2 (\ln M_{\text{obs}} - \ln M_{\text{pivot}}) \quad (3)$$

$$\sigma_{\ln M}^2(M_{\text{obs}}, z) = \sigma_0^2 + \sum_{i=1}^3 b_i z^i + \sum_{i=1}^3 c_i (\ln M_{\text{obs}} - \ln M_{\text{pivot}})^i \quad (4)$$



## additional improvements from prior in Mobs proxy



# further exploration into Mobs and Mass Function/Bias

Cunha et al (2009)

$$p(M_{\text{obs}}|M) = \frac{1}{\sqrt{2\pi\sigma_{\ln M}^2}} \exp[-x^2(M_{\text{obs}})], \quad (11)$$

where

$$x(M_{\text{obs}}) \equiv \frac{\ln M_{\text{obs}} - \ln M - \ln M_{\text{bias}}(M, z)}{\sqrt{2\sigma_{\ln M}^2(M, z)^2}}. \quad (12)$$

We model systematic error in the mass proxy by introducing a redshift-dependent bias and variance

$$\ln M_{\text{bias}}(z) = B_0 + B_1(1+z), \quad (13)$$

$$\sigma_{\ln M}^2(z) = \sigma_0^2 + \sum_{i=1}^3 s_i z^i, \quad (14)$$

We write the space density of halos as

$$\frac{dn}{dM} = f(\sigma) \frac{\bar{\rho}_m}{M} \frac{d \ln \sigma^{-1}}{dM} \quad (15)$$

and adopt the Tinker parameterization of  $f(\sigma)$  [? ]

$$f(\sigma) = A \left[ \left( \frac{\sigma}{b} \right)^{-a} + 1 \right] e^{-c/\sigma^2}. \quad (16)$$

Following [? ], we allow the first three parameters of  $f(\sigma)$  to vary with redshift, so that

$$A(z) = A_0(1+z)^{A_z} \quad (17)$$

$$a(z) = a_0(1+z)^{a_z} \quad (18)$$

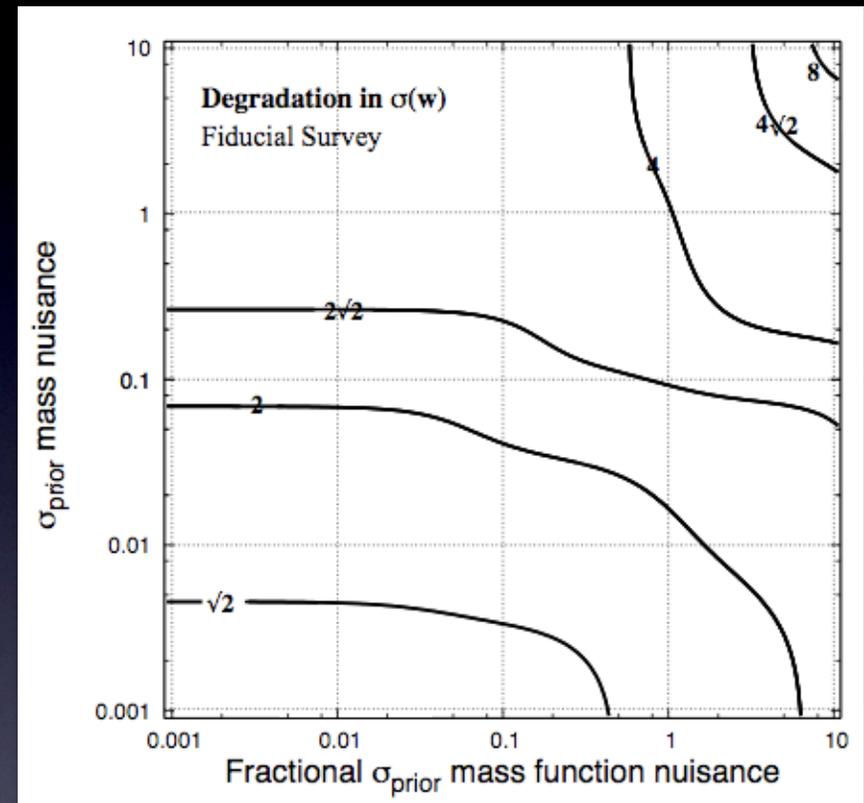
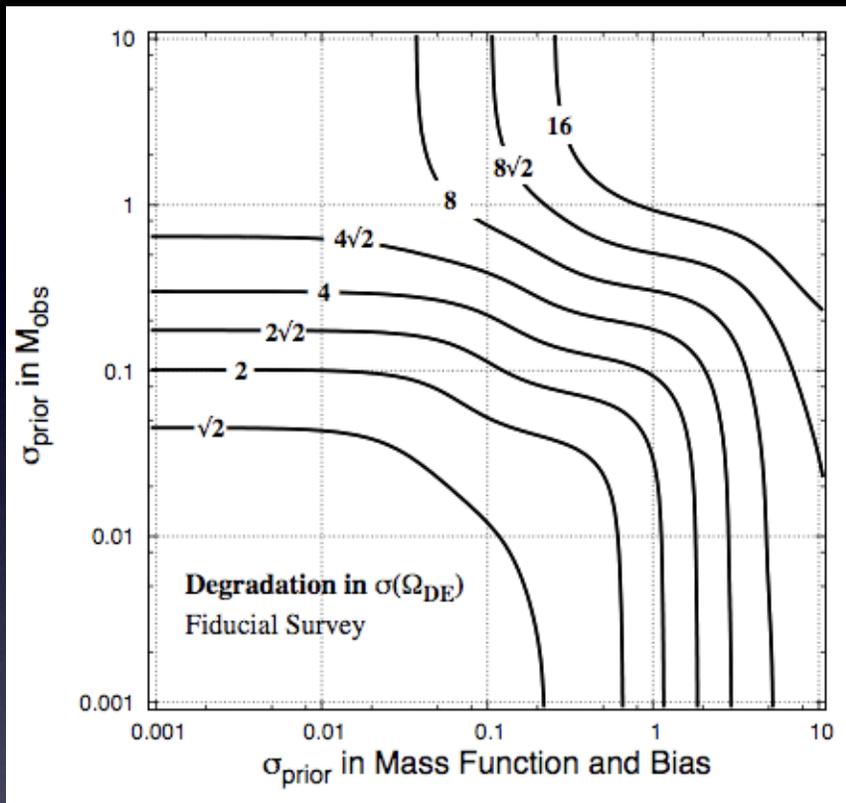
$$b(z) = b_0(1+z)^{-\alpha} \quad (19)$$

TABLE I: Fiducial constraints on cosmological parameters for perfectly known nuisance parameters

Survey	$M_{\text{th}}[h^{-1}M_{\odot}]$	$N_{\text{tot}}$	Sharp priors			No priors	
			$\sigma_0$	$\sigma(\Omega_{\text{DE}})$	$\sigma(w)$	$\sigma(\Omega_{\text{DE}})$	$\sigma(w)$
Fid.	$10^{14.2}$	8,400	0.2	0.010	0.050	0.91	2.19
1	$10^{14.2}$	16,400	0.5	0.0083	0.039	0.82	1.81
2	$10^{13.5}$	229,200	0.2	0.0025	0.011	0.098	0.23
3	$10^{13.5}$	287,200	0.5	0.0023	0.0097	0.22	0.35



# sensitivity to Mobs and Mass Function/Bias priors



**the (near) future...**

## Dark Energy Survey is approaching

### An NSF/DOE-funded study of dark energy using four techniques

- 1) Galaxy cluster surveys (with SPT)
- 2) Galaxy angular power spectrum
- 3) Weak lensing/cosmic shear
- 4) SN Ia distances

### Two linked, multiband optical surveys

5000 deg<sup>2</sup> *g r i z* colors to ~24<sup>th</sup> mag  
Repeated observations of 40 deg<sup>2</sup>

### Development and schedule

Construction: 2007-2011

New 3 deg<sup>2</sup> camera on Blanco 4m, Cerro Tololo

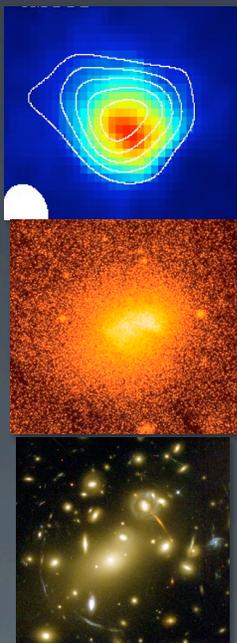
Data management system at NCSA

Survey Operations: 2011-2016

510 nights of telescope time over 5 years

John Peoples, Director

Fermilab, U Illinois, U Chicago, LBNL, U Michigan  
CTIO/NOAO, Barcelona, UCL, Cambridge, Edinburgh



# KITP Workshop

## Galaxy Clusters: The Crossroads of Astrophysics and Cosmology

January 31 – April 22, 2011

Organizers:

Andrey Kravtsov  
Dan Marrone  
Peng Oh

Advisors:

Dick Bond  
John Carlstrom  
Megan Donahue  
Gus Evrard  
Maxim Markevitch  
Mark Voit

